

LINEAR STABILITY IMPLIES NONLINEAR STABILITY FOR FABER-KRAHN TYPE INEQUALITIES

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ABSTRACT. For a domain $\Omega \subseteq \mathbb{R}^n$ and a small number $\mathfrak{T} > 0$, let

$$\mathcal{E}_0(\Omega) = \lambda_1(\Omega) + \mathfrak{T}\text{tor}(\Omega) = \inf_{u, w \in H_0^1(\Omega) \setminus \{0\}} \frac{\int |\nabla u|^2}{\int u^2} + \mathfrak{T} \int \frac{1}{2} |\nabla w|^2 - w$$

be a modification of the first Dirichlet eigenvalue of Ω . It is well-known that over all Ω with a given volume, the only sets attaining the infimum of \mathcal{E}_0 are balls B_R ; this is the Faber-Krahn inequality. The main result of this paper is that, if for all Ω with the same volume and barycenter as B_R and whose boundaries are parametrized as small C^2 normal graphs over ∂B_R with bounded C^2 norm,

$$\int |u_\Omega - u_{B_R}|^2 + |\Omega \Delta B_R|^2 \leq C[\mathcal{E}_0(\Omega) - \mathcal{E}_0(B_R)]$$

(i.e. the Faber-Krahn inequality is *linearly stable*), then the same is true for *any* Ω with the same volume and barycenter as B_R without any smoothness assumptions (i.e. it is *nonlinearly stable*). Here u_Ω stands for an L^2 -normalized first Dirichlet eigenfunction of Ω . Related results are shown for Riemannian manifolds. The proof is based on a detailed analysis of some critical perturbations of Bernoulli-type free boundary problems. The topic of when linear stability is valid, as well as some applications, are considered in a companion paper.

1. INTRODUCTION

For a domain $\Omega \subseteq \mathbb{R}^n$, let

$$\lambda_1(\Omega) = \inf_{u \in H_0^1(\Omega)} \frac{\int |\nabla u|^2}{\int u^2}$$

be the first Dirichlet eigenvalue, and let u_Ω be a nonnegative function with $\int u_\Omega^2 = 1$ attaining the infimum. The Faber-Krahn inequality asserts that

$$(1.1) \quad \lambda_1(\Omega) \geq \lambda_1(B_R),$$

where B_R is a ball with $|B_R| = |\Omega|$, and moreover this inequality is strict unless Ω is a translate of B_R . In recent years, the topic of *stability* for this inequality has received considerable attention: if $\lambda_1(\Omega) - \lambda_1(B_R)$ is small, how closely must Ω resemble B_R ? For a survey of the recent literature on this question, see [8]. In particular, *quantitative* stability refers to inequalities of the form

$$(1.2) \quad d^\alpha(B_R(x_\Omega), \Omega) \leq C[\lambda_1(\Omega) - \lambda_1(B_R)],$$

where d is some notion of distance between sets, $\alpha > 0$ is a positive number, and $x_\Omega = \int_\Omega x$ is the barycenter. Here it is desirable to find the strongest d , as well as the smallest number α , for which such an inequality may hold.

The best result of this type in the literature is due to Brasco, De Philippis, and Velichkov [9], which establishes (1.2) with

$$(1.3) \quad d = d_0(E, E') = |E \Delta E'|$$

and $\alpha = 2$. This theorem is sharp, in the sense that (1.2) is false with $d = d_0$ and $\alpha < 2$. However, it does not fully resolve the issue of quantitative stability, as it seems unlikely that d_0 is the “strongest” distance for which (1.2) holds. To understand why, it is useful to separate the stability question into two “regimes” where it may be studied: we say that (1.1) is *linearly stable* with respect to d, α if (1.2) holds for any Ω whose boundary may be expressed as a C^2 normal graph over ∂B_R , and we say (1.1) is *nonlinearly stable*

with respect to d, α if (1.2) holds for any Ω , unconditionally. For example, the result of [9] is a nonlinear stability result.

It is not difficult to see that d_0 is not the strongest norm for which *linear* stability is valid. Indeed, if $\partial\Omega$ is given by a normal graph g , then d_0 is comparable to the L^2 norm of ξ , while the right-hand side of (1.2) is comparable to $\int_{\partial B_R} \xi L\xi$, where L is a first-order differential operator (this is implicit in the proof of [9], and considered more carefully in the companion paper [3]). In particular, Sobolev norms of ξ up to $\|\xi\|_{H^{1/2}(\partial\Omega)}$ appear to be controllable. Unfortunately, this kind of observation does not easily lead to improved nonlinear stability statements, as it is unclear how to define a stronger Sobolev-type distance for arbitrary sets Ω .

In the nonlinear context, d_0 has a different sort of drawback, which is discussed in [8, Open Problem 3]: it is possible to modify Ω by removing a set of measure 0 in a way which strictly increases its eigenvalue (consider, for example, B_1 vs. $B_1 \setminus \{te_1 : t > 0\}$, a slit domain in \mathbb{R}^2). Any change to Ω which modifies the local capacity of its complement should be reflected in $\lambda_1(\Omega)$. On the other hand, d_0 is clearly insensitive to any such changes.

We propose a distance

$$d_1(E, E') = \sqrt{\int |u_E - u_{E'}|^2}$$

where $u_E, u_{E'}$ are the normalized first eigenfunctions as above extended by zero to be defined on all of \mathbb{R}^n , partially in an attempt to formulate nonlinear stability statements for (1.1) which address these issues. This has advantages and disadvantages in comparison to d_0 . In the nonlinear context, it is sensitive to capacity-zero perturbations of Ω . It also turns out to be specifically relevant to some applications to monotonicity formulas in free boundary theory, which we discuss in [3]. On the other hand, d_1 is a kind of indirect measure of set difference and may be difficult to use in practice; we also do not know whether $d_1(E, B_R)$ controls $d_0(E, B_R)$ for all sets.

The results of this paper, in conjunction with our companion paper [3], establish sharp quantitative stability for the Faber-Krahn inequality with respect to the distance d_1 . The result is proven on the round sphere and hyperbolic space as well as in Euclidean space: we prove in [3, Theorem 1.1] that (1.2) holds with $d = d_1 + d_2$ and $\alpha = 2$ on simply connected space forms. The proof is based on a *selection principle*. This scheme, introduced by Cicalese and Leonardi in [17] in the context of the isoperimetric inequality, has been implemented in various settings in recent years (including [9]) to establish quantitative forms of geometric inequalities. A selection principle has two main steps:

- (1) The reduction step to show that linear stability implies nonlinear stability, based on regularity theory for a penalized functional.
- (2) The local step to prove that linear stability holds.

This paper is dedicated to carrying out Step (1): if (1.1) is linearly stable with $d = d_1$ and $\alpha = 2$, then it is also nonlinearly stable. Step (2) on simply connected space forms is established in [3]. Let us describe how the classical selection principle argument for Step (1) might go in this context, and where the significant challenges arise in this setting. Suppose by way of contradiction that (1.2) fails. We may thus find a sequence of sets Ω_k with $|\Omega_k| = |B_R|$ and $\lambda_1(\Omega_k) - \lambda_1(B_R) \rightarrow 0$, while

$$(1.4) \quad c_k := d_1^2(B_R(x_{\Omega_k}), \Omega_k) \geq k[\lambda_1(\Omega_k) - \lambda_1(B_R)]$$

The goal, then, is to replace each Ω_k with a set U_k that also satisfies (1.4), but such that ∂U_k a smooth normal graph over ∂B_R , in this way contradicting the linear stability assumption. Such a set is obtained by choosing U_k to be a minimizer of a functional that roughly takes the form

$$\mathcal{F}_\tau^0(U) = \lambda_1(U) + \tau \sqrt{c_k^2 + (c_k - d_1(U, B_R(x_U)))^2}$$

for a small parameter $\tau > 0$ and establishing regularity estimates for such minimizers. The second term in this functional forces $d_1(U_k, B_R(x_{U_k}))$ to be close to $d_1(\Omega_k, B_R(x_{\Omega_k}))$ so that (1.4) will be satisfied by U_k , while the first term is hopefully regularizing the boundary of U_k .

The Euler-Lagrange equation for minimizers of \mathcal{F}_τ^0 leads to a free boundary condition along ∂U_k of Bernoulli type, similar to the one studied in [5]. When working with d_0 or other weaker distances, the

second term in \mathcal{F}_τ^0 is of lower order, and can easily be seen to be a minor perturbation to the functional. In particular, the regularity theory of Alt and Caffarelli [5] applies more or less directly (at least in conjunction with more recent literature, such as [20, 18]). In the case of d_1 , however, this is no longer a sufficient heuristic: the second term of \mathcal{F}_τ^0 is now a *critical*, or same-order, perturbation of $\lambda_1(U)$. The key point to establish regularity becomes controlling the distance $d_1(U', U)$ between a minimizing set U and various competitors U' . This amounts to needing estimates like

$$(1.5) \quad \int |u_U - u_{U'}| \leq C|\lambda_1(U) - \lambda_1(U')|,$$

at least for the competitors U' necessary to obtain estimates on U . These competitors were essentially laid out in [5], and closely follow minimal surface theory; they consist of $U \cup B_r(x)$, $U \setminus B_r(x)$, and $\phi_t(U)$ for smooth diffeomorphisms ϕ_t .

We do not know if (1.5) is valid for these competitors, and believe this is an interesting problem. This means that we are unable to establish the necessary regularity theory for \mathcal{F}_τ^0 to carry out Step (1) of the selection principle as described above. In order to overcome this challenge, we *modify the functional*: let

$$(1.6) \quad \mathcal{F}_\tau^*(U) = \lambda_1(\Omega) + \mathfrak{T}\text{tor}(\Omega) + \tau\sqrt{c_k^2 + (c_k - d_1(U, B_R(x_U)))^2},$$

where $\mathfrak{T} > 0$ is a small parameter and

$$\text{tor}(\Omega) = \int \frac{1}{2}|\nabla w|^2 - w$$

is the torsional rigidity of Ω . The torsional rigidity has a long history of appearing when considering spectral optimization problems (see [11, 9]), and shares with λ_1 the ball as the unique volume-constrained minimizer. The functional (1.6) corresponds to a *vectorial* free boundary problem. The crucial benefit of introducing this torsion term is that we *can* establish the analogue of the key estimate (1.5) corresponding to this functional:

$$(1.7) \quad \int |u_U - u_{U'}| \leq C(|\lambda_1(U) - \lambda_1(U')| + |\text{tor}(U) - \text{tor}(U')|)$$

for all the relevant competitors U' . This is shown in Proposition 3.9 and is the starting point toward establishing an existence and regularity theory for \mathcal{F}_τ^* and its generalizations, which constitutes the core analysis of this paper.

Theorem 1.1. *There exists a $\mathfrak{T}_0 > 0$ such that for each $\mathfrak{T} < \mathfrak{T}_0$, there exists a $\tau_0(\mathfrak{T}) > 0$ such that if $\tau < \tau_0$:*

- (1) *There exists a minimizer U of \mathcal{F}_τ^* over the class of all open sets U with $|U| = |B_R|$.*
- (2) *This U has $C^{2,\alpha}$ boundary for any $\alpha < 1$, and ∂U may be expressed as a normal graph ξ over ∂B_R with $\|\xi\|_{C^2} = o_\tau(1)$ (i.e. for every $\varepsilon > 0$, there is a $\tau(\varepsilon)$ such that if $\tau < \tau(\varepsilon)$, then $\|\xi\|_{C^{2,\alpha}} \leq \varepsilon$).*

Beyond establishing the key estimate (1.7), further difficulties arise in proving Theorem 1.1, which we discuss below. Theorem 1.1 holds for a more general class of critical perturbations of the Alt-Caffarelli functional that is introduced in Section 2.2. One generalization is that Theorem 1.1 holds for the functional (1.6) with d_1 replaced by $\tilde{d}_0 + d_1$ where \tilde{d}_0 is an essentially equivalent regularization of d_0 in (1.3) above.

Returning to the discussion of the selection principle, we use the functional \mathcal{F}_τ^* to establish Step (1) of the selection principle corresponding to the inequality $\lambda_1(\Omega) + \mathfrak{T}\text{tor}(\Omega) \geq \lambda_1(B_R) + \mathfrak{T}\text{tor}(B_R)$ instead of (1.1). This linear stability-implies-nonlinear stability result takes the following form.

Corollary 1.2. *Assume that, for a $\mathfrak{T} < \mathfrak{T}_0$ and $\varepsilon > 0$, the inequality*

$$d_1^2(B_R(x_\Omega), \Omega) + d_0^2(B_R(x_\Omega), \Omega) \leq C[\lambda_1(\Omega) - \lambda_1(B_R) + \mathfrak{T}(\text{tor}(\Omega) - \text{tor}(B_R))]$$

holds for all Ω with $|\Omega| = |B_R|$ whose boundaries may be expressed as a $C^{2,\alpha}$ normal graph over ∂B_R with norm bounded by ε . Then it also holds for all open sets Ω with $|\Omega| = |B_R|$ (with a possibly larger constant).

Although it may appear that we have departed from the central aim of establishing Step (1) of the selection principle for the Faber-Krahn inequality, we show in the companion paper [3] that the statement of Corollary 1.2 actually implies that the same is valid with $\mathfrak{T} = 0$ after combining with the Kohler-Jobin inequality ([27, 28], see also [29, 7]), an inequality that relates $\text{tor}(\Omega) - \text{tor}(B_R)$ to $\lambda_1(\Omega) - \lambda_1(B_R)$. It is also

shown in [3] that (1.1) is linearly stable with $d = d_1 + d_2$ and $\alpha = 2$, which when combined with Corollary 1.2 gives nonlinear stability as well.

While the preceding discussion was carried out in \mathbb{R}^n to more clearly illustrate the concepts in play, the results proved here also apply to Riemannian manifolds, and are stated accordingly (see Theorem 2.13 and Theorem 2.14 for the general forms of Theorem 1.1 and Corollary 1.2 respectively). In particular, Corollary 1.2 is also valid for any manifold for which balls are the isoperimetric sets (i.e. they are the unique sets to attain the minimum in the isoperimetric inequality), or with suitable modifications to the statement, on any manifold on which minimizers to (1.1) are unique up to isometry. The former condition is restrictive and only known for a handful of manifolds (see [34] for a survey of the known results), but includes the round sphere and hyperbolic space with standard metrics, which will be studied in [3] in greater detail:

Remark 1.3. Corollary 1.2 holds on hyperbolic space and the round sphere, with the barycenter x_Ω replaced by the set centers defined in Examples 2.8 and 2.9 respectively.

The outline of our approach and the structure of the rest of the paper is as follows. In Section 2 we give careful definitions of the quantities discussed here on manifolds and set up notation for everything to follow. In Section 2.4, we discuss how to pass from Theorem 1.1 to Corollary 1.2 (and state it more carefully and in greater generality). In Section 3 we establish basic properties of the first eigenvalue and the torsional rigidity and discuss Faber-Krahn inequalities on manifolds to the extent needed. The introduction of the torsional rigidity is the first key idea in the paper, and in Section 3 we also prove (1.7). As previously discussed, this inequality will be necessary to prove even basic estimates.

Typical approaches to Bernoulli free boundary problems involve first proving existence of minimizers, usually via the direct method, and then proving nondegeneracy and growth estimates at the free boundary. Due to the presence of the eigenfunction penalization term, which is of the same order as the other terms in the functional, a direct approach to the existence of minimizers fails as the functional is not lower semicontinuous with respect to topologies for which compactness is available. For this reason, we must introduce a new approach to proving existence. In Section 4 we prove a priori nondegeneracy estimates for the free boundary for inward minimizers. In Section 5 we prove a priori Lipschitz growth estimates at the free boundary for outward minimizers. In Section 6 we construct a minimizing sequence of outward minimizers and utilize the Lipschitz estimate to prove the limit Ω is a minimizer. This minimizer will then have both nondegeneracy and Lipschitz estimates since it is also both an inward and outward minimizer. One benefit of this approach is that the entire argument takes place in the class of open sets; the a priori estimates allow us to avoid relaxing the problem to the class of quasi-open sets as is common in the literature [12].

From here, we would like to follow the recent approaches of [14, 32] to establish regularity for vectorial free boundary problems. The first step in doing so is to apply the boundary Harnack principle in order to establish a scalar form of the free boundary condition. Substantial difficulties arise here due to the nonlinear term, because the contribution from d_1 to the Euler-Lagrange equation along $\partial\Omega$ is of the same order as that from $\lambda_1(\Omega)$ or $\text{tor}(\Omega)$: it may be expressed as $|\nabla v|^2$ for some function v satisfying an elliptic PDE on Ω . In particular, there is no favorable sign appearing in this term (it is not *elliptic*, in some sense), and so it must be controlled by the other contributions. The key tool toward controlling this term is a careful analysis of the Green's function G_Ω of Ω , which we carry out in Section 7. This involves utilizing an inhomogeneous boundary Harnack principle recently shown in [4].

We begin Section 8 by considering more carefully how d_1 and u_Ω change under smooth deformations of domains. Applying the estimates from the previous section we obtain several forms of the free boundary condition satisfied along $\partial\Omega$ and arrive at a pointwise version. The final fundamental point in the paper is to again apply the inhomogeneous boundary Harnack principle to rewrite the free boundary condition in a form suitable to known methods.

In Section 9 we apply recent free boundary theory to conclude that $\partial\Omega$ is $C^{1,\alpha}$. In Section 10 we apply a higher order inhomogeneous boundary Harnack principle [22] to prove Theorem 1.1 (that $\partial\Omega$ may be parametrized as a $C^{2,\alpha}$ normal graph over the sphere). Appendix A gives further examples of functionals to which our results apply. Finally, Appendix B shows that domains supporting solutions to elliptic PDE with linear growth away from ∂U are nontangentially accessible, a result which is known in the literature but,

to the best of our knowledge, not in the generality required here. This is a technical point relevant in Section 7.

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2. SETUP AND DEFINITIONS

Throughout this paper, we let (M, g) be a complete, n -dimensional smooth Riemannian manifold without boundary that is connected and oriented. We further assume that (M, g) has bounded geometry, that is, $\text{inj}_M := \inf_{x \in M} \text{inj}_x > 0$ and $\sup_{x \in M} |\text{Rm}| < \infty$. All constants will depend on the metric g ; it is likely that these constants are uniform for Riemannian manifolds with uniformly bounded geometry but we do not track this dependence.

Unless otherwise specified, all integrals are taken with respect to the volume measure m associated to g . We use the notation $|\Omega|$ for the volume of a measurable set $\Omega \subseteq M$ and \mathcal{H}^{n-1} for $(n-1)$ -dimensional Hausdorff measure. We let $d(x, y)$ denote the geodesic distance between $x, y \in M$ and let

$$B_r(x) = \{y \in M : d(y, x) < r\}$$

denote a geodesic ball centered at x of radius r . Given a set Ω and a point $x \in M$, we let $d(x, \Omega) = \inf\{d(x, y) : y \in \Omega\}$. We let ∇ and div denote the gradient and divergence with respect to g and let $\Delta u = \text{div } \nabla u$ be the Laplace-Beltrami operator with respect to g .

Let G denote the (possibly trivial) group of isometries of (M, g) . Several global existence results in the sequel will hold for Riemannian manifolds such that either M is compact or M/G is compact; examples of the latter to keep in mind are Euclidean space, hyperbolic space, cylinders $S^{n-k} \times \mathbb{R}^k$ equipped with the standard product metric, and more generally products of space forms.

2.1. Base energy. Given a bounded open set $\Omega \subseteq M$, we may define the *torsional rigidity* of Ω :

$$(2.1) \quad \text{tor}(\Omega) = \inf_{u \in H_0^1(\Omega)} \int_{\Omega} \frac{1}{2} |\nabla u|^2 - u.$$

We note that the definition of the torsional rigidity given here differs by a sign from the typical definition. This is a negative quantity, and it is straightforward to check (using the Sobolev inequality and the Lax-Milgram theorem) that the infimum is finite and uniquely attained for any open bounded $\Omega \neq M$ by a function $w_{\Omega} \geq 0$. We will refer to w_{Ω} as the *torsion function*. Note also that $\text{tor}(\Omega)$ is decreasing with respect to set inclusion.

The *first (Dirichlet) eigenvalue* of Ω is defined as

$$(2.2) \quad \lambda_1(\Omega) = \inf_{u \in H_0^1(\Omega) \setminus \{0\}} \frac{\int_{\Omega} |\nabla u|^2}{\int_{\Omega} u^2}.$$

This infimum is also attained by a nonnegative function u_{Ω} , the *first eigenfunction*, and $\lambda_1(\Omega)$ is also a decreasing function with respect to set inclusion. The first eigenfunction is unique up to scalar multiples so long as $\lambda_1(\Omega) < \lambda_2(\Omega)$, which in particular is true if Ω is connected. The notation u_{Ω} will be used for the unique first eigenfunction with $u_{\Omega} \geq 0$ and $\int_{\Omega} u_{\Omega}^2 = 1$, whenever it exists and is unique. We frequently will extend u_{Ω} and w_{Ω} by zero to be defined on all of M , using the same notation to indicate the functions $u_{\Omega}, w_{\Omega} : M \rightarrow \mathbb{R}$.

In principle, we are interested in minimizing energy functionals of the type (1.6) among sets $\Omega \subset M$ of a fixed volume $|\Omega| = v$. In practice, it is more convenient to replace the volume constraint with a volume penalization term. Following [1, 9], we define the volume penalization

$$(2.3) \quad f_{v,\eta}(t) = \begin{cases} \eta(t - v) & t \leq v \\ \frac{1}{\eta}(t - v) & t > v. \end{cases}$$

The idea of the volume penalization (2.3) is that by choosing the parameter η to be sufficiently small, a set Ω whose volume exceeds the prescribed volume v will have a large energy contribution coming from the term $f_{v,\eta}(|\Omega|)$. We define the *base energy* of Ω as

$$(2.4) \quad \mathcal{E}(\Omega) = \lambda_1(\Omega) + \mathfrak{T}\text{tor}(\Omega) + f_{v,\eta}(|\Omega|),$$

where $\mathfrak{T} \in [0, 1]$. The principal goal of this paper is to study certain critical perturbations of the base energy introduced in the next subsection. In order to do so, we will need to first consider the minimization problem of the base energy itself. We discuss the problem of minimizing $\mathcal{E}(\Omega)$ over sets briefly here and in more detail in Section 3. Many of the existence and regularity results in later sections will apply to this base energy problem as well. Let us begin with two observations regarding the existence theory for this problem:

Remark 2.1 (A hard volume constraint is needed). If we attempt to minimize the base energy \mathcal{E} among all bounded open subsets of M , then the energy may fail to be bounded from below and minimizers may fail to exist or $\Omega = M$ may be a minimizer. For instance, on Euclidean space, a simple scaling computation shows that $\mathcal{E}(B_R) \rightarrow -\infty$ as $R \rightarrow \infty$ for any fixed triple of parameters $v, \eta, \mathfrak{T} > 0$. Similarly, if M has finite volume, then $\mathcal{E}(M) = -\infty$ since $\text{tor}(M) = -\infty$ while the other quantities are finite.

To remedy this issue, we will introduce a “hard volume constraint” that is much larger than v . Given the desired volume constraint $v \in (0, |M|)$, we fix a larger number $v_{\max} \in (v, |M|)$ and minimize with respect to competitors Ω with $|\Omega| \leq v_{\max}$. Ultimately this hard constraint will never be saturated: in Proposition 9.5 we show that for small values of η (depending on v and v_{\max}) and \mathfrak{T} (depending on v, v_{\max} and η), any minimizer over this class actually satisfies the original desired volume constraint $|\Omega| = v$. In particular, this means the volume-constrained problem (minimizing only over $|\Omega| = v$) is equivalent to the volume-penalized problem where $|\Omega| \leq v_{\max}$.

Remark 2.2 (Compactness issues on non-compact manifolds). On noncompact manifolds, a minimizing sequence Ω_k for the base energy \mathcal{E} can “drift to infinity” in the sense that $d(x_0, \Omega_k) \rightarrow \infty$ for any reference point x_0 . In view of this issue, we first minimize over the class of Ω contained in (smoothed out) balls of large radius R . Depending on the geometry of the manifold at infinity, global minimizers may simply fail to exist when $R \rightarrow \infty$. However, when M/G is compact, the compactness issues are solely caused by the symmetries of the metric and we may send the radius R to infinity using concentration compactness techniques (see [31]). This will be addressed in Section 3.4 for the base energy and then more completely in Theorem 6.3.

Since geodesic balls need not have smooth boundary on an arbitrary Riemannian manifold, it is more convenient from a technical standpoint to work with an exhaustion Q_R of smooth bounded open sets that play the role of smoothed-out balls. More specifically, we let Q_R be a fixed 1-parameter family of open sets which have the following properties:

- (1) $\bar{Q}_R \subseteq Q_S$ for any $S > R$
- (2) $\text{diam}(Q_R) \leq 2R$
- (3) ∂Q_R is either smooth (a finite union of disjoint C^2 codimension-1 submanifolds) or empty
- (4) $M = \cup_R Q_R$.

In view of Remarks 2.1 and 2.2, we will fix parameters $R > 0$ and $v_{\max} \in (0, |M|)$ and minimize the base energy (as well as the main energy) over open sets in Q_R of volume at most v_{\max} . We use the notation

$$(2.5) \quad \mathcal{H} = \mathcal{H}_{R, v_{\max}} = \{\Omega \subseteq M \text{ open} : |\Omega| \leq v_{\max}, \Omega \subseteq Q_R\}$$

to denote this collection of sets. We will assume throughout the paper, without further note, that R is chosen to be large enough so that $|Q_R| \geq v_{\max}$. In this way, the container Q_R does not obstruct the hard volume constraint and the collection \mathcal{H} is nonempty. Moreover, we will always assume that $v_{\max} < |M|$ when the volume of M is finite.

Given parameters $0 < v < v_{\max} < |M|, \eta > 0, \mathfrak{T} \in (0, 1]$ and $R > 0$, we consider the following minimization problem for the base energy:

$$(2.6) \quad \mathcal{E}_{\min} = \mathcal{E}_{\min}(v, v_{\max}, \eta, \mathfrak{T}, R) = \inf\{\mathcal{E}(\Omega) : \Omega \subseteq \mathcal{H}\};$$

we will drop the parameters unless ambiguous. The collection of minimizers for this problem will be denoted by

$$(2.7) \quad \mathcal{M} = \mathcal{M}(v, v_{\max}, \eta, \mathfrak{T}, R) = \{\Omega \subseteq \mathcal{H} : \mathcal{E}(\Omega) = \mathcal{E}_{\min}\},$$

which we will show in Lemma 3.10 is always nonempty. While several of our intermediate results will hold for larger parameter ranges, the reader should typically think of first fixing the desired and hard volume constraints v, v_{\max} and the radius R , then fixing the volume penalization parameter η , and finally fixing the coefficient \mathfrak{T} in front of the torsional rigidity.

In the following example, we consider the minimization problem (2.6) on simply connected space forms; generally speaking these are important motivating examples to keep in mind throughout the paper. Since these spaces are compact after modding out by the isometry group, we may send the parameter R to infinity and so the following example has no R parameter.

Example 2.3. Let (M, g) be Euclidean space, hyperbolic space, or the round sphere. Fix parameters v, v_{\max} with $0 < v < v_{\max}$, choosing $v_{\max} < |M|$ in the case of the round sphere. For $\eta \leq \eta_0(v, v_{\max})$ and $\mathfrak{T} \leq \mathfrak{T}_0(v, v_{\max}, \eta)$, geodesic balls of volume v are the unique minimizers of the base energy \mathcal{E} among all open bounded subsets of M with volume at most v_{\max} .

For general Riemannian manifolds, we cannot hope to explicitly characterize minimizers as in the example above. However, the following proposition summarizes the basic facts that will be proven about minimizers of the base energy in this paper. The proof follows from combining Lemma 3.10 and Theorem 6.3 (local and global existence respectively), Lemma 3.13 (connectedness), and Lemma 9.5 (volume constraint satisfied).

Proposition 2.4 (Minimizers of the base energy). *Fix $R > 0$ and $v \in (0, |M|)$. Fix any $v_{\max} \in (v, |M|)$. There exists $\eta_0 = \eta_0(R, v, v_{\max})$ such that for $\eta \leq \eta_0$, there exists $\mathfrak{T}_0 = \mathfrak{T}_0(R, v, v_{\max}, \eta)$ such that the following holds. If $\mathfrak{T} \leq \mathfrak{T}_0$, then $\mathcal{E}_{\min} > -\infty$ and a minimizer of the base energy \mathcal{E} exists among sets in \mathcal{H} , i.e. the set \mathcal{M} is nonempty. Moreover, any minimizer in \mathcal{M} is connected and has volume equal to v .*

If M is compact or M/G is compact, then constants above may be taken independent of R and the minimization of \mathcal{E} may be taken among all open bounded subsets of M with volume at most v_{\max} .

2.2. Main energy functional. Let us now introduce the class of energy functionals whose existence and regularity theory constitute the heart of this paper and using which we can carry out a selection principle as described in the introduction (in Section 2.4 below). Following the discussions in the introduction and in the previous subsection, we are interested in functionals of the type (1.6), though as in the previous subsection we relax the volume constraint. So, the general form of the functionals we consider is

$$(2.8) \quad \mathcal{F}_\tau(\Omega) = \mathcal{E}(\Omega) + \tau \mathfrak{h}(\Omega).$$

Here $\mathfrak{h}(\Omega)$ is a functional mapping open bounded sets to \mathbb{R} and $\tau < \tau_0$ is a parameter that will be chosen sufficiently small depending on the parameters of the base energy \mathcal{E} . The existence and regularity results of this paper will hold whenever \mathfrak{h} is an *admissible nonlinearity* as defined in Definition 2.11 below. Before defining this general class of admissible nonlinearities, we motivate the definition with some important concrete examples that will be used for our selection principle applications. In these examples, the relevant nonlinearities take the form

$$(2.9) \quad \mathfrak{h}(\Omega) = \sqrt{c^2 + (c - d_*(\Omega)^2)^2} - c$$

for a constant $c \in (0, 1]$, where $d_*(\Omega)$ is a suitably defined distance of Ω to the class of minimizers \mathcal{M} of the base energy. Note that (1.6) takes this form with $d_*(\Omega) = d_1(\Omega, B_R(x_\Omega))$ on Euclidean space; the subtraction of the constant c in (2.9) is immaterial for the minimization problem but conveniently normalizes the functional. We are interested in generalizing this in two directions: (1) we want a distance d_* that measures the size of the symmetric difference as well as the eigenfunction difference, and (2) we wish to suitably generalize this to certain Riemannian manifolds.

In order to establish Step (1) of the selection principle (and ultimately quantitative stability) in the form of Corollary 1.2 in which both the eigenfunction distance d_1 and the asymmetry distance d_0 are controlled (say on Euclidean space), a first approach might be to let $d_*(\Omega) = d_1(\Omega, B_R(x_\Omega)) + d_0(\Omega, B_R(x_\Omega))$ in (2.9), where d_0 is the symmetric difference defined in (1.3). The issue here, already understood in [9], is that

the distance d_0 is not sufficiently smooth (as the integral of a characteristic function) and consequently minimizers of such a functional will not be smooth.

To deal with this, we follow the approach of [9] and consider a smoothed out version of the asymmetry term d_0 . More specifically, consider a bounded, open set U with smooth boundary (say ∂U is a disjoint union of C^2 embedded submanifolds); in the relevant applications U will be a minimizer of the base energy \mathcal{E} , so in Euclidean space, hyperbolic space, or the round sphere, take U to be a geodesic ball. Take f to be a smooth nondecreasing function with $f(t) = t$ for $|t| \leq c_0$ and $f(t)$ is locally constant for $|t| > 2c_0$, where c_0 is selected so that $d(x, \partial U)$ is a C^2 function on $\{d(x, \partial U) < 4c_0\}$. One can always find such a c_0 ; see [24, Lemma 14.16]). Now, let ψ_U be given by

$$\psi_U(x) = \begin{cases} f(d(x, \partial U)) & x \in U \\ -f(d(x, \partial U)) & x \in M \setminus U. \end{cases}$$

The essential point of the function ψ_U is that $\int_{\Omega} \psi_U - \int_U \psi_U$ is comparable to $|U \Delta \Omega|^2$ for Ω whose boundary is a (small) normal graph over ∂U , and controls $|U \Delta \Omega|^2$ in all cases. More precisely, for any Ω we have

$$\int_{\Omega} \psi_U - \int_U \psi_U = \int_{\Omega \Delta U} |\psi_U| \geq \int_{\Omega \Delta U} \min\{|d(x, \partial U)|, c_0\} dx \geq c|U \Delta U|^2,$$

where the last step can be seen from the fact that the integral is minimized by tubular neighborhoods of U , and then by changing variables and using the smoothness of ∂U . On the other hand, if the boundary of Ω is a C^2 normal graph ξ over ∂U with $|\xi| \leq c_0$, we may change variables and integrate

$$\int_{\Omega \Delta U} |\psi_U| \leq C \int_{\partial U} \int_0^{|\xi(x)|} t dt d\mathcal{H}^{n-1} = C \int_{\partial U} \int_0^{|\xi(x)|^2/2} 1 dt d\mathcal{H}^{n-1} \leq C|U \Delta \Omega|^2.$$

So, given any two bounded open sets Ω and U with U smooth as above, we let

$$(2.10) \quad d_*(\Omega, U)^2 = \int_{\Omega} \psi_U - \int_U \psi_U + \int |u_{\Omega} - u_U|^2$$

where u_{Ω} denotes the first eigenfunction of Ω , normalized so that $u \geq 0$ and $\|u_{\Omega}\|_{L^2} = 1$ and extended by zero to be defined on all of M . In each of the following examples, we will use $d_*(\Omega, U)$ to define a distance $d_*(\Omega)$ to the collection \mathcal{M} of minimizers of the base energy and consider the functional (2.8) with the nonlinearity (2.9) for this suitable definition of $d_*(\Omega)$.

Remark 2.5 (Unique first eigenfunction). Lemma 3.14 below guarantees that if $\mathcal{E}(\Omega) \leq \mathcal{E}_{min} + \tau_0$ for τ_0 is sufficiently small, then Ω has a unique first eigenfunction and so $d_*(\Omega, U)$ is well-defined for such a set. This means that by choosing the parameter τ in (2.8) accordingly, u_{Ω} and $d_*(\Omega, U)$ will be well-defined for any minimizing sequence or minimizer of (2.8). As such, we will generally only define \mathfrak{h} for sets with $\mathcal{E}(\Omega) \leq \mathcal{E}_{min} + \tau_0$. There is no loss of generality, at least when considering minimizers, to define $\mathfrak{h}(\Omega) = 1$ if $\mathcal{E}(\Omega) \geq \mathcal{E}_{min} + \tau_0$.

The simplest example is the case when only one minimizer of the base energy exists.

Example 2.6 (Unique minimizer). Suppose that the collection of minimizers \mathcal{M} of the base energy consists of a single set U , and U has boundary of class C^2 . Then the distance $d_*(\Omega)$ of Ω to the collection of minimizers of the base energy with $v = |\Omega|$ is defined by $d_*(\Omega) = d_*(\Omega, U)$.

In the case when minimizers of the base energy are not unique, to define a distance of a given set Ω to the collection of minimizers, we must select the nearest minimizer to Ω . On Euclidean space, where translation invariance gives rise to nonuniqueness of minimizers, we do this by choosing the ball with the same barycenter as Ω .

Example 2.7 (Euclidean space). Let (M, g) be Euclidean space. Given a bounded open set Ω , let R be the unique radius so that $|\Omega| = |B_R|$ and let $x_{\Omega} = \int_{\Omega} x dx$ be the barycenter of Ω . Then the distance $d_*(\Omega)$ of Ω to the collection of balls B_R (i.e. the minimizers of the base energy with $v = |\Omega|$) is defined by $d_*(\Omega) = d_*(\Omega, B_R(x_{\Omega}))$.

We saw in Example 2.3 that the situation on hyperbolic space and the round sphere is similar to Euclidean space: the collection of minimizers of the base energy comprises geodesic balls of suitable radius centered at *any* point. In these cases, we define an appropriate “set center” in analogy to the Euclidean barycenter.

Example 2.8 (Hyperbolic space). Let (M, g) be hyperbolic space. Given a set bounded open set $\Omega \subset M$, define the set center x_Ω of Ω by $x_\Omega = \operatorname{argmin}_x \int_\Omega d^2(x, y) dm(y)$, which is well-defined. Then the distance $d_*(\Omega)$ of Ω to the collection of balls B_R (i.e. the minimizers of the base energy with $v = |\Omega|$) is defined by $d_*(\Omega) = d_*(\Omega, B_R(x_\Omega))$.

On the round sphere, we cannot define the notion of set center in analogy to Example 2.8 as the minimum may not be uniquely achieved. We instead consider the following.

Example 2.9 (The round sphere). Let (M, g) be the round sphere. Consider its standard embedding in \mathbb{R}^{n+1} . Given any open $\Omega \subseteq M$, let $y_\Omega = \int_\Omega y d\mathcal{H}^n(y)$, where $y, y_\Omega \in \mathbb{R}^{n+1}$ and this is a (vector-valued) surface integral. If $y_\Omega \neq 0$, then define the set center of Ω as $x_\Omega = y_\Omega/|y_\Omega|$. The distance $d_*(\Omega)$ of Ω to the collection of balls B_R (i.e. the minimizers of the base energy with $v = |\Omega|$) is defined by $d_*(\Omega) = d_*(\Omega, B_R(x_\Omega))$.

In all three of the preceding examples, uniqueness of minimizers of the base energy fails solely due to the isometries of the space, and modulo the isometry group G , there is a unique minimizer U with boundary of class C^2 . In more general situations of this type, we define a suitable notion of *set center*, which is a mapping $\Omega \mapsto x_\Omega$ for which there is a unique minimizer U_Ω of the base energy with the same set center as Ω . Because the definition is somewhat technical, we postpone it until Definition A.2 in Appendix A. In this situation, we define $d_*(\Omega)$ analogously to the examples above; the Euclidean barycenter and the sets centers for hyperbolic space and the round sphere above are examples of set centers in the sense of Definition A.2.

Example 2.10 (Minimizer is unique modulo isometries). Assume that \mathcal{M} consists of a single set U up to the action of the isometry group G . Also assume that it has smooth boundary, and that there exists a set center $\Omega \mapsto x_\Omega$ adapted to U . We define

$$(2.11) \quad d_*(\Omega) = d_*(\Omega, U_\Omega)$$

where U_Ω the unique set in \mathcal{M} with the same set center as Ω .

It may be tempting to give generic constructions of $d_*(\Omega)$, independent of set centers, by solving minimization problems along the lines of $U_\Omega = \operatorname{argmin} \{d_*(\Omega, U) : U \in \mathcal{M}\}$ and letting $d_*(\Omega) = d_*(\Omega, U_\Omega)$. However, such variational problems have issues with uniqueness and regularity of the solution map: these are closely related to the convexity and smoothness of the “distance” functional being minimized, as well as the structure of the family \mathcal{M} .

All of the examples given above fall into the general class of admissible nonlinearities \mathfrak{h} for which our main existence and regularity results are valid. Below, u_Ω is used to denote the normalized, nonnegative first eigenfunction of Ω (see (2.2)); recalling Remark 2.5, this function exists and is unique when $\mathcal{E}(\Omega) \leq \mathcal{E}_{\min} + \tau_0$.

Definition 2.11 (Admissible nonlinearities). *Fix a (possibly trivial) closed subgroup $G_0 \leq G$ of the isometry group of (M, g) . A function \mathfrak{h} mapping bounded open subsets of M to \mathbb{R} is said to be an admissible nonlinearity with respect to G_0 if the following properties hold.*

- (N1) $\mathfrak{h} \in [0, 1]$.
- (N2) \mathfrak{h} is invariant under isometries in G_0 .
- (N3) If Ω, Ω' are bounded open sets with $\mathcal{E}(\Omega) \leq \mathcal{E}_{\min} + \tau_0$, then

$$|\mathfrak{h}(\Omega) - \mathfrak{h}(\Omega')| \leq |\Omega \Delta \Omega'| + \int |u_\Omega - u_{\Omega'}|.$$

- (N4) For any $x_0 \in M$ and $r \leq r_0$, take a one-parameter family ϕ_t of diffeomorphisms with $\phi_0(x) = x$ and $|\partial_t \phi_t| \leq 1$ such that $|\{x : \phi_t(x) \neq x\}| \subseteq B_r(x_0)$. Assume $\phi_t(\Omega)$ is a set with $\mathcal{E}(\phi_t(\Omega)) \leq \mathcal{E}_{\min} + \tau_0$ and $\mathcal{H}^{n-1}(\partial\Omega) < \infty$; then

$$\limsup_{t \rightarrow 0} \frac{1}{|t|} \left| \mathfrak{h}(\phi_t(\Omega)) - \mathfrak{h}(\Omega) - \int (u_{\phi_t(\Omega)} - u_\Omega) a_\Omega - \left[\int_{\phi_t(\Omega)} b_\Omega - \int_\Omega b_\Omega \right] \right| \leq C_{\mathfrak{h}} r^n,$$

where a_Ω, b_Ω are continuous functions $M \rightarrow \mathbb{R}$ with $|a_\Omega|, |b_\Omega| \leq 1$. Moreover, $\|b_\Omega\|_{C^{1,1}(Q_R)} \leq 1$ and $\|a_\Omega\|_{C^{0,1}(\Omega)} \leq 1$.

Assumption (N1) simply says that \mathfrak{h} is nonnegative and uniformly bounded on bounded open sets, and it is convenient to normalize any such functional so that the upper bound is 1. Assumption (N2) should be understood with Example 2.10 in mind; when the base energy \mathcal{E} has a nonunique minimizer due only to a subgroup of isometries G_0 , then (N2) is used to compare a set Ω to the nearest minimizer of the base energy. Assumption (N3) heuristically states that \mathfrak{h} is a Lipschitz function $\Omega \mapsto \mathbb{R}$, where the topology placed on the space of sets is either the L^1 distance $|\Omega \Delta \Omega'|$ or a distance governed by the L^1 difference of their eigenfunctions. Assumption (N4) should be viewed as a higher regularity assumption on \mathfrak{h} ; it is asking that \mathfrak{h} is not only Lipschitz but also C^1 , with derivative represented by the functions a_Ω and b_Ω .

Only the first three assumptions (N1), (N2), and (N3) are needed for the existence theory and core estimates up through Section 7. The last assumption (N4) is used to derive the Euler-Lagrange equation in Section 8 (indeed, the functions a_Ω and b_Ω appear in the first variation) and for the regularity of minimizers in Sections 9 and 10. In particular the extra regularity of a_Ω, b_Ω is only needed in Section 10 to obtain better smoothness of $\partial\Omega$ for minimizers.

In Appendix A, we verify that the examples given above are admissible nonlinearities. Given an admissible nonlinearity, it is easy to construct others via composition with functions of one variable:

Remark 2.12. Given any admissible nonlinearity \mathfrak{h} , and any C^1 function $\phi : [0, 1] \rightarrow [0, 1]$ with $|\phi'| \leq 1$, the composition $\phi \circ \mathfrak{h}$ is again an admissible nonlinearity. No monotonicity or structure of ϕ is required here.

2.3. Minimizers of the main energy. As we did for the base energy, we will fix $R > 0$ and v_{\max} and minimize the main energy over the set $\mathcal{H} = \mathcal{H}_{R, v_{\max}}$ defined in (2.5). We say that a bounded open set Ω is a *minimizer* if $\Omega \in \mathcal{H}$ and for any $\Omega' \in \mathcal{H}$,

$$(2.12) \quad \mathcal{F}_\tau(\Omega) \leq \mathcal{F}_\tau(\Omega').$$

We say that Ω is an *inward minimizer* if (2.12) holds for any $\Omega' \in \mathcal{H}$ with $\Omega' \subseteq \Omega$. We say that Ω is an *outward minimizer* if (2.12) holds for any $\Omega' \in \mathcal{H}$ with $\Omega' \supset \Omega$.

The following theorem summarizes the main existence and regularity properties for \mathcal{F}_τ that are established in this paper. The proof follows by combining Theorem 6.2 and Theorem 6.3 (existence), Proposition 9.5 (volume constraint), Theorem 8.1 (free boundary condition), Theorem 9.4 ($C^{1,\alpha}$ regularity), and Corollary 10.5 (higher regularity).

We say that a collection \mathcal{C} of bounded open subsets of M is *uniformly C^k* if there is a constant $C > 2/\text{inj}_M$ such that for every $U \in \mathcal{C}$ and every $x \in \partial U$, the set $\partial U \cap B_{1/C}(x)$ may be expressed as a graph with C^k norm at most C over a hyperplane in normal coordinates.

Theorem 2.13. *Fix $R > 0$ and $v \in (0, |M|)$. Fix any $v_{\max} \in (v, |M|)$. There exist positive constants $\eta, \mathfrak{T}, \tau_0 > 0$ depending on R, v, v_{\max} such that the following holds. Let \mathfrak{h} be an admissible nonlinearity with respect to the isometry group G of (M, g) and for fixed $\tau < \tau_0$, consider the energy functional \mathcal{F}_τ defined in (2.8) with the parameters v, η , and \mathfrak{T} in the base energy \mathcal{E} of (2.8).*

- (1) *There exists an open set Ω that minimizes \mathcal{F}_τ among sets in $\mathcal{H} = \mathcal{H}_{R, v_{\max}}$. Any such minimizer satisfies the volume constraint $|\Omega| = v$, is a set of finite perimeter, and has a unique first eigenfunction u_Ω up to scaling.*
- (2) *There exist $C > 0$ and $\alpha \in (0, 1)$ depending only on v, v_{\max}, R , a function $\rho \geq -C\tau$ with $\|\rho\|_{C^{0,\alpha}(\partial\Omega)} \leq C$, and a constant $A_0 \in [1/C, C]$ such that the normalized first eigenfunction u_Ω satisfies the free boundary condition*

$$|\nabla u_\Omega(x)|^2(1 + \rho(x)) = A_0$$

for \mathcal{H}^{n-1} -a.e. $x \in \partial^*\Omega \cap Q_R$.

- (3) *Assume further that the collection \mathcal{M} of minimizers of the base energy is uniformly C^4 . For every $\alpha \in (0, 1)$ and $r_1 > 0$ there is a $\tau_1 = \tau_1(v, v_{\max}, \eta, r_1, \alpha) > 0$ such that the following holds. If $\tau < \tau_1$, then for any minimizer Ω of \mathcal{F}_τ , the set $\partial\Omega \cap \{x \in Q_R : d(x, \partial Q_R) > r_1\}$ may be parametrized as a $C^{2,\alpha}$ normal graph (with $C^{2,\alpha}$ norm bounded by r_1) over ∂U for some $U \in \mathcal{M}$.*

If M/G is compact, then all constants may be taken to be independent of R , the minimization of \mathcal{F}_τ may be taken among all bounded open subsets of M with volume at most v_{\max} , and in (3) the entire boundary of any minimizer Ω may be written as a $C^{2,\alpha}$ graph over ∂U for some $U \in \mathcal{M}$.

Theorem 2.13(1) implies that the minimization of \mathcal{F}_τ among sets in \mathcal{H} is equivalent to the minimization problem

$$\inf\{\mathcal{F}_\tau(\Omega) : \Omega \in \mathcal{H}^*, |\Omega| = v\}$$

where we let \mathcal{H}^* denote either the collection of open bounded subsets of Q_R or, in the case that M/G is compact, the collection of open bounded subsets of M .

In order to better track the dependence of constants in various estimates in the sections to follow, we define two quantities which measure growth properties of the eigenfunction and torsion function of Ω . Let Ω be an open subset of Q_R with $\lambda_1(\Omega) < \lambda_2(\Omega)$, and u_Ω the first eigenfunction (normalized so that $\int u_\Omega^2 = 1$ and $u_\Omega \geq 0$ as usual), while w_Ω is the torsion function. Then

$$(2.13) \quad \text{UP}(\Omega) = \sup \left\{ \frac{u_\Omega(x) + \sqrt{\mathfrak{T}} w_\Omega(x)}{d(x, \partial\Omega)} : x \in \Omega, d(x, \partial\Omega) \in (0, 1) \right\}$$

and

$$(2.14) \quad \text{DO}(\Omega) = \inf \left\{ \frac{1}{r} \sup_{x \in B_r(y)} u_\Omega(x) + \sqrt{\mathfrak{T}} w_\Omega(x) : y \in \bar{\Omega}, r \in (0, 1) \right\}.$$

In general, $\text{UP}(\Omega) \in (0, \infty]$ while $\text{DO}(\Omega) \in [0, \infty]$, though we will show in Section 3.1 that $\text{DO}(\Omega) < \infty$ by showing that u_Ω and w_Ω are bounded. We also always have $\text{DO}(\Omega) \leq \text{UP}(\Omega)$. The core nonlinear estimates for this type of variational problem, following [5], are that outward minimizers have $\text{UP}(\Omega) \leq C < \infty$, while inward minimizers have $\text{DO}(\Omega) \geq c > 0$.

2.4. Selection Principle. In this section, we establish the main linear stability-implies-nonlinear stability result of this paper, which is Theorem 2.14 below. This is a generalization of Corollary 1.2. The proof is essentially self-contained assuming Proposition 2.4 and Theorem 2.13, with the exception of some continuity statements of Lemma 3.10 and Lemma 3.16. We operate under one of the following set of assumptions:

Case 1: M/G is compact.

Case 2: M is arbitrary and $R > 0$ is fixed.

Let $\mathcal{H}^* = \{\Omega \text{ open, bounded}\}$ in Case 1 and $\mathcal{H}^* = \{\Omega \text{ open} : \Omega \subseteq Q_R\}$ in Case 2. Fix $0 < v < v_{\max} < |M|$, with $v_{\max} < |Q_R|$ in Case 2. Let $\eta = \eta(v, v_{\max}, R)$ be chosen small enough according to Proposition 2.4 as well as Lemma 3.16 in Case 1 (Lemma 3.10 in Case 2). Then choose $\mathfrak{T}_* = \mathfrak{T}_*(v, v_{\max}, R, \eta) > 0$ small enough according to Proposition 2.4. Fix any $\mathfrak{T} < \mathfrak{T}_*$. All constants and parameters below depend on these fixed values.

We are interested in looking at the quantitative stability of minimizers of the base energy \mathcal{E} with a volume constraint:

$$(2.15) \quad \inf\{\lambda_1(\Omega) + \mathfrak{T} \text{tor}(\Omega) : \Omega \in \mathcal{H}^*, |\Omega| = v\},$$

By Proposition 2.4, the infimum in (2.15) is equal to the infimum $\mathcal{E}_{\min}(v)$ of the unconstrained problem (2.6), and the two minimization problems are equivalent in the sense that if a set $\Omega \in \mathcal{H}^*$ attains the infimum in (2.6), then it has $|\Omega| = v$ and it attains the infimum in (2.15), and conversely. Let \mathcal{M}^* denote the collection of minimizers of this variational problem, which is nonempty by Proposition 2.4. In Theorem 2.14, we make the following assumptions about \mathcal{M}^* :

In Case 1, assume that \mathcal{M}^* may be represented as $\{e(U)\}_{e \in G}$ for a set $U \in \mathcal{M}^*$, the boundary ∂U is smooth, and there exists a set center adapted to U defined for sets in $\{\Omega : \mathcal{E}(\Omega) \leq \mathcal{E}_{\min}(v) + \delta\}$ for δ small enough. Recall that set centers are defined in Definition A.2 and that examples to keep in mind are those on simply connected. space forms in Examples 2.7–2.9.

In Case 2, assume that \mathcal{M}^* contains only one element U and $U \subset\subset Q_R$ and that ∂U is smooth.

More general situations, such as finitely many minimizers (modulo isometries in Case 1) may be considered with suitable modifications to the arguments here. Given an open bounded set Ω , let $d_*(\Omega)$ be the distance to the collection of minimizers as defined in Example 2.10. In Case 1, we let U_Ω denote the unique minimizer with the same set center as Ω .

Theorem 2.14. *For M and \mathcal{M}^* as above and $\varepsilon > 0$ small, assume that the inequality*

$$(2.16) \quad \lambda_1(\Omega) + \mathfrak{T}\text{or}(\Omega) - \mathcal{E}_{\min} \geq cd_*(\Omega)^2$$

holds for all $\Omega \in \mathcal{H}^$ with $|\Omega| = v$ and $\partial\Omega$ expressible as a C^2 normal graph over ∂U_Ω with $C^{2,\alpha}$ norm bounded by ε . Then, up to possibly replacing c with a smaller constant, (2.16) holds for all $\Omega \in \mathcal{H}^*$ with $|\Omega| = v$.*

Proof. Suppose by way of contradiction that the theorem is false. We may find a sequence of sets $\{\Omega_j\}$ with $\Omega_j \in \mathcal{H}^*$ and $|\Omega_j| = v$ such that

$$(2.17) \quad \mathcal{E}(\Omega_j) = \lambda_1(\Omega_j) + \mathfrak{T}\text{or}(\Omega_j) \leq \mathcal{E}_{\min} + \frac{d_j^2}{j},$$

where here and in the remainder of the proof, we let $d_j = d_*(\Omega_j)$. The quantity $d_*(\Omega)$ is bounded above uniformly for any $\Omega \in \mathcal{H}^*$, so (2.17) implies that

$$\lim_{j \rightarrow \infty} \mathcal{E}(\Omega_j) = \mathcal{E}_{\min}(v).$$

Applying Lemma 3.16 in Case 1 or Lemma 3.10 in Case 2, we see that $d_j \rightarrow 0$, and so there exists $U \in \mathcal{M}^*$ such that, after passing to a subsequence, $|\Omega_j \Delta e_j(U)| \rightarrow 0$ as $j \rightarrow \infty$ for some $e_j \in G$.

For each $j \in \mathbb{N}$, consider the minimization problem

$$(2.18) \quad \inf\{\mathcal{F}_\tau(\Omega) = \mathcal{E}(\Omega) + \tau\mathfrak{h}(\Omega) : \Omega \in \mathcal{H}^*\},$$

where we let

$$\mathfrak{h}(\Omega) = \sqrt{d_j^4 + (d_*(\Omega)^2 - d_j^2)^2} - d_j^2$$

where $\tau < \tau_0$ is a small parameter, and $\tau_0 = \tau_0(v, v_{\max}, R, \eta, \mathfrak{T})$ is chosen so that Theorem 2.13 holds. By Theorem 2.13, a minimizer V_j of this variational problem exists and satisfies $|V_j| = v$. Let us take Ω_j as a competitor in (2.18) in order to relate the deficits and distances of Ω_j and V_j . We find

$$(2.19) \quad \mathcal{E}(V_j) + \tau \left(\sqrt{d_j^4 + (d_*(V_j)^2 - d_j^2)^2} - d_j^2 \right) \leq \mathcal{E}(\Omega_j)$$

$$(2.20) \quad \leq \mathcal{E}(U) + \frac{d_j^2}{j}.$$

In (2.20) we have applied (2.17). Subtracting $\mathcal{E}(U) = \mathcal{E}_{\min}(v)$ throughout (2.19) and (2.20) (and recalling that $\mathcal{E}_{\min}(v) \leq \mathcal{E}(V_j)$ by definition), we directly see that

$$(2.21) \quad \mathcal{E}(V_j) - \mathcal{E}_{\min}(v) \leq \mathcal{E}(\Omega_j) - \mathcal{E}_{\min} \rightarrow 0.$$

Furthermore, (2.20) tells us that

$$\sqrt{d_j^4 + (d_*(V_j)^2 - d_j^2)^2} \leq \left(1 + \frac{1}{\mathfrak{T}j}\right) d_j^2.$$

After squaring both sides of this inequality, we find that $(d_*(V_j)^2 - d_j^2)^2 \leq 3d_j^4/\tau j$. This guarantees, first of all, that $V_j \notin \mathcal{M}^*$ for j sufficiently large. Moreover, taking the square root of both sides, we find that for sufficiently large j ,

$$(2.22) \quad \frac{1}{2}d_j \leq d_*(V_j).$$

So, combining the contradiction hypothesis (2.17) with the deficit and distance comparisons (2.21) and (2.22), we see that

$$(2.23) \quad \mathcal{E}(V_j) - \mathcal{E}_{\min}(v) \leq \mathcal{E}(\Omega_j) - \mathcal{E}_{\min}(v) \leq \frac{d_j^2}{j} \leq \frac{2d_*(V_j)^2}{j}.$$

By Theorem 2.13, V_j may be represented as a $C^{2,\alpha}$ normal graph over ∂U_j for some $U_j \in \mathcal{M}$, with $C^{2,\alpha}$ norm controlled by $o_j(1)$. Moreover, the set centers of U_j and V_j converge, since $d(x_{U_j}, x_{V_j}) \leq C|U_j \Delta V_j| = o_j(1)$, which immediately implies that $|U_j \Delta U_{V_j}| = o_j(1)$ from the properties of set centers. Here U_{V_j} is the unique minimizer with the same set center as V_j . Using the smoothness of the minimizers U and the tubular neighborhood theorem, this means for j large ∂U_j (and then also ∂V_j) may be written as a $C^{2,\alpha}$ normal graph over ∂U_{V_j} with small $C^{2,\alpha}$ norm. Applying the hypothesis (2.16) to V_j , we see that

$$cd_*(V_j)^2 \leq \mathcal{E}(V_j) - \mathcal{E}_{min},$$

which for j large is a contradiction. This completes the proof. \square

Recall from the discussion in Section 2.2 that $d_*(\Omega)^2$ is comparable to

$$\bar{d}_*(\Omega)^2 := |\Omega \Delta U_\Omega|^2 + \int |u_\Omega - u_{U_\Omega}|^2$$

for Ω whose boundary is a (small) normal graph over ∂U , and controls $\bar{d}_*(\Omega)^2$ in all cases. This leads to the following corollary, which in particular proves Corollary 1.2.

Corollary 2.15. *Theorem 2.14 remains valid with \bar{d}_* in place of d_* .*

3. EIGENFUNCTIONS, TORSION FUNCTIONS, AND THE BASE ENERGY

This section collects some useful facts about eigenfunctions and torsion functions of bounded open subsets of M , as well some initial properties about the base energy \mathcal{E} and its optimizers. Section 3.1 is devoted to general properties of eigenfunctions and torsion functions.

Section 3.2 contains Proposition 3.9, the key estimate that will allow us to control the nonlinear perturbation in the main energy in the remainder of the paper. This lemma estimates the difference of eigenfunctions on nested domains in terms of their eigenvalue and torsional rigidity differences. Fundamentally, this lemma highlights why we are able to establish existence and regularity of minimizers of the main energy when the coefficient \mathfrak{T} in front of the torsional rigidity is positive (and suitably small), but not when $\mathfrak{T} = 0$ and the leading term in the base energy is only the first eigenfunction.

In Section 3.3, we establish the existence and initial properties of minimizers of the base energy in Q_R , while Section 3.4 extends these some of these properties of minimizers of the base energy globally.

We recall the set $\mathcal{H} = \mathcal{H}_{R, v_{\max}}$ defined in (2.5). Recall that we always implicitly assume that v_{\max} and R are chosen so that $v_{\max} \leq |Q_R|$ and, when $|M|$ is finite, so that $v_{\max} < |M|$.

3.1. Eigenfunctions and torsion functions. Our first lemma collects some basic facts about eigenfunctions and torsion functions. The proof is elementary and so we omit it.

Lemma 3.1. *Let $\Omega \subset M$ be a bounded open set. Then the following properties hold.*

- (1) w_Ω is nonnegative and satisfies $\Delta w_\Omega \geq -1_\Omega$ on M in the sense of distributions.
- (2) $\Delta w_\Omega = -1$ on Ω .
- (3) $\int_\Omega |\nabla w_\Omega|^2 = \int_\Omega w_\Omega = -2\text{tor}(\Omega)$.
- (4) If $\lambda_1(\Omega) < \lambda_2(\Omega)$, the normalized eigenfunction u_Ω , up to replacing with $-u_\Omega$, is nonnegative and $\Delta u_\Omega \geq -\lambda_1(\Omega)u_\Omega$ on M in the sense of distributions.
- (5) If u_Ω is as in (4), $\Delta u_\Omega = -\lambda_1(\Omega)u_\Omega$ on Ω .

Next, we show that eigenfunctions of a domain $\Omega \subset M$ are bounded uniformly in terms of an eigenvalue upper bound.

Lemma 3.2 (Eigenfunction upper bound). *For each $\lambda_0 > 0$, there is a constant $C_1 = C_1(\lambda_0)$ such that the following holds. Let Ω be an open bounded set and fix $\lambda \leq \lambda_0$. For any function $u \in H_0^1(\Omega)$ with $\|u\|_{L^2} = 1$ satisfying $\Delta|u| \geq -\lambda|u|$ on Ω ,*

$$|u| \leq C_1.$$

In particular, this applies to any eigenfunction of Ω .

Proof. Let $K(x, y, t)$ be the heat kernel on M (see [19]), and set $C_1 = e^{\lambda_0} \sup_{x \in M} \sqrt{K(x, x, 2)}$. Since (M, g) has bounded geometry, C_1 is finite. Extending u by zero to be defined on all of M , the function $e^{-t\lambda}|u(x)|$ is a subsolution to the heat equation on M . Using the comparison principle,

$$e^{-t\lambda}|u(x)| \leq \int K(x, y, t)|u(y)|dy \leq \left(\int K^2(x, y, t)dy \right)^{1/2} \|u\|_{L^2} = \sqrt{K(x, x, 2t)}.$$

Set $t = 1$ to conclude. \square

Next, the torsion function of an open bounded set is bounded in terms of its torsional rigidity.

Lemma 3.3 (Torsion function upper bound). *For each $T > 0$, there exists a constant $C_2 = C_2(T)$ such that the following holds. Let Ω be an open bounded set with $\text{tor}(\Omega) \geq -T$.*

$$w_\Omega \leq C_2.$$

Proof. In local coordinates on a ball $B_r(x)$ around any point $x \in M$, we have that $w_\Omega \geq 0$ and $\Delta w_\Omega \geq -1$, so applying the local maximum principle ([24], Theorem 8.17) gives

$$\sup_{B_{r/2}(x)} w_\Omega^2 \leq C(r, x) \left[\int_{B_r(x)} |\nabla w_\Omega|^2 + 1 \right] \leq C(r, x)[- \text{tor}(\Omega) + 1].$$

In the second inequality we used Lemma 3.1(3). Since (M, g) has bounded geometry, the values of r and the constant $C(r, x) = C$ may be taken uniformly for all $x \in M$. We conclude by letting $C_2 = C[T + 1]$. \square

The next lemma shows that the eigenfunction is controlled by the torsion function.

Lemma 3.4 (Torsion function controls eigenfunction). *Let Ω be an open bounded set with $|\Omega| < |M|$ such that $\lambda_1(\Omega) < \lambda_2(\Omega)$. Then there is a constant C_3 depending only on C_1 of Lemma 3.2 and $\lambda_1(\Omega)$ such that*

$$u_\Omega \leq C_3 w_\Omega.$$

Proof. This follows directly by the maximum principle and Lemma 3.2: $-\Delta u_\Omega \leq \lambda_1(\Omega)C_1 = -\Delta \lambda_1(\Omega)C_1 w_\Omega$ on Ω , and both are in $H_0^1(\Omega)$, giving $u_\Omega \leq C_1 \lambda_1(\Omega) w_\Omega$ on Ω . \square

As a consequence of this lemma, if we write $\Delta u_\Omega = -\lambda_1(\Omega)u_\Omega + \mu$ in the sense of distributions, with μ a nonnegative measure, and similarly $\Delta w_\Omega = -1_\Omega + \nu$ where ν is a nonnegative measure, then $\mu \leq C\nu$. The opposite inequality $w_\Omega \leq C u_\Omega$ is much more subtle, and will be discussed for *minimizing* Ω in Proposition 7.8.

We now establish two forms of a Poincaré inequality on the sets Q_R defined at the beginning of Section 2; the key point here is that they depend only on the parameters R and v_{\max} . We let $\dot{H}_0^1(Q_R)$ and $\dot{H}^1(Q_R)$ denote the completions of $C_0^\infty(Q_R)$ and $C^\infty(Q_R)$ respectively with respect to the seminorm $\|\nabla u\|_{L^2(Q_R)}$.

Lemma 3.5 (Poincaré Inequality). *For every $v_{\max} < |M|$ and any $R > 0$, the embedding*

$$X = \{u \in \dot{H}_0^1(Q_R) : |\{u > 0\}| \leq v_{\max}\} \rightarrow L^p(Q_R)$$

is continuous for $p \in [1, \frac{2n}{n-2}]$ and compact for $p < \frac{2n}{n-2}$. If $v_{\max} < |Q_R|$, then the same is true as well for $X = \{u \in \dot{H}^1(Q_R) : |\{u > 0\}| \leq v_{\max}\}$.

Proof. It is well known (see for instance [25, Theorem 10.2] or [6, Theorems 2.30 and 2.34]) that on Q_R (or more generally, any compact Riemannian manifold with boundary), $H^1(Q_R)$ embeds continuously into $L^p(Q_R)$ for $p \in [1, \frac{2n}{n-2}]$ and the embedding is compact for $p < \frac{2n}{n-2}$.

Suppose by way of contradiction that the second embedding stated in the lemma fails for $p = 2n/(n-2)$. Take a sequence of functions $\{u_k\}$ with $\|u_k\|_{L^p(Q_R)} = 1$ and $|\{u_k > 0\}| \leq v_{\max}$ such that $\|\nabla u_k\|_{L^2(Q_R)} \rightarrow 0$. We see from Hölder's inequality that $\|u_k\|_{L^2(Q_R)} \leq C(Q_R)$. So, from the continuous embedding $H^1(Q_R) \hookrightarrow L^p(Q_R)$ we find that $u_k \rightarrow c$ in $H^1(Q_R)$, $L^p(Q_R)$ and pointwise a.e. in Q_R for some constant $c \in \mathbb{R}$. On one hand, $\|c\|_{L^p(Q_R)} = 1$, while on the other hand, Fatou's lemma implies that $|\{c > 0\}| \leq v_{\max}$ and so $c = 0$. We reach a contradiction.

Now, for $p < 2n/(n-2)$, the continuous embedding then follows by Hölder's inequality, while the compactness of the embedding follows from the compact embedding of $H^1(Q_R)$ into $L^p(Q_R)$ and the continuous embedding of $\{u \in \dot{H}^1(Q_R) : |\{u > 0\}| \leq v_{\max}\}$ into $L^2(Q_R)$ that we have just established.

We prove the first embedding of the lemma analogously. If M is compact, then the argument above carries over directly. If M is noncompact, we again argue by way of contradiction and take a sequence of functions with $\|u_k\|_{L^p(Q_R)} = 1$ and $\|\nabla u_k\|_{L^2(Q_R)} \rightarrow 0$. Extend each u_k by zero to be defined on Q_{2R} . By Hölder's inequality, $\|u_k\|_{L^2(Q_{2R})} \leq C$ and so from the continuous embedding $H^1(Q_{2R}) \hookrightarrow L^p(Q_{2R})$ we see that $u_k \rightarrow c$ in $H^1(Q_{2R})$, $L^p(Q_{2R})$ and pointwise a.e. in Q_{2R} . On one hand, $\|c\|_{L^p(Q_{2R})} = 1$, while on the other hand, $u_k \equiv 0$ on $Q_{2R} \setminus Q_R$ and so we see that $c = 0$. We reach a contradiction. Again, the continuous embedding for $p < 2n/(n-2)$ then follows by Hölder's inequality, while the compactness of the embedding follows from the compact embedding of $H_0^1(Q_R)$ embeds continuously into $L^p(Q_R)$ and the continuous embedding of $\dot{H}_0^1(Q_R)$ into $L^2(Q_R)$ that we have just established. \square

As an immediate consequence of Lemma 3.5, we may bound the torsional rigidity $\text{tor}(\Omega)$ from below uniformly for all domains in $\mathcal{H} = \mathcal{H}_{R, v_{\max}}$, the class of admissible domains defined in (2.5). Naturally, if M has bounded diameter, then the lower bound is independent of R .

Corollary 3.6 (Torsional rigidity lower bound). *Fix $R > 0$ and $v_{\max} < |M|$. There is a constant $C(R, v_{\max}) < \infty$ such that for any $\Omega \in \mathcal{H}$,*

$$\text{tor}(\Omega) \geq -C(R, v_{\max}).$$

Proof. Fix $\Omega \in \mathcal{H}$ and let w_Ω be the corresponding torsion function, so that $\text{tor}(\Omega) = \int_\Omega \frac{1}{2} |\nabla w_\Omega|^2 - w_\Omega$. Then for any $\varepsilon > 0$,

$$\int |w_\Omega| \leq C(R, v_{\max}) \|\nabla w_\Omega\|_{L^2} \leq \varepsilon \int |\nabla w_\Omega|^2 + C(\varepsilon, R, v_{\max}),$$

where the first inequality used Lemma 3.5 on w_Ω extended by 0 and the fact that $|\{w_\Omega > 0\}| \leq |\Omega| \leq v_{\max}$. Take $\varepsilon = \frac{1}{4}$ and reabsorb to get

$$\text{tor}(\Omega) = \int_\Omega \frac{1}{2} |\nabla w_\Omega|^2 - w_\Omega \geq -C(R, v_{\max}) + \int_\Omega \frac{1}{4} |\nabla w_\Omega|^2.$$

This concludes the proof. \square

Remark 3.7. Notice that from Corollary 3.6, for any $\Omega \in \mathcal{H} = \mathcal{H}_{R, v_{\max}}$, the constant C_2 in Lemma 3.3 depends only on R and v_{\max} .

The final lemma of this section is an elementary fact about the energy of the difference between any normalized nonnegative function and the first eigenfunction.

Lemma 3.8. *Let Ω be an open bounded set with $\lambda_1(\Omega) + \alpha \leq \lambda_2(\Omega)$ for some $\alpha > 0$, and $v \in H_0^1(\Omega)$ with $\int v^2 = 1$ and $v \geq 0$. Then*

$$\int |\nabla(u_\Omega - v)|^2 \leq \left(1 + \frac{2\lambda_1(\Omega)}{\alpha}\right) \left[\int |\nabla v|^2 - \lambda_1(\Omega)\right].$$

Proof. Let us use the shorthand $\lambda_1 = \lambda_1(\Omega)$ and $\lambda_2 = \lambda_2(\Omega)$, and write $v = \beta u_\Omega + v_\perp$ where $\beta = \int u_\Omega v$. Because $\int u_\Omega v_\perp = 0$ by definition, we observe that

$$\lambda_2(1 - \beta^2) = \lambda_2 \int v_\perp^2 \leq \int |\nabla v_\perp|^2 = \int |\nabla v|^2 - \lambda_1 \beta^2 = \int |\nabla v|^2 - \lambda_1 + \lambda_1(1 - \beta^2).$$

In particular, $\alpha(1 - \beta)^2 \leq \int |\nabla v|^2 - \lambda_1$ and $\int |\nabla v_\perp|^2 \leq (1 + \lambda_1/\alpha)(\int |\nabla v|^2 - \lambda_1)$. So, since $u_\Omega - v = (1 - \beta)u_\Omega + v_\perp$ and $0 \leq \beta \leq 1$, we see that

$$\int |\nabla(u_\Omega - v)|^2 = (1 - \beta)^2 \lambda_1 + \int |\nabla v_\perp|^2 \leq \left(1 + \frac{2\lambda_1}{\alpha}\right) \left(\int |\nabla v|^2 - \lambda_1\right).$$

This concludes the proof. \square

3.2. The key estimate on the nonlinear perturbation. In order to study the existence and regularity of minimizers of the main energy \mathcal{F}_τ , we will require some sharp estimates on how an admissible nonlinearity \mathfrak{h} changes under a change in domain from Ω to Ω' . From assumption (N3) in Definition 2.11, this is controlled by the sum of the symmetric difference $|\Omega \Delta \Omega'|$, which is easy to estimate, and the L^1 difference of eigenfunctions $\int |u_\Omega - u_{\Omega'}|$, which is not significantly more difficult to estimate. In fact, we do not know how to control this quantity in terms of the eigenvalue difference $|\lambda_1(\Omega') - \lambda_1(\Omega)|$ alone, at least in the case of outward perturbations. The following lemma does bound $\int |u_\Omega - u_{\Omega'}|$, but the difference $|\text{tor}(\Omega) - \text{tor}(\Omega')|$ appears on the right-hand side of the estimate as well. This is the fundamental reason for introducing the torsional rigidity term into the main energy functional \mathcal{F}_τ .

This lemma will be the crucial ingredient in proving the main estimates of Sections 4 and 5 and establishing our main existence result in Section 6. At least for the estimates, the specific form of the lemma, with *linear* dependence on $|\text{tor}(\Omega) - \text{tor}(\Omega')|$ and $|\lambda_1(\Omega') - \lambda_1(\Omega)|$, is essential; the (much) simpler sublinear versions of this inequality are insufficient.

Proposition 3.9. *Fix $0 < v_{\max} < |M|$, $\alpha > 0$ and $\lambda_0 > 0$. Let $\Omega \subseteq \Omega'$ be bounded open sets with $|\Omega'| \leq v_{\max}$, $\lambda_2(\Omega') \geq \lambda_1(\Omega') + \alpha$, and $\lambda_1(\Omega) \leq \lambda_0$. Let $f : M \rightarrow \mathbb{R}$ be a bounded function. Then there exists a constant $C_4 = C_4(\alpha, v_{\max}, \lambda_0)$ such that*

$$(3.1) \quad \left| \int f(u_{\Omega'} - u_\Omega) \right| \leq C_4 \|f\|_{L^\infty} (\text{tor}(\Omega) - \text{tor}(\Omega') + \lambda_1(\Omega) - \lambda_1(\Omega')).$$

In the case that the first eigenvalue of Ω is not simple, then Proposition 3.9 holds trivially for any first eigenfunction u_Ω . To see this, note that $\lambda_1(\Omega) \geq \lambda_1(\Omega')$ and $\lambda_2(\Omega) \geq \lambda_2(\Omega')$ from $\Omega \subseteq \Omega'$. So, if $\lambda_1(\Omega)$ is not simple, then $\lambda_1(\Omega) = \lambda_2(\Omega)$ and so $\lambda_1(\Omega) - \lambda_1(\Omega') \geq \alpha$, meaning the estimate here holds automatically for *every* normalized first eigenfunction u_Ω (using Lemma 3.2). So, in the proof below, we will assume that $\lambda_1(\Omega) < \lambda_2(\Omega)$, in which case u_Ω is uniquely defined.

Proof. We first consider the case of $\partial\Omega \cap \Omega'$ smooth, but all constants will be independent of the regularity of $\partial\Omega$. Let $\beta = \int f u_{\Omega'}$ and write $f = \beta u_{\Omega'} + f_\perp$. We estimate the contributions of $\beta u_{\Omega'}$ and f_\perp to the left-hand side of (3.1) separately. To estimate the contribution from the first term, use the normalization of the eigenfunctions $u_\Omega, u_{\Omega'}$:

$$\int \beta u_{\Omega'}(u_{\Omega'} - u_\Omega) = \beta \int u_{\Omega'}^2 - u_{\Omega'} u_\Omega = \frac{\beta}{2} \int u_{\Omega'}^2 + u_\Omega^2 - 2u_{\Omega'} u_\Omega = \frac{\beta}{2} \int |u_\Omega - u_{\Omega'}|^2.$$

Applying Lemma 3.8 and assuming without loss of generality that $\alpha \leq \lambda_1(\Omega')$, we have

$$\int |u_\Omega - u_{\Omega'}|^2 \leq \frac{1}{\lambda_1(\Omega')} \int |\nabla(u_\Omega - u_{\Omega'})|^2 \leq \frac{3}{\alpha} \left[\int |\nabla u_\Omega|^2 - \lambda_1(\Omega') \right] = \frac{3}{\alpha} [\lambda_1(\Omega) - \lambda_1(\Omega')]$$

and therefore

$$(3.2) \quad \left| \int \beta u_{\Omega'}(u_{\Omega'} - u_\Omega) \right| \leq \frac{3}{\alpha} \|f\|_{L^2} [\lambda_1(\Omega) - \lambda_1(\Omega')].$$

Now, the remainder f_\perp has $\|f_\perp\|_{L^2} \leq \|f\|_{L^2} \leq \sqrt{v_{\max}} \|f\|_{L^\infty}$ and $\int f_\perp u_{\Omega'} = 0$. This means we may solve for the potential

$$(3.3) \quad \begin{cases} -\Delta q - \lambda_1(\Omega')q = f_\perp & \text{on } \Omega' \\ q = 0 & \text{on } \partial\Omega', \end{cases}$$

which has a unique solution $q \in H_0^1(\Omega')$ with $\int q u_{\Omega'} = 0$, and enjoys the estimate $\|q\|_{H_0^1(\Omega')} \leq C(\alpha) \|f_\perp\|_{L^2}$ (this follows from the Fredholm alternative). From Lemma 3.2 we have that $\|u_{\Omega'}\|_{L^\infty} \leq C$, so

$$\|f_\perp\|_{L^\infty} \leq \|f\|_{L^2} \|u_{\Omega'}\|_{L^\infty} + \|f\|_{L^\infty} \leq C \|f\|_{L^\infty}.$$

So, from (3.3), we see that $-\Delta q - \lambda_1(\Omega')q \in L^\infty$, we have from applying [24, Theorem 8.17] in charts that

$$\|q\|_{L^\infty} \leq C [\|f_\perp\|_{L^\infty} + \|q\|_{L^2}] \leq C \|f\|_{L^\infty}.$$

Now rewriting the PDE as $-\Delta q = \lambda_1(\Omega')q + f_\perp \in L^\infty$, applying the comparison principle with a multiple of the torsion function $w_{\Omega'}$ gives

$$|q| \leq C \|f\|_{L^\infty} w_{\Omega'}.$$

This is the key fact about q .

We proceed in our estimate:

$$\int f_{\perp}(u_{\Omega'} - u_{\Omega}) = \int f_{\perp}u_{\Omega} = \int (-\Delta q - \lambda_1(\Omega')q)u_{\Omega} = \int q(-\Delta u_{\Omega} - \lambda_1(\Omega')u_{\Omega}),$$

where the last equality holds *in the sense of distributions*. From Lemma 3.1, on Ω' we have $\Delta u_{\Omega} = -\lambda_1(\Omega)u_{\Omega} + \mu$, where μ is a nonnegative measure supported on $\partial\Omega \cap \Omega'$. Continuing on,

$$\int q(-\Delta u_{\Omega} - \lambda_1(\Omega')u_{\Omega}) = \int -qd\mu + (\lambda_1(\Omega) - \lambda_1(\Omega')) \int qu_{\Omega}.$$

The second term is controlled by

$$(\lambda_1(\Omega) - \lambda_1(\Omega')) \left| \int qu_{\Omega} \right| \leq (\lambda_1(\Omega) - \lambda_1(\Omega')) \|q\|_{L^2} \|u_{\Omega}\|_{L^2} \leq C(\lambda_1(\Omega) - \lambda_1(\Omega')) \|f\|_{L^{\infty}},$$

so we need only focus on the first. Recall from Lemma 3.4 that writing $\Delta w_{\Omega} = -1_{\Omega} + \nu$ on Ω' for a nonnegative measure ν , then $\mu \leq C\nu$. As such,

$$\left| \int qd\mu \right| \leq \int |q|d\mu \leq C\|f\|_{L^{\infty}} \int w_{\Omega'}d\nu = C\|f\|_{L^{\infty}} \int (w_{\Omega'} + w_{\Omega})d\nu,$$

in the final identity using that w_{Ω} vanishes on the support of ν . Then using the equation for w_{Ω} ,

$$\begin{aligned} \int (w_{\Omega'} + w_{\Omega})d\nu &= \int (w_{\Omega'} + w_{\Omega})(\Delta w_{\Omega} + 1_{\Omega}) \leq \int (w_{\Omega'} + w_{\Omega})(\Delta w_{\Omega} + 1_{\Omega'}) \\ &= \int (w_{\Omega'} + w_{\Omega})(\Delta w_{\Omega} - \Delta w_{\Omega'}) = \int |\nabla w_{\Omega'}|^2 - |\nabla w_{\Omega}|^2. \end{aligned}$$

Recall that $\text{tor}(\Omega) = -\frac{1}{2} \int |\nabla w_{\Omega}|^2$ and similarly for Ω' , so we have shown that

$$\left| \int f_{\perp}(u_{\Omega'} - u_{\Omega}) \right| \leq C\|f\|_{L^{\infty}} [\lambda_1(\Omega) - \lambda_1(\Omega') + \text{tor}(\Omega) - \text{tor}(\Omega')].$$

Together with (3.2) this gives the conclusion.

If $\partial\Omega \cap \Omega'$ is not smooth, let $\Omega_t = \{u_{\Omega} > t\}$, which has smooth boundary for almost every t , and apply (3.1) to Ω_t . We have that $\lambda_1(\Omega_t) \geq \lambda_1(\Omega)$ and $\text{tor}(\Omega_t) \geq \text{tor}(\Omega)$ from set inclusion, and using $\frac{(u_{\Omega}-t)_+}{\|(u_{\Omega}-t)_+\|_{L^2}} \rightarrow u_{\Omega}$ in H^1 as a competitor for $\lambda_1(\Omega_t)$ gives that $\lim_t \lambda_1(\Omega_t) = \lambda_1(\Omega)$. On the other hand, taking first eigenfunctions $u_{\Omega_t} \rightharpoonup u \in H_0^1(\Omega)$ weakly and locally in C^k topology for some u along a subsequence, and similarly $w_{\Omega_k} \rightharpoonup w$. Passing the PDEs satisfied by these to the limit, we see that $-\Delta u = \lambda_1(\Omega)u$ and $-\Delta w = 1$ on Ω . This implies that $u_{\Omega} = u$ (recall we are assuming that $\lambda_1(\Omega)$ is simple) and $w_{\Omega} = w$. In particular, both sides of (3.1) pass to the limit. \square

3.3. Base energy minimizers in Q_R . We now move toward minimizing the base energy \mathcal{E} among subsets of a fixed Q_R . In this section we establish existence, volume bounds and connectedness properties of minimizers in the class $\mathcal{H} = \mathcal{H}_{R, v_{\max}}$.

We start with an existence and compactness theorem for minimizers. The statement is formulated in a way to give information on w_{Ω} if and only if $\mathfrak{T} > 0$. We recall that the base energy \mathcal{E} , the collection of sets over which we minimize \mathcal{H} , and the collection of minimizers \mathcal{M} are defined in (2.4), (2.5), and (2.7) respectively. Recall from (2.6) that $\mathcal{E}_{\min} = \mathcal{E}_{\min}(v, v_{\max}, \eta, \mathfrak{T}, R)$ denotes the infimum of the base energy.

Lemma 3.10. *Fix $R > 0$ and v, v_{\max} with $0 < v < v_{\max}$, as well as the parameters $\eta > 0$ and $\mathfrak{T} > 0$ in the base energy. Then:*

- (1) $\mathcal{E}_{\min} > -\infty$.
- (2) Let $\Omega_k \in \mathcal{H}$ and $u_k, w_k \in H_0^1(\Omega_k)$ with $\int u_k^2 = 1$ and

$$(3.4) \quad \lim_k \int |\nabla u_k|^2 + \mathfrak{T} \int \frac{1}{2} |\nabla w_k|^2 - w_k + f_{v, \eta}(|\Omega_k|) = \mathcal{E}_{\min}.$$

Then there is a subsequence Ω_{k_j} and a set $\Omega \in \mathcal{H}$ such that

$$\|u_{k_j} - u_{\Omega}\|_{H_0^1(Q_R)} + \sqrt{\mathfrak{T}} \|w_{k_j} - w_{\Omega}\|_{H_0^1(Q_R)} + |\Omega \Delta \Omega_k| \rightarrow 0,$$

where u_Ω is a normalized first eigenfunction of Ω and w_Ω is the torsion function of Ω .
(3) \mathcal{M} is nonempty, and for any $\Omega \in \mathcal{M}$, u_Ω and $\sqrt{\mathfrak{T}}w_\Omega$ are continuous functions (for any choice of u_Ω normalized first eigenfunction).

Proof. We first check that $\mathcal{E}_{min} > -\infty$. This amounts to showing that $\text{tor}(\Omega)$ is bounded from below uniformly in $\Omega \in \mathcal{H}$ (the other two terms are bounded from below), so it follows directly from Corollary 3.6.

Step 1: Basic compactness. Let Ω_k be a minimizing sequence as in claim (2) of the lemma. As \mathcal{E}_{min} is bounded from below, so are each of the terms $\int |\nabla u_k|^2$, $\mathfrak{T} \int \frac{1}{2} |\nabla w_k|^2 - w_k$, and $f_{v,\eta}(|\Omega_k|)$ individually. Because \mathcal{E}_{min} is also bounded from above, these three terms are bounded above individually as well. This immediately gives that $\|u_k\|_{H_0^1(Q_R)}$ is uniformly bounded, while applying Lemma 3.5 as in the proof of Corollary 3.6 gives that

$$C \geq \mathfrak{T} \int \frac{1}{2} |\nabla w_k|^2 - w_k \geq \mathfrak{T} \left[-C(R, v_{\max}) + \frac{1}{4} \int |\nabla w_k|^2 \right],$$

giving $\sqrt{\mathfrak{T}} \|\nabla w_k\|_{L^2(Q_R)} \leq C$ and so $\sqrt{\mathfrak{T}} \|w_k\|_{H_0^1(Q_R)} \leq C$ from Lemma 3.5 again.

Passing to subsequences, we have that $u_k \rightarrow u$ weakly in $H_0^1(Q_R)$, strongly in L^2 , and a.e., and similarly for $w_k \rightarrow w$ when $\mathfrak{T} > 0$. This implies that $1_{\{|u| + \mathfrak{T}|w| > 0\}} \leq \liminf_k 1_{\{|u_k| + \mathfrak{T}|w_k| > 0\}}$, so by Fatou's lemma $|\{ |u| + \mathfrak{T}|w| > 0 \}| \leq \liminf_k |\Omega_k|$ and

$$f_{v,\eta}(|\{ |u| + \mathfrak{T}|w| > 0 \}|) \leq \liminf_{k \rightarrow \infty} f_{v,\eta}(|\Omega_k|).$$

Using the lower semicontinuity of the norm under weak convergence, we have that

$$(3.5) \quad \int |\nabla u|^2 + \mathfrak{T} \int \frac{1}{2} |\nabla w|^2 - w + f_{v,\eta}(|\{ |u| + \mathfrak{T}|w| > 0 \}|) \leq \mathcal{E}_{min}$$

and $\int u^2 = 1$. From Lemmas 3.2 and 3.3, we have $|u_k| \leq C_1$ and $|w_k| \leq C_2$; these pass to the limit to give $|u| + \sqrt{\mathfrak{T}}|w| \leq C$.

Step 2: Continuous representatives. We will now show that u, w admit continuous representatives; this will imply that the set $\Omega := \{|u| + \mathfrak{T}|w| > 0\}$ is open (and hence in \mathcal{H}), u is a first eigenfunction of Ω , $w = w_\Omega$, and the inequality in (3.5) is an equality.

Step 2a. To this end, we first claim that u and w have the following ‘‘almost minimality’’ property. Fix a small ball $B_{2r}(x)$, $x \in Q_R$. We claim that

$$(3.6) \quad \int |\nabla u|^2 + \mathfrak{T} \int \frac{1}{2} |\nabla w|^2 - w \leq \frac{\int |\nabla a|^2}{\int a^2} + \mathfrak{T} \int \frac{1}{2} |\nabla b|^2 - b + Cr^n$$

for any pair of functions a, b with $a = u$ and $b = w$ on $B_r(x)^c$ and such that $u - a$ and $w - b$ are in $H_0^1(B_r(x) \cap Q_R)$.

Indeed, let ψ be a smooth cutoff function which is 1 on $B_r(x)$ and supported on $B_{2r}(x)$. Let $\Omega'_k = \Omega_k \cup B_r(x)$ and define the functions $a_k, b_k \in H_0^1(\Omega'_k)$ by

$$\begin{aligned} a_k &= \psi a + (1 - \psi)u_k, \\ b_k &= \psi b + (1 - \psi)w_k. \end{aligned}$$

Note that $a_k \rightarrow a$ and $b_k \rightarrow b$ in L^2 . Moreover, $f_{v,\eta}(|\Omega'_k|) \leq f_{v,\eta}(|\Omega_k|) + Cr^n$. So, we may use $a_k/\|a_k\|_{L^2}$ as a competitor for $u_{\Omega'_k}$ and b_k for $w_{\Omega'_k}$ inside of (2.2) and (2.1) to give

$$\mathcal{E}_{min} \leq \mathcal{E}(\Omega'_k) \leq \frac{\int |\nabla a_k|^2}{\int a_k^2} + \mathfrak{T} \int \frac{1}{2} |\nabla b_k|^2 - b_k + f_{v,\eta}(|\Omega_k|) + Cr^n.$$

Recalling that u_k, w_k were chosen to satisfy (3.4), this gives that

$$(3.7) \quad 0 \leq \liminf_{k \rightarrow \infty} \frac{\int |\nabla a_k|^2}{\int a_k^2} - \int |\nabla u_k|^2 + \mathfrak{T} \left[\int \frac{1}{2} |\nabla b_k|^2 - b_k - \int \frac{1}{2} |\nabla w_k|^2 - w_k \right] + Cr^n,$$

We first focus on the term in brackets. We have

$$\int \frac{1}{2} |\nabla b_k|^2 - \int \frac{1}{2} |\nabla w_k|^2 = \frac{1}{2} \int |\nabla w_k|^2 [(1 - \psi)^2 - 1] + \psi^2 |\nabla b|^2 + 2\psi(1 - \psi) \nabla b \cdot \nabla w_k + o_k(1),$$

where the $o_k(1)$ term contains all error terms containing $\nabla \psi$ and can be seen to converge to 0 using the weak convergence of $\nabla w_k \rightarrow \nabla w$, the strong convergence of $w_k \rightarrow w$, and the fact that $b = w$ on the support of $\nabla \psi$. Therefore, taking limits in the term in brackets in (3.7), since $w = b$ when $1 - \psi$ is nonzero, we have

$$\begin{aligned} \limsup_{k \rightarrow \infty} \int \frac{1}{2} |\nabla b_k|^2 - b_k - \int \frac{1}{2} |\nabla w_k|^2 - w_k &\leq \frac{1}{2} \int |\nabla w|^2 [(1 - \psi)^2 - 1] + \psi^2 |\nabla b|^2 + 2\psi(1 - \psi) |\nabla w|^2 + \int w - b \\ &= \frac{1}{2} \int \psi^2 [|\nabla b|^2 - |\nabla w|^2] + \int w - b \\ &= \int \frac{1}{2} |\nabla b|^2 - b - \int \frac{1}{2} |\nabla w|^2 - w, \end{aligned}$$

In the inequality we have used that $\int \phi |\nabla w|^2 \leq \liminf_{k \rightarrow \infty} \int \phi |\nabla w_k|^2$ for $\phi \geq 0$ from weak convergence (and setting $\phi = 1 - (1 - \psi^2)$). A similar computation for the a_k terms gives

$$\limsup_{k \rightarrow \infty} \frac{\int |\nabla a_k|^2}{\int a_k^2} - \int |\nabla u_k|^2 \leq \frac{\int |\nabla a|^2}{\int a^2} - \int |\nabla u|^2.$$

Substituting these conclusions into (3.7), we arrive at (3.6), thus proving the claim.

Step 2b. Now let a be the harmonic replacement of u on $B_r(x) \cap Q_R$ (i.e. the minimizer of $\int_{B_r(x) \cap Q_R} |\nabla a|^2$ with data u on $\partial[B_r(x) \cap Q_R]$), and b the harmonic replacement of w . Then

$$\int_{B_r(x)} |\nabla u|^2 + \mathfrak{I} \frac{1}{2} |\nabla w|^2 \leq Cr^n + \int_{B_r(x)} |\nabla a|^2 + \mathfrak{I} \frac{1}{2} |\nabla b|^2,$$

estimating the non-gradient terms using $|u|, |w|, |a|, |b| \leq C$ and absorbing them into the Cr^n term. Rewriting and using that $\int g(\nabla a, \nabla(u - a)) = 0$ from the equation on a (and similarly with b),

$$\int_{B_r(x)} |\nabla(u - a)|^2 + \mathfrak{I} \frac{1}{2} |\nabla(w - b)|^2 \leq Cr^n.$$

This is valid for any $B_{2r}(x)$, and implies that $u, \sqrt{\mathfrak{I}}w$ are $C^{0,\alpha}(Q_R)$ (see [15]).

Step 3: Conclusion. Finally, as Ω is open, $\Omega \in \mathcal{H}$ and $\mathcal{E}_{min} \leq \mathcal{E}(\Omega) \leq \int |\nabla u|^2 + \int \frac{1}{2} |\nabla w|^2 - w + f_{v,\eta}(|\Omega|) \leq \mathcal{E}_{min}$. This implies that $w = w_\Omega$, u is a first eigenfunction of Ω , and

$$\lim_{k \rightarrow \infty} \int |\nabla u_k|^2 = \int |\nabla u|^2 \quad \lim_{k \rightarrow \infty} \int |\nabla w_k|^2 = \int |\nabla w|^2 \quad \lim_{k \rightarrow \infty} |\Omega_k| = |\Omega|.$$

This gives the strong convergence of u_k and w_k in $H_0^1(Q_R)$. To see the convergence of Ω_k , recall that $1_\Omega \leq \liminf_k 1_{\Omega_k}$; integrating over Ω and applying Fatou's lemma gives $|\Omega| \leq \liminf_k |\Omega \cap \Omega_k|$, or $|\Omega \setminus \Omega_k| \rightarrow 0$. Together with $|\Omega_k| \rightarrow |\Omega|$ this guarantees $|\Omega_k \setminus \Omega| \rightarrow 0$ as well.

For (3), (1) and (2) immediately give that \mathcal{M} is nonempty by applying with $u_k = u_{\Omega_k}$, an eigenfunction of Ω_k , and $w_k = w_{\Omega_k}$, where Ω_k is any sequence in \mathcal{H} with $\mathcal{E}(\Omega_k) \rightarrow \mathcal{E}_{min}$. Choosing $\Omega_k = \Omega$ and $u_{\Omega_k} = u_\Omega$ any first eigenfunction and applying the continuity argument above gives that u_Ω and $\sqrt{\mathfrak{I}}w_\Omega$ are continuous. \square

Recall that the volume penalization term $f_{v,\eta}$ in the base energy does not immediately guarantee that a minimizer $\Omega \in \mathcal{M}$ satisfies the desired volume constraint $|\Omega| \leq v$. However, the following lemma provides an initial upper bound on the volume of a minimizer, provided that the parameter η in the volume penalization term is taken to be sufficiently small. We will eventually show in Proposition 9.5 that and $\Omega \in \mathcal{M}$ actually satisfies the desired volume constrain $|\Omega| = v$.

Lemma 3.11. *Fix $R > 0$ and v, v_{max} with $0 < v < v_{max}$, and $0 < \mathfrak{T} \leq 1$. There exists an $\eta_0 = \eta_0(v, v_{max}, R) > 0$ so that for $\eta < \eta_0$, if $\Omega \in \mathcal{M}$ then*

$$|\Omega| < \frac{v_{max} + v}{2}.$$

Proof. We have already verified in the proof of Lemma 3.10 that $\mathcal{E}(\Omega) \geq -C + f_{v,\eta}(|\Omega|)$ for any $\Omega \in \mathcal{H}$; this estimate is uniform for any $\mathfrak{T} \leq 1$. By choosing some smooth Ω with precisely $|\Omega| = v$, we see that $\mathcal{E}(\Omega) \leq C$ with C independent of η (and \mathfrak{T}). We therefore have, for $\Omega \in \mathcal{M}$,

$$\frac{|\Omega| - v}{\eta} \leq f_{v,\eta}(|\Omega|) \leq C.$$

Choosing $\eta \leq \frac{v_{max} - v}{2C}$ gives the conclusion. \square

Remark 3.12. In the remainder of the paper, we will *always* assume, without further comment, that the parameter in the base energy \mathcal{E} and the main energy \mathcal{F}_τ is taken such that $\eta \leq \eta_0$, where $\eta_0 = \eta_0(v, v_{max}, R)$ is the constant obtained in Lemma 3.11.

In the next lemma, we show that the coefficient \mathfrak{T} in front of the torsional rigidity in the base energy can be taken to be sufficiently small to guarantee that a minimizer $\Omega \in \mathcal{M}$ is connected and its first eigenvalue is simple. Although the simplicity of the first eigenvalue follows from the connectedness, we prove these two properties separately so that the proof can be immediately generalized to prove Lemma 3.14 below.

Lemma 3.13. *Fix $R > 0$, $v < v_{max}$, and $\eta > 0$. There exists a $\mathfrak{T}_0(v, v_{max}, \eta, R) > 0$ such that if $\Omega \in \mathcal{M}$ with $\mathfrak{T} < \mathfrak{T}_0$, then Ω is connected, $u_\Omega > 0$ on Ω (up to changing sign), and*

$$\lambda_2(\Omega) > \lambda_1(\Omega).$$

Proof. Consider first the case of $\mathfrak{T} = 0$. If $\lambda_2(\Omega) = \lambda_1(\Omega)$, then there are two functions u_1, u_2 with $\int |\nabla u_i|^2 = \lambda_1(\Omega)$, $\int u_i^2 = 1$, and $\int u_1 u_2 = 0$. We know from Lemma 3.10 that they are both continuous. We may assume without loss of generality that $u_i \geq 0$ by replacing u_i with $|u_i|$ in the Rayleigh quotient characterization of eigenfunctions. So, the orthogonality shows that $\{u_1 > 0\}$ and $\{u_2 > 0\}$ are disjoint open sets.

Now let $\Omega' = \{u_1 > 0\}$, noting that $\lambda_1(\Omega') \leq \int |\nabla u_1|^2 = \lambda_1(\Omega)$. On the other hand, $|\Omega'| < |\Omega|$, so $f_{v,\eta}(|\Omega'|) < f_{v,\eta}(|\Omega|)$ and $\mathcal{E}(\Omega') < \mathcal{E}(\Omega) = \mathcal{E}_{min}$. This is a contradiction. If Ω is disconnected, then $u_\Omega = u_E$ where E is one of its connected components; the same argument shows that $\mathcal{E}(E) < \mathcal{E}(\Omega)$.

To handle the case of $\mathfrak{T} > 0$, argue by contradiction. Assume there is a sequence of sets $\Omega_k \in \mathcal{M}_{\mathfrak{T}_k}$, $\mathfrak{T}_k \rightarrow 0$, with $\lambda_2(\Omega_k) = \lambda_1(\Omega_k)$ or Ω_k disconnected. Then

$$\begin{aligned} \lambda_1(\Omega_k) + \mathfrak{T}_k \text{tor}(\Omega_k) + f_{v,\eta}(\Omega_k) &\leq \lambda_1(\Omega) + \mathfrak{T}_k \text{tor}(\Omega) + f_{v,\eta}(\Omega) \\ &\rightarrow \lambda_1(\Omega) + f_{v,\eta}(\Omega) \end{aligned}$$

for any $\Omega \in \mathcal{H}$. As $\text{tor}(\Omega) \geq -C$ from Corollary 3.6, each of $\lambda_1(\Omega_k)$, $\text{tor}(\Omega_k)$, and $f_{v,\eta}(\Omega_k)$ are bounded uniformly in k , so the middle term $\mathfrak{T}_k \text{tor}(\Omega_k) \rightarrow 0$. This means

$$\begin{aligned} \mathcal{E}_{min}(\mathfrak{T}_k) &= \lambda_1(\Omega_k) + f_{v,\eta}(\Omega_k) + o_k(1) \\ &\leq \inf\{\lambda_1(\Omega) + f_{v,\eta}(\Omega) : \Omega \in \mathcal{H}\} + o_k(1) = \mathcal{E}_{min}(0) + o_k(1). \end{aligned}$$

In particular, $\mathcal{E}_{min}(\mathfrak{T}_k) \rightarrow \mathcal{E}_{min}(0)$ and $\lim_k \lambda_1(\Omega_k) + f_{v,\eta}(\Omega_k) = \mathcal{E}_{min}(0)$. Fix first eigenfunctions u_{Ω_k} for Ω_k ; as above we may take them nonnegative. Applying the compactness claim of Lemma 3.10, there is an $\Omega \in \mathcal{M}_0$ with $\|u_{\Omega_k} - u_\Omega\|_{H_0^1(Q_R)} \rightarrow 0$ and $|\Omega_k \triangle \Omega| \rightarrow 0$.

First, suppose that $\lambda_2(\Omega_k) = \lambda_1(\Omega_k)$ along some subsequence. Let v_k be functions in $H_0^1(\Omega_k)$ with $v_k \geq 0$, $\int v_k u_{\Omega_k} = 0$, $\int v_k^2 = 1$, and $\int |\nabla v_k|^2 = \lambda_1(\Omega_k) \rightarrow \lambda_1(\Omega)$. We may apply Lemma 3.10 to v_k instead of u_{Ω_k} to obtain another set $\Omega' \in \mathcal{M}_0$ and a continuous first eigenfunction of Ω' , $u_{\Omega'}$, such that $\|v_k - u_{\Omega'}\|_{H_0^1(Q_R)} \rightarrow 0$ and $|\Omega_k \triangle \Omega'| \rightarrow 0$. This immediately gives that $|\Omega \triangle \Omega'| = 0$, and so $\Omega \cup \Omega'$ is in \mathcal{H} ,

$\lambda_1(\Omega' \cup \Omega) \leq \lambda_1(\Omega)$, and $|\Omega' \cup \Omega| = |\Omega|$. In particular, $\Omega \cup \Omega' \in \mathcal{M}_0$ and $\lambda_1(\Omega' \cup \Omega) = \lambda_1(\Omega)$. This means both $u_{\Omega'}$ and u_{Ω} are first eigenfunctions of $\Omega \cup \Omega'$, and also

$$\int u_{\Omega} u_{\Omega'} = \lim_{k \rightarrow \infty} \int u_{\Omega_k} v_k = 0.$$

Thus $\lambda_2(\Omega \cup \Omega') = \lambda_1(\Omega \cup \Omega')$, which is a contradiction.

Next, assume Ω_k is disconnected, and let E_k^i be its connected components. We claim that $|E_k^i|$ converges either to $|\Omega|$ or to 0. Indeed, let $u_k \geq 0$ be the first eigenfunction of Ω_k , and E_k^1 be the (unique) connected component of Ω_k on which it is nonzero. We know that $u_k \rightarrow u_{\Omega}$ a.e., which is strictly positive on Ω ; this gives that $|E_k^1| \rightarrow |\Omega|$ while $|E_k^i| \rightarrow 0$ for any other i .

It remains to check that for any Ω_k , if $|E_k^i|$ is sufficiently small, then it is empty. The basic idea is that the torsional rigidity is additive on connected components, and a component of small volume contributes too much torsional rigidity for a set to be minimizing. More specifically, we consider the quantity

$$a(t) := \inf\{\text{tor}(E) : E \in \mathcal{H}, |E| \leq t\}.$$

From Lemma 3.1(3), we have $\int |\nabla w_E|^2 = \int w_E = -2\text{tor}(E)$ (using w_E as a test function for itself), and therefore

$$-\text{tor}(E) = \frac{1}{2} \int w_E \leq |E|^{\frac{n+2}{2n}} \|w_E\|_{L^{\frac{2n}{n-2}}} \leq C|E|^{\frac{n+2}{2n}} \|\nabla w_E\|_{L^2} \leq C|E|^{\frac{n+2}{2n}} \sqrt{-\text{tor}(E)},$$

where the constant $C = C(R, v_{\max})$ introduced in the second-to-last inequality comes from the Poincaré inequality of Lemma 3.5. Dividing and squaring, $\text{tor}(E) \geq -C|E|^{\frac{n+2}{n}}$, so

$$(3.8) \quad a(t) \geq -Ct^{\frac{n+2}{n}}.$$

Let E_k^1 be the large connected component of Ω_k and set $|\Omega_k| = |E_k^1| + t_k$. We know that $t_k \rightarrow 0$ and assume for contradiction that $t_k > 0$. Since $\lambda_1(\Omega_k) = \lambda_1(E_k^1)$, we take E_k^1 as a competitor for Ω_k in the minimization of $\mathcal{E}_{\mathfrak{T}_k}$ to find

$$\mathfrak{T}_k \text{tor}(\Omega_k) + f_{v,\eta}(|\Omega_k|) \leq \mathfrak{T}_k \text{tor}(E_k^1) + f_{v,\eta}(|E_k^1|).$$

Moreover, for disjoint A, B , $\text{tor}(A \cup B) = \text{tor}(A) + \text{tor}(B)$, which gives

$$\mathfrak{T}_k \text{tor}(\Omega_k \setminus E_k^1) + f_{v,\eta}(|\Omega_k|) \leq f_{v,\eta}(|E_k^1|).$$

Consequently, we deduce that

$$\mathfrak{T}_k a(t_k) \leq f_{v,\eta}(|E_k^1|) - f_{v,\eta}(|\Omega_k|) \leq -\eta t_k,$$

so applying (3.8) gives $\eta t_k \leq C\mathfrak{T}_k t_k^{\frac{2+n}{n}} \leq C t_k^{\frac{2+n}{n}}$. For large k , this is a contradiction. \square

The same compactness argument can be applied to “approximate” minimizers instead, except for the full connectedness conclusion. We omit the proof.

Lemma 3.14. *Fix $R > 0$, $0 < v < v_{\max}$, $\eta > 0$ and $\varepsilon > 0$. There exist $\mathfrak{T}_0 = \mathfrak{T}_0(v, v_{\max}, \eta, R) > 0$ and $\delta = \delta(v, v_{\max}, R, \eta, \varepsilon) > 0$ such that if $\mathfrak{T} < \mathfrak{T}_0$, $\Omega \in \mathcal{H}$, and*

$$\mathcal{E}(\Omega) \leq \mathcal{E}_{\min} + \delta,$$

then:

- (1) *There is an $\Omega' \in \mathcal{M}$ with $|\Omega' \triangle \Omega| < \varepsilon$, $\|u_{\Omega'} - u_{\Omega}\|_{H^1} \leq \varepsilon$, and $\sqrt{\mathfrak{T}}\|w_{\Omega'} - w_{\Omega}\|_{H^1} \leq \varepsilon$.*
- (2) *$\lambda_2(\Omega) > \lambda_1(\Omega) + c(v, v_{\max}, R)$.*
- (3) *One connected component E of Ω has $|E| \geq |\Omega| - \varepsilon$, while every other (if any) has $|E| \leq \varepsilon$.*

Note that under the assumptions of this lemma, it is simply false that Ω must be connected: one may always add a ball of small measure to Ω , which perturbs \mathcal{E} continuously.

3.4. Global minimizers of the base energy. In general, the constants in the previous section necessarily depend on R , and there is no reason to expect, for instance, existence of global minimizers on noncompact Riemannian manifolds with arbitrary (bounded geometry) behavior at infinity. In this section, we turn our attention to Riemannian manifolds (M, g) where M/G is compact. In this case, most of the above estimates remain valid with constants independent of R , and minimizers of the base energy \mathcal{E} among all open bounded sets exist and have bounded diameter. For the present section, we focus only on estimates which will be relevant in later sections; in particular, the existence of minimizers is not treated here (but does follow from Theorem 6.3 later).

Of course, in the case when (M, g) is compact, then for $M \subset Q_R$ for sufficiently large R , and thus the constants in the previous section may be taken independent of R trivially. So, to simplify statements, in this section we will assume that (M, g) is noncompact but $(M/G, g)$ is compact. Note in this case that M has infinite volume.

For our first estimate, we show some initial bounds on the infimum of the energy and on the torsion function of any open bounded set in terms of its volume. Observe that $\mathcal{E}_{\min}(R)$ is a nonincreasing function of R , as the class \mathcal{H} is increasing in R . Set $\mathcal{E}_{\min}(\infty) = \lim_{R \rightarrow \infty} \mathcal{E}_{\min}(R)$.

Lemma 3.15. *Assume M/G is compact and fix $v < v_{\max} < \infty$. Then there exists $C = C(v_{\max}) < \infty$ such that for any bounded open set Ω with $|\Omega| \leq v_{\max}$, we have $\mathcal{E}_{\min}(\infty) \geq -C$ and $w_\Omega \leq C$.*

Proof. First, we claim there exists a large $S \gg 0$ such that

$$\bigcup_{e \in G} e(Q_S) = M.$$

If not, for every k there is an $x_k \in M$ such that $e(x_k) \notin Q_k$ for any $e \in G$. There exists, however, a subsequence x_{k_j} and isometries e_j with $e_j(x_{k_j}) \rightarrow x \in M$. If $x \in Q_J$ for some J , then as Q_J is open $e_j(x_{k_j}) \in Q_J$ for j large enough; this is impossible if $k_j \geq J$, so $x \notin Q_J$. This contradicts $M = \bigcup_R Q_R$. Using the Vitali covering lemma and possibly taking S larger, there exists a countable collection $L \subseteq G$ such that $\bigcup_{e \in L} e(Q_S) = M$ and $\{e(Q_S)\}_{e \in L}$ have finite overlap. Up to possibly increasing S depending on v_{\max} , we also assume that $|Q_S| > v_{\max}$.

Now, fix any Ω open, bounded, with $|\Omega| \leq v_{\max}$. For any $e \in L$, we apply the Poincaré inequality of Lemma 3.5 to $e(Q_S)$ to obtain that

$$(3.9) \quad \|w_\Omega\|_{L^2(\Omega \cap e(Q_S))} \leq C \|\nabla w_\Omega\|_{L^2(\Omega \cap e(Q_S))}$$

for a constant $C = C(S, v_{\max}) = C(v_{\max})$, using that $|\Omega \cap e(Q_S)| \leq v_{\max} < |Q_S|$. Summing over all $e \in L$ and using the finite overlapping property gives $\int_\Omega w_\Omega^2 \leq C \int_\Omega |\nabla w_\Omega|^2$. Then

$$\int_\Omega w_\Omega \leq C |\Omega|^{1/2} \|\nabla w_\Omega\|_{L^2(\Omega)} \leq C \|\nabla w_\Omega\|_{L^2(\Omega)},$$

so $\text{tor}(\Omega) = \int \frac{1}{2} |\nabla w_\Omega|^2 - w_\Omega \geq -C$. On the other hand $\lambda_1(\Omega) > 0$, while $f_{v, \eta}(\Omega) \geq -v_{\max}$, so $\mathcal{E}(\Omega) \geq -C(v_{\max})$. Taking the infimum over all Ω gives the first conclusion.

To see that $w_\Omega \leq C$, we now have from (3.9), elliptic estimates [24, Theorem 8.17], and a basic covering argument that

$$\|w_\Omega\|_{L^\infty(\Omega \cap e(Q_S))} \leq C[\|w_\Omega\|_{L^2(Q_{2S})} + 1] \leq C.$$

Applying to every e gives the estimate. \square

We now establish an analogue of Lemma 3.14 that is independent of R : that any low energy set can be well approximated by a minimizer (on some Q_S for S uniform) and has one “large” connected component. We use a concentration compactness argument to handle the loss of compactness coming from the isometries of the space.

Lemma 3.16. *Assume M/G is compact. Fix $v < v_{\max} < \infty$, $\varepsilon > 0$, and $\eta > 0$. There exist $\mathfrak{T}_0 = \mathfrak{T}_0(v, v_{\max}, \eta) > 0$, and a $\delta = \delta(v, v_{\max}, \eta, \varepsilon) > 0$ such that if $\mathfrak{T} < \mathfrak{T}_0$, then the following holds. For any open bounded set $\Omega \subset M$ with $|\Omega| \leq v_{\max}$ and $\mathcal{E}(\Omega) \leq \mathcal{E}_{\min}(\infty) + \delta$:*

- (1) There is an $S = S(v, v_{max}, \eta, \varepsilon)$ and a $U \in \mathcal{M}(S)$ (i.e. a minimizer of \mathcal{E} over open $\Omega \subseteq Q_S$ with $|\Omega| \leq v_{max}$) such that $|U \Delta \Omega| < \varepsilon$, $\|u_U - u_\Omega\|_{H^1} \leq \varepsilon$, and $\sqrt{\mathfrak{T}} \|w_U - w_\Omega\|_{H^1} \leq \varepsilon$.
- (2) $\lambda_2(\Omega) > \lambda_1(\Omega) + c(v, v_{max})$
- (3) One connected component E of Ω has $|E| \geq |\Omega| - \varepsilon$, while every other (if any) has $|E| \leq \varepsilon$.

Proof. Let $L \subseteq G$ and S be as in the proof of Lemma 3.15. We use the shorthand $\bar{Q} = Q_S$. Let $\{\psi_e\}_{e \in L}$ a partition of unity subordinate to the cover of M by $\{e(\bar{Q})\}_{e \in L}$ with $\psi_e = \psi \circ e^{-1} \in C_c^\infty(\bar{Q})$ translates of each other. Throughout this proof, C and c denote constants depending only on v_{max} (and, as usual, on g) whose values change from line to line. Let $\delta \leq 1$ be a fixed number to be specified in the proof and let Ω be as above. Applying Lemma 3.15, we know that

$$(3.10) \quad \int_{\Omega} |\nabla u_\Omega|^2 + u_\Omega^2 + |\nabla w_\Omega|^2 + w_\Omega^2 \leq C.$$

The first step of the proof is a concentration compactness argument that allows us to replace Ω by a uniformly bounded set, losing only a small amount of mass and increasing the base energy only a small amount.

Step 1: There exist \mathfrak{T}_0 and $\bar{J} = \bar{J}(v_{max}, \delta) \in \mathbb{N}$ such that we may choose $x_0 \in M$ and $j \leq \bar{J}$ depending on Ω such that

$$(3.11) \quad \begin{cases} \mathcal{E}(\Omega \cap B_{8Sj}(x_0)) \leq \mathcal{E}_{min} + 2\delta \\ |\Omega \setminus B_{8Sj}(x_0)| \leq 4\eta\delta. \end{cases}$$

Because this step is rather involved, we divide it into several substeps.

Step 1a: Selection of x_0 . We show that Ω has a (uniformly) nontrivial amount of mass in $e_0(\bar{Q})$ for some $e_0 \in L$, and choose x_0 in this set. More specifically, for any $e \in L$, we apply Lemma 3.5 to $\Omega \cap e(\bar{Q})$ to find

$$\int_{e(\bar{Q})} u_\Omega^2 \leq |\Omega \cap e(\bar{Q})|^{\frac{2}{n}} \left(\int_{e(\bar{Q})} u_\Omega^{\frac{2n}{n-2}} \right)^{\frac{n-2}{n}} \leq C |\Omega \cap e(\bar{Q})|^{\frac{2}{n}} \int_{e(\bar{Q})} |\nabla u_\Omega|^2.$$

Summing over $e \in L$, recalling the finite overlap property of L and (3.10), we have

$$1 = \int_{\Omega} u_\Omega^2 \leq C \left(\sup_{e \in L} |\Omega \cap e(\bar{Q})|^{\frac{2}{n}} \right) \int_{\Omega} |\nabla u_\Omega|^2 \leq C \left(\sup_{e \in L} |\Omega \cap e(\bar{Q})|^{\frac{2}{n}} \right).$$

Consequently, $|\Omega \cap e_0(\bar{Q})| \geq c$ for some e_0 . We let x_0 be any fixed point in $e_0(\bar{Q})$.

Step 1b: Selecting j for each $J \in \mathbb{N}$. Next, we show how, for each $J \in \mathbb{N}$, to suitably choose $j \leq J$ that will ultimately lead to estimates of the form (3.11) but with a C/J error; we will later choose J depending on δ to absorb this term. For any $k \in \mathbb{N}$, we denote by L_k the finite collection of isometries such that $e(\bar{Q})$ lies in the annulus $B_{16kS}(x_0) \setminus B_{8kS}(x_0)$. That is,

$$L_k = \{e \in L : e(\bar{Q}) \cap [B_{8kS}(x_0) \cup (M \setminus B_{16kS}(x_0))] = \emptyset\}.$$

Since $\text{diam } \bar{Q} \leq 2S$ by assumption, the collections L_k are pairwise disjoint, and so from (3.10) we find that

$$\sum_{k=1}^{\infty} \sum_{e \in L_k} \int_{e(\bar{Q}) \cap \Omega} |\nabla u_\Omega|^2 + u_\Omega^2 + |\nabla w_\Omega|^2 + w_\Omega^2 \leq C.$$

Since this sum is finite, it follows that for every $J > 0$, we may find some $j \leq J$ such that

$$\sum_{e \in L_j} \int_{e(\bar{Q}) \cap \Omega} |\nabla u_\Omega|^2 + u_\Omega^2 + |\nabla w_\Omega|^2 + w_\Omega^2 \leq \frac{C}{J}.$$

In other words, the functions u_Ω and w_Ω have most of their mass and energy outside of the annulus corresponding to L_j . Let this $j = j(J)$ be fixed.

Step 1c: Truncation and intermediate bounds. In this step, we consider $J \in \mathbb{N}$ fixed and $j = j(J)$ as determined in the previous step. Let us now define two smooth truncations of the eigenfunction u_Ω , the “inner part” u_I in the ball B_{8jS} and the “outer part” u_O in the complement of B_{16jS} , by setting

$$\begin{aligned} u_I &= \sum_{e \in L_I} \psi_e u_\Omega & L_I &= \{e \in L : e(\bar{Q}) \cap B_{8jS}(x_0) \neq \emptyset\}, \\ u_O &= \sum_{e \in L_O} \psi_e u_\Omega & L_O &= \{e \in L : e(\bar{Q}) \cap (M \setminus B_{16jS}(x_0)) \neq \emptyset\}. \end{aligned}$$

Notice that the functions u_I, u_O have disjoint support and the remaining “annular part” of u_Ω is given by $u_\Omega - u_I + u_O = \sum_{e \in L_k} \psi_e u_\Omega$. The truncations u_I and u_O capture most of the energy of u_Ω in the following sense:

$$\begin{aligned} \left| \int |\nabla u_\Omega|^2 - |\nabla u_O|^2 - |\nabla u_I|^2 \right| &= \left| \int g(\nabla[u_\Omega - u_O - u_I], \nabla[u_\Omega + u_O + u_I]) \right| \\ &\leq \sum_{e \in L_j} (|\nabla \psi_e| |u_\Omega| + \psi_e |\nabla u_\Omega|) \cdot C |\nabla u_\Omega| \\ &\leq C \sum_{e \in L_j} \int_{e(\bar{Q})} |u_\Omega|^2 + |\nabla u_\Omega|^2 \leq \frac{C}{J}. \end{aligned}$$

In particular, this together with the Poincaré inequality applied to u_O provide the following upper bound for the energy of u_I :

$$(3.12) \quad \int |\nabla u_I|^2 \leq \lambda_1(\Omega) - \int |\nabla u_O|^2 + \frac{C}{J} \leq \lambda_1(\Omega) \left(1 - \int u_O^2\right) + \frac{C}{J}.$$

In a similar fashion, we find that

$$(3.13) \quad \int u_\Omega^2 - u_O^2 - u_I^2 \leq \frac{C}{J}.$$

The first main estimate of this sub-step is the following: there exists C such that

$$(3.14) \quad \begin{aligned} \text{if } \int u_I^2 \geq \frac{1}{4} & \quad \text{then } \lambda_1(\Omega \cap B_{8jS}(x_0)) \leq \lambda_1(\Omega) + \frac{C}{J}, \\ \text{if } \int u_O^2 \geq \frac{1}{4} & \quad \text{then } \lambda_1(\Omega \setminus B_{8jS}(x_0)) \leq \lambda_1(\Omega) + \frac{C}{J}. \end{aligned}$$

Note that (3.13) guarantees that one of the two alternatives in (3.14) must hold provided J is sufficiently large. If $\int u_I^2 \geq 1/4$ then (3.13) implies that $(1 - \int u_O^2) / \int u_I^2 \leq 1 + C/J$. So, if we multiply (3.12) by $(1 - \int u_O^2) / \int u_I^2$ and divide both sides by $1 - \int u_O^2$, we arrive at

$$\frac{\int |\nabla u_I|^2}{\int u_I^2} \leq \lambda_1(\Omega) + \frac{C}{J}.$$

Here we have also used the fact that $\lambda_1(\Omega)$ is uniformly bounded; recall (3.10). So, the first estimate of (3.14) follows by taking u_I as a test function for $\lambda_1(\Omega \cap B_{8jS}(x_0))$. The proof of the second estimate of (3.14) is the same, using the analogous estimate to (3.12) for u_O .

The second main estimate of this substep is the following:

$$(3.15) \quad \begin{aligned} \text{tor}(\Omega \cap B_{8jS}(x_0)) &\leq \text{tor}(\Omega) + \frac{C}{J} + C|\Omega \setminus B_{8jS}(x_0)|, \\ \text{tor}(\Omega \setminus B_{8jS}(x_0)) &\leq \text{tor}(\Omega) + \frac{C}{J} + C|\Omega \cap B_{8jS}(x_0)| \end{aligned}$$

The estimate (3.15) follows similarly to (3.14). Splitting the torsion function w_Ω into an “inner part” w_I , and “outer part” w_O , and an “annular part” $w_\Omega - w_I - w_O$ exactly as above leads to

$$-\frac{1}{2} \int w_I + w_O \leq \text{tor}(\Omega) + \frac{C}{J};$$

we omit the details. This estimate is potentially much worse than the eigenvalue one due to the presence of the w_O term (the torsional rigidity has poorer concentration-compactness behavior than the eigenvalue), but recalling that $w_\Omega \leq C$ from Lemma 3.15, we find that

$$\begin{aligned} \operatorname{tor}(\Omega \cap B_{8Sj}(x_0)) &\leq -\frac{1}{2} \int w_I \leq \operatorname{tor}(\Omega) + \frac{C}{J} + \frac{1}{2} \int w_O \\ &\leq \operatorname{tor}(\Omega) + \frac{C}{J} + C|\Omega \setminus B_{8Sj}(x_0)|, \end{aligned}$$

thus proving the first estimate in (3.15). The second estimate in (3.15) is proven in the same way.

Step 1d: Energy comparison and selection of \bar{J} and \mathfrak{T}_0 . Fix any $J \in \mathbb{N}$ and let $j = j(J)$ be as determined in Step 1b. Let us first consider the case that $\int u_I^2 \geq \frac{1}{4}$. In this case, we may use the first estimates in (3.14) and (3.15) to compare $\mathcal{E}(\Omega)$ to $\mathcal{E}(\Omega \cap B_{8Sj}(x_0))$, which gives

$$\begin{aligned} \mathcal{E}(\Omega) &\leq \mathcal{E}_{\min} + \delta \leq \mathcal{E}(\Omega \cap B_{8Sj}(x_0)) + \delta \\ &\leq \lambda_1(\Omega) + \mathfrak{T} \operatorname{tor}(\Omega) + \frac{C}{J} + f_{v,\eta}(|\Omega \cap B_{8Sj}(x_0)|) + \delta + C\mathfrak{T}|\Omega \setminus B_{8Sj}(x_0)| \\ &\leq \mathcal{E}(\Omega) + \frac{C}{J} + f_{v,\eta}(|\Omega \cap B_{8Sj}(x_0)|) - f_{v,\eta}(|\Omega|) + \delta + C\mathfrak{T}|\Omega \setminus B_{8Sj}(x_0)| \\ &\leq \mathcal{E}(\Omega) + \frac{C}{J} + \delta - [\eta - C\mathfrak{T}]|\Omega \setminus B_{8Sj}(x_0)|. \end{aligned}$$

At this point, we can fix our remaining parameters. Choosing $\mathfrak{T}_0 < \frac{\eta}{2C}$, the last term is negative and this leads to

$$(3.16) \quad \begin{cases} \mathcal{E}(\Omega \cap B_{8Sj}(x_0)) \leq \mathcal{E}_{\min} + \frac{C}{J} + \delta \leq 2\delta \\ |\Omega \setminus B_{8Sj}(x_0)| \leq 4\eta\delta. \end{cases}$$

Finally, we choose $J = \bar{J}$ large enough in terms of δ so that $\frac{C}{J} \leq \delta$, and (3.11) follows. This concludes the proof of Step 1 in the case when $\int u_I^2 \geq \frac{1}{4}$. Now, assume instead that $\int u_O^2 \geq \frac{1}{4}$. We may argue in the analogous fashion to find that

$$(3.17) \quad \begin{cases} \mathcal{E}(\Omega \setminus B_{8Sj}(x_0)) \leq \mathcal{E}_{\min} + \frac{C}{J} + \delta \leq 2\delta \\ |\Omega \cap B_{8Sj}(x_0)| \leq 4\eta\delta. \end{cases}$$

This contradicts Step 1a for δ chosen sufficiently small. We conclude that $\int u_I^2 \geq \frac{1}{4}$ holds and thus have completed the proof of Step 1.

Step 2: Conclusion. We are now in a position to prove the first conclusion of the lemma. Let $\Omega' = \Omega \cap B_{8Sj}(x_0)$ be the truncation of Ω obtained in Step 1. Let S' be chosen large enough that $e_0^{-1}(B_{8Sj}(x_0)) \subseteq Q_{S'}$; note that this depends only on J and S , so only on v_{\max} and δ . Applying Lemma 3.10 to $e_0^{-1}(\Omega')$ on $Q_{S'}$, there is a $U \in \mathcal{M}(S')$ with $\|u_U - u_I\|_{H^1}, \sqrt{\mathfrak{T}}\|w_U - w_I\|_{H^1} \leq \varepsilon/2$. On the other hand, from the Caccioppoli inequality and letting $L' = \{e \in L : e(Q_S) \cap B_{8jS}(x_0) = \emptyset\}$,

$$\int |\nabla(u_\Omega - u_I)|^2 \leq C \sum_{e \in L'} \int |\nabla \psi_e|^2 |u_\Omega|^2 + \psi_e^2 |\nabla u_\Omega|^2 \leq C \int_{\Omega \setminus B_{8Sj}(x_0)} |u_\Omega|^2 \leq C\delta,$$

using that $|u_\Omega| \leq C$ and the volume estimate in (3.11) at the end. Similarly, $\int |\nabla(w_\Omega - w_I)|^2 \leq C\delta$, and we obtain conclusion (1) from the triangle inequality.

To prove (2), note that applying Lemma 3.14 gives that $\lambda_2(\Omega') \geq \lambda_1(\Omega) + c_0$ for $c_0 = c_0(S', v, v_{\max})$, and so it remains to show that $\lambda_2(\Omega') \leq \lambda_2(\Omega) + c_0/2$. To this end, let u_2 be a second eigenfunction for Ω . From Lemma 3.2 we have that $|u_2| \leq C(\lambda_1(\Omega)) \leq C$ unless $\lambda_2(\Omega) > 1 + \lambda_1(\Omega)$ (in which case we would be done). Then set $u_{2,I} = \sum_{e \in L_I} \psi_e u_2$ to be the ‘‘inner part,’’ i.e. the analogue of u_I for u_2 . As $\Delta|u_2| \geq -\lambda_2(\Omega)|u_2|$ in the sense of distributions, the same Caccioppoli inequality argument gives that

$$\int |\nabla(u_2 - u_{2,I})|^2 \leq C \int_{\Omega \setminus B_{8Sj}(x_0)} |u_2|^2 \leq C\delta.$$

Hence

$$\frac{\int |\nabla u_{2,I}|^2}{\int u_{2,I}^2} \leq \lambda_2(\Omega) + C\delta \quad \text{and} \quad \frac{|\int u_{2,I} u_I|}{\|u_{2,I}\|_{L^2} \|u_I\|_{L^2}} \leq C\delta,$$

implying that $\lambda_2(\Omega') \leq \lambda_2(\Omega) + C\delta$. Taking δ small enough completes the proof of (2).

To check the final conclusion (3) of the lemma, note that from Lemma 3.14, we know that one connected component A of Ω' has $|A| > |\Omega'| - \varepsilon/2$, while every other has $|A| \leq \varepsilon/2$ (choosing δ small enough in terms of ε). As $|\Omega \setminus B_{8S_j}(x_0)| \leq C\delta$, we may conclude by choosing $C\delta < \varepsilon/2$. \square

4. THE LOWER BOUND

In this section we consider inward minimizers of the main functional \mathcal{F}_τ defined in (2.8). The main result of this section is Theorem 4.1, which provides a linear lower bound for the growth of the function $u_\Omega + \sqrt{\mathfrak{T}}w_\Omega$ away from the boundary of an inward minimizer Ω . The theorem is phrased in terms of a lower bound on the quantity $DO(\Omega)$ defined in (2.14). One important consequence of this theorem is a lower volume density estimate for inward minimizers shown in Corollary 4.4. Throughout this section, we fix $R > 0$, $0 < v < v_{\max}$, and $0 < \eta \leq \eta_0$ (recall Remark 3.12).

Theorem 4.1. *There are constants $\delta_m, \mathfrak{T}_m, c_m > 0$ depending only on R, v, v_{\max} , and η and a constant $\tau_m(R, v, v_{\max}, \eta, \mathfrak{T}) > 0$ such that if we fix $\mathfrak{T} < \mathfrak{T}_m$ and then $\tau < \tau_m$, the following holds. Let Ω be an inward minimizer of \mathcal{F}_τ on Q_R satisfying $\mathcal{E}(\Omega) \leq \mathcal{E}_{\min} + \delta_m$. Then we have*

$$DO(\Omega) \geq c_m$$

where $DO(\Omega)$ is defined in (2.14). If M/G is compact, all constants may be taken independent of R .

It is worth noting that Theorem 4.1 and Lemma 3.4 show that the torsion function w_Ω grows at least linearly away from the boundary, but we cannot immediately deduce an analogous lower bound for the growth of the eigenfunction u_Ω . It will require some delicate Green's function estimates to eventually show in Section 7 that, for minimizers, the first eigenfunction also satisfies this type of linear lower bound.

The core arguments in the proof of Theorem 4.1 below are from David and Toro [18]. We present the details to show how to apply the key estimate of Proposition 3.9 in order to handle the nonlinear term \mathfrak{h} , as well as to verify the dependence on all of the parameters. The proof will be established by iteratively applying the following lemma.

Lemma 4.2. *There exist constants $\delta_m, \mathfrak{T}_m, c_m > 0$ depending only on R, v, v_{\max} , and η and a constant $\tau_m(R, v, v_{\max}, \eta, \mathfrak{T}) > 0$ such that if we fix $\mathfrak{T} < \mathfrak{T}_m$ and then $\tau < \tau_m$, the following holds. Let Ω be an inward minimizer on Q_R satisfying $\mathcal{E}(\Omega) \leq \mathcal{E}_{\min} + \delta_m$. For any $x \in Q_R$, and $\sqrt{r} \leq c_m$, if*

$$(4.1) \quad \sup_{B_r(x)} u_\Omega + \sqrt{\mathfrak{T}}w_\Omega \leq c_m r$$

then

$$\sup_{B_{r/2}(x)} u_\Omega + \sqrt{\mathfrak{T}}w_\Omega \leq \frac{1}{4}c_m r.$$

If M/G is compact, all constants may be taken independent of R .

Proof. The basic idea of the proof is make use of the inward minimality property of Ω using an energy competitor Ω' obtained by removing a small ball from Ω .

Step 1: The competitor Ω' . Let $c_m \leq \min\{\text{inj}_M, 1\}$ be a fixed number to be specified later in the proof. For $r \leq c_m^2$ and $x \in Q_R$, we consider the competitor $\Omega' = \Omega \setminus B_{3r/4}(x)$. Choose δ_m , and \mathfrak{T}_m according to Lemma 3.14, so that Ω has a spectral gap $\lambda_2(\Omega) > \lambda_1(\Omega) + c(R, v_{\max}, v)$ of a definite size. This allows us to apply the key estimate of Proposition 3.9 to $\Omega' \subseteq \Omega$, which yields

$$\|u_{\Omega'} - u_\Omega\|_{L^1} \leq C[\text{tor}(\Omega') - \text{tor}(\Omega) + \lambda_1(\Omega') - \lambda_1(\Omega)],$$

where the constant depends only on R, v_{\max}, v . (Note that the constant a priori also depends on an upper bound for $\lambda_1(\Omega')$; a basic cutoff function argument like the one in step 3 below and the assumption on $\mathcal{E}(\Omega)$

show that $\lambda_1(\Omega') \leq C(R, v_{\max}, v)$. If M/G is compact, we may take the constants here independent of R by using Lemma 3.16 in place of Lemma 3.14. This leads to

$$|\mathfrak{h}(\Omega) - \mathfrak{h}(\Omega')| \leq |\Omega \Delta \Omega'| + \int |u_{\Omega'} - u_{\Omega}| \leq |\Omega \setminus \Omega'| + C[\text{tor}(\Omega') - \text{tor}(\Omega) + \lambda_1(\Omega') - \lambda_1(\Omega)]$$

using property (N3) of \mathfrak{h} . From the inward minimality of Ω , we have $\mathcal{F}_{\tau}(\Omega) \leq \mathcal{F}_{\tau}(\Omega')$ and so

$$\begin{aligned} 0 &\leq \mathcal{F}_{\tau}(\Omega') - \mathcal{F}_{\tau}(\Omega) \\ &\leq \lambda_1(\Omega') - \lambda_1(\Omega) + \mathfrak{T}[\text{tor}(\Omega') - \text{tor}(\Omega)] + f_{v,\eta}(\Omega') - f_{v,\eta}(\Omega) \\ &\quad + \tau \left(|\Omega \setminus \Omega'| + C[\text{tor}(\Omega') - \text{tor}(\Omega) + \lambda_1(\Omega') - \lambda_1(\Omega)] \right) \\ &\leq 2[\mathfrak{T}(\text{tor}(\Omega') - \text{tor}(\Omega)) + \lambda_1(\Omega') - \lambda_1(\Omega)] + (\tau - \eta)|\Omega \setminus \Omega'|, \end{aligned}$$

where the final inequality holds provided τ_m is chosen small enough in terms of C and \mathfrak{T} . So long as $\tau < \eta/2$, we may rewrite this as

$$(4.2) \quad |\Omega \cap B_{3r/4}(x)| = |\Omega \setminus \Omega'| \leq C[\mathfrak{T}(\text{tor}(\Omega') - \text{tor}(\Omega)) + \lambda_1(\Omega') - \lambda_1(\Omega)].$$

Step 2: Estimates for the eigenfunction and torsion function. From Lemma 3.1 as well as the bounds from Lemmas 3.2 and 3.3, we have that $\Delta u_{\Omega}, \Delta w_{\Omega} \geq -C$ on M . Fix ϕ a cutoff function which is 1 on $B_{\frac{7r}{8}}(x)$, compactly supported on B_1 , and has $|\nabla \phi| \leq C/r$. Applying the Caccioppoli inequality to w_{Ω} (i.e. use $\phi^2 w_{\Omega}$ as a test function for w_{Ω} and rearrange terms) and the assumption (4.1),

$$\mathfrak{T} \int \phi^2 |\nabla w_{\Omega}|^2 \leq \mathfrak{T} C \int (|\nabla \phi|^2 w_{\Omega}^2 + \phi^2 w_{\Omega}) \leq C[r^{n-2} c_m^2 r^2 + r^n \sqrt{\mathfrak{T}} c_m r] \leq C r^n c_m^2.$$

In the final inequality we have used that $r \leq \sqrt{r} \leq c_m$. We may estimate u_{Ω} in a similar manner:

$$\int \eta^2 |\nabla u_{\Omega}|^2 \leq C \int (|\nabla \eta|^2 u_{\Omega}^2 + \eta^2 u_{\Omega}) \leq C r^n c_m^2.$$

All together, this implies

$$\int_{B_{7r/8}(x)} |\nabla u_{\Omega}|^2 + \mathfrak{T} |\nabla w_{\Omega}|^2 \leq C r^n c_m^2, \quad \int_{B_{7r/8}(x)} u_{\Omega}^2 + \mathfrak{T} w_{\Omega}^2 \leq C r^{n+2} c_m^2.$$

Step 3: Estimates for the eigenvalue and torsional rigidity of Ω' . Now take another cutoff function ϕ_1 which is 1 on $M \setminus B_{7r/8}(x)$, vanishes on $B_{3r/4}(x)$, and has $|\nabla \phi_1| \leq C/r$. We use $\phi_1 u_{\Omega}$ as a competitor for $\lambda_1(\Omega')$ to find $\lambda_1(\Omega') \leq \int |\nabla(\phi_1 u_{\Omega})|^2 / \int (\phi_1 u_{\Omega})^2$. Estimating the denominator,

$$\int \phi_1^2 u_{\Omega}^2 \geq \int u_{\Omega}^2 - \int_{\Omega \setminus B_{7r/8}(x)} u_{\Omega}^2 \geq 1 - C r^n c_m^2.$$

For the numerator,

$$\int |\nabla(\phi_1 u_{\Omega})|^2 \leq \int |\nabla u_{\Omega}|^2 + \int |\nabla u_{\Omega}|^2 \phi^2 + |\nabla \phi|^2 u_{\Omega}^2 \leq \lambda_1(\Omega) + C r^n c_m^2,$$

so

$$\lambda_1(\Omega') \leq \lambda_1(\Omega) + C r^n c_m^2.$$

An identical estimate on w_{Ω} gives

$$\mathfrak{T} \text{tor}(\Omega') \leq \mathfrak{T} \text{tor}(\Omega) + C r^n c_m^2.$$

Step 4: Conclusion. Substituting the estimates of Step 3 into (4.2), we arrive at

$$(4.3) \quad |\Omega \cap B_{3r/4}(x)| \leq C r^n c_m^2.$$

We now apply the local maximum principle for subsolutions, [24, Theorem 8.17], to $u_{\Omega} + \sqrt{\mathfrak{T}} w_{\Omega}$:

$$\begin{aligned} \sup_{B_{r/2}(x)} u_{\Omega} + \sqrt{\mathfrak{T}} w_{\Omega} &\leq C \left[r^{-n/2} \|u_{\Omega} + \sqrt{\mathfrak{T}} w_{\Omega}\|_{L^2(B_{3r/4}(x))} + r^2 \right] \\ &\leq C \left[c_m r r^{-n/2} |\Omega \cap B_{3r/4}(x)|^{1/2} + r c_m^2 \right] \leq C r c_m^2, \end{aligned}$$

where we used (4.3) and that $r \leq c_m^2$ by assumption. Finally, we choose c_m small enough so that $Cc_m \leq \frac{1}{4}$ to conclude the proof. \square

Theorem 4.1 now follows by iteratively applying Lemma 4.2:

Proof of Theorem 4.1. Set all constants as in Lemma 4.2. Take any $y \in \bar{\Omega}$ and any $r < r_0 := c_m^2/16$. We claim that

$$\sup_{B_r(y)} u_\Omega + \sqrt{\mathfrak{T}} w_\Omega \geq c_m r.$$

Indeed, if this is not the case, we apply Lemma 4.2 to obtain

$$\sup_{B_{r/2}(y)} u_\Omega + \sqrt{\mathfrak{T}} w_\Omega \leq \frac{c_m r}{2}.$$

Then take any $x \in B_{r/4}(y)$:

$$\sup_{B_{r/4}(x)} u_\Omega + \sqrt{\mathfrak{T}} w_\Omega \leq c_m \frac{r}{4},$$

so repeatedly applying Lemma 4.2 gives

$$\sup_{B_{r/4 \cdot 2^{-k}}(x)} u_\Omega + \sqrt{\mathfrak{T}} w_\Omega \leq \frac{c_m r}{4} 2^{-k}.$$

In particular, x is a Lebesgue point of u_Ω , w_Ω and $u_\Omega(x) = w_\Omega(x) = 0$. This is true of all points in $B_{r/4}(y)$, so $u_\Omega \equiv w_\Omega \equiv 0$ there. Using $\Omega \setminus B_{r/4}(x)$ as a competitor for Ω implies that $|\Omega \cap B_{r/4}| = 0$, contradicting that $y \in \bar{\Omega}$ and Ω open.

For $r \in [r_0, 1]$, we simply use $\sup_{B_r(y)} u_\Omega + \sqrt{\mathfrak{T}} w_\Omega \geq c_m r_0 \geq c_m r_0 r$. This gives $\text{DO}(\Omega) \geq c_m r_0$. \square

Note that while for an absolute minimizer (should one exist) $\mathcal{E}(\Omega) \leq \mathcal{E}_{\min} + \delta$ follows from τ being small (see Section 6), there is no reason for that to be the case for an inward minimizer in general.

Remark 4.3. The careful reader may observe that it is possible to modify the argument in this section to give $\sup_{B_r(y)} w_\Omega \geq c_m r$ as long as $\tau < \tau_m(R, v, v_{\max})$ independent of \mathfrak{T} (by using Lemma 3.4 to estimate $u_\Omega \leq Cw_\Omega$). Estimates in this spirit were used by Bucur [11] for other problems. We do not pursue this point further here because the opposite estimate on $\text{UP}(\Omega)$, discussed in the next section, cannot be altered in this way without a version of Proposition 3.9 which omits the torsional rigidity terms on the right-hand side.

The following corollary will help us bound the diameter of minimizers when taking $R \rightarrow \infty$. We will discuss much sharper estimates in Section 7.1, using the upper bound $\text{UP}(\Omega)$ as well.

Corollary 4.4. *Let Ω be as in Theorem 4.1. Then for any $x \in \Omega$ and $r < c_m^4$,*

$$|\Omega \cap B_r(x)| \geq c(r, v, v_{\max}, \eta, R).$$

If M/G is compact, c may be taken independent of R .

Proof. As in the proof of Lemma 4.2, apply the local maximum principle on $B_r(x)$ to obtain

$$\begin{aligned} \sup_{B_{r/2}(x)} u_\Omega^2 + \mathfrak{T} w_\Omega^2 &\leq Cr^{-n} \int_{B_r(x) \cap \Omega} (u_\Omega^2 + \mathfrak{T} w_\Omega^2) + Cr^4 \leq C|\Omega \cap B_r(x)|r^{-n} \sup_{\Omega} (u_\Omega^2 + \mathfrak{T} w_\Omega^2) + Cr^4 \\ &\leq C|\Omega \cap B_r(x)|r^{-n} + Cr^4, \end{aligned}$$

with the last step using Lemmas 3.2 and 3.3. The quantities $|\text{tor}(\Omega)|$ and $\lambda_1(\Omega)$ are controlled by $\mathcal{E}(\Omega)$, which is bounded from the assumption $\mathcal{E}_{\min} \leq \mathcal{E}(\Omega) \leq \mathcal{E}_{\min} + \delta$ and Lemma 3.10 or 3.15. The constant C can be taken to depend only on M, v, v_{\max}, η if M/G compact, and otherwise also depends on R . Then applying Theorem 4.1 to the left-hand side, $c_m^2 r^2 / 2 \leq C|\Omega \cap B_r(x)|r^{-n} + Cr^4$, so reabsorbing the r^4 implies the conclusion. \square

5. THE UPPER BOUND

In this section we consider outward minimizers of the main functional \mathcal{F}_τ defined in (2.8). The main result of the section is Theorem 5.1 below, which shows that the function $u_\Omega + \sqrt{\mathfrak{T}}w_\Omega$ grows at most linearly away from the boundary of an outer minimizer Ω . The result is phrased in terms of the quantity $UP(\Omega)$ defined in (2.13). Throughout this section, we fix $R > 0$, $0 < v < v_{\max}$, and $0 < \eta \leq \eta_0$ (recall Remark 3.12).

Theorem 5.1. *There are constants $\delta_M, \mathfrak{T}_M, C_M > 0$ depending only on R, v, v_{\max} , and η and a constant $\tau_M(\mathfrak{T}, R, v, v_{\max}, \eta) > 0$ such that if we fix $\mathfrak{T} < \mathfrak{T}_M$ and then $\tau < \tau_M$, the following holds. Let Ω be an outward minimizer on Q_R and suppose $\mathcal{E}(\Omega) \leq \mathcal{E}_{\min} + \delta_M$. Then we have*

$$UP(\Omega) \leq C_M$$

where $UP(\Omega)$ is defined in (2.13).

The proof of Theorem 5.1 will require a few initial lemmas. We start by checking that, for an outward minimizer Ω , u_Ω and w_Ω are close to their harmonic replacements in small balls. By the harmonic replacement of u_Ω on $B_r(x) \cap Q_R$, we mean the unique function u with $u - u_\Omega \in H_0^1(B_r(x) \cap Q_R)$ and $\Delta u = 0$ on $B_r(x) \cap Q_R$ in the weak sense, extended so that $u = u_\Omega$ on $Q_R \setminus B_r(x)$.

Lemma 5.2. *There are constants $\delta_M, \mathfrak{T}_M, C_M, r_0 > 0$ depending only on R, v, v_{\max} , and η and a constant $\tau_M(\mathfrak{T}, R, v, v_{\max}) > 0$ such that if we fix $\mathfrak{T} < \mathfrak{T}_M$ and then $\tau < \tau_M$, the following holds. Let Ω be an outward minimizer on Q_R satisfying $\mathcal{E}(\Omega) \leq \mathcal{E}_{\min} + \delta_M$. Then for any $x \in Q_R$ and $r < r_0$,*

$$\int |\nabla(u_\Omega - u)|^2 + \mathfrak{T}|\nabla(w_\Omega - w)|^2 \leq C_M \left[|B_r(x) \setminus \Omega| + r^n \sup_{B_r(x)} (u_\Omega^2 + \mathfrak{T}w_\Omega^2) \right] \leq C_M r^n,$$

where u, w are the harmonic replacements of u_Ω, w_Ω on $B_r(x) \cap Q_R$.

Proof. The basic idea of the proof is make use of the inward minimality property of Ω using an energy competitor Ω' obtained by adding a small ball to Ω .

Step 1: The competitor Ω' . Set $\Omega' = \Omega \cup (B_r(x) \cap Q_R)$ and $A = \sup_{B_r(x)} u_\Omega^2 + \mathfrak{T}w_\Omega^2$. First, we may use that $\lambda_1(\Omega') \leq \lambda_1(\Omega)$ and $\text{tor}(\Omega') \leq \text{tor}(\Omega)$ to deduce that

$$\mathcal{E}(\Omega') \leq \mathcal{E}(\Omega) + \frac{C}{\eta} r^n \leq \delta_M + \frac{C}{\eta} r_0^n \leq 2\delta_M$$

if r_0 is small enough. Choosing δ_M and \mathfrak{T}_M sufficiently small and applying Lemmas 3.11 and 3.14 shows that $\Omega' \in \mathcal{H}$ and $\lambda_2(\Omega') > \lambda_1(\Omega') + c(v, v_{\max}, R)$. The latter fact allows us to apply the key Proposition 3.9 to Ω, Ω' to find that

$$\int |u_\Omega - u_{\Omega'}| \leq C[\lambda_1(\Omega) - \lambda_1(\Omega') + \text{tor}(\Omega) - \text{tor}(\Omega')],$$

where C depends only on R and v, v_{\max} . So, from assumption (N3) on \mathfrak{h} , we know that

$$|\mathfrak{h}(\Omega) - \mathfrak{h}(\Omega')| \leq |\Omega \Delta \Omega'| + \int |u_\Omega - u_{\Omega'}| \leq |B_r(x) \setminus \Omega| + C[\lambda_1(\Omega) - \lambda_1(\Omega') + \text{tor}(\Omega) - \text{tor}(\Omega')].$$

Now, we use Ω' as an energy competitor for Ω' ; the outward minimizing property gives that

$$\begin{aligned} 0 \leq \mathcal{F}_\tau(\Omega') - \mathcal{F}_\tau(\Omega) &\leq \lambda_1(\Omega') - \lambda_1(\Omega) + \mathfrak{T}(\text{tor}(\Omega') - \text{tor}(\Omega)) + \frac{C}{\eta} |\Omega' \setminus \Omega| + \tau |\mathfrak{h}(\Omega) - \mathfrak{h}(\Omega')| \\ &\leq [\lambda_1(\Omega') - \lambda_1(\Omega) + \mathfrak{T}(\text{tor}(\Omega') - \text{tor}(\Omega))] \left(1 - C \frac{\tau}{\mathfrak{T}}\right) + \frac{C}{\eta} |B_r(x) \setminus \Omega|. \end{aligned}$$

Select τ_M small enough that $1 - C \frac{\tau_M}{\mathfrak{T}} < \frac{1}{2}$; then

$$(5.1) \quad \lambda_1(\Omega) - \lambda_1(\Omega') + \mathfrak{T}(\text{tor}(\Omega) - \text{tor}(\Omega')) \leq C |B_r(x) \setminus \Omega|.$$

Step 2: Estimates for the eigenvalue and torsional rigidity of Ω' . Next, we use u as a competitor for $\lambda_1(\Omega')$ to find $\lambda_1(\Omega') \leq \int |\nabla u|^2 / \int u^2$. To estimate this quantity further, from the maximum principle we know that $u^2 + \mathfrak{T}w^2 \leq CT$ on $B_r(x)$ as well. This can be used to bound the denominator:

$$\int u^2 = \int u_\Omega^2 + \int_{B_r(x) \cap Q_R} u^2 - u_\Omega^2 \geq 1 - CA|B_r(x)| \geq 1 - CAr^n.$$

For the energy term,

$$\begin{aligned} \int |\nabla u|^2 &= \lambda_1(\Omega) + \int_{B_r(x) \cap Q_R} |\nabla u|^2 - |\nabla u_\Omega|^2 \\ &= \lambda_1(\Omega) + \int_{B_r(x) \cap Q_R} g(\nabla(u - u_\Omega), \nabla(u + u_\Omega)) \\ &= \lambda_1(\Omega) - \int_{B_r(x) \cap Q_R} g(\nabla(u - u_\Omega), \nabla(u - u_\Omega)) = \lambda_1(\Omega) - \int_{B_r(x) \cap Q_R} |\nabla(u - u_\Omega)|^2, \end{aligned}$$

where the second-to-last step used that $\int g(\nabla u, \nabla(u - u_\Omega)) = 0$ since $\Delta u = 0$ and $u - u_\Omega \in H_0^1(B_r(x) \cap Q_R)$. Putting these together leads to

$$\lambda_1(\Omega') \leq \lambda_1(\Omega) - \int |\nabla(u - u_\Omega)|^2 + CAr^n.$$

A similar computation (using Lemma 3.3) gives

$$\mathfrak{T}\text{tor}(\Omega') \leq \mathfrak{T}\text{tor}(\Omega) - \frac{\mathfrak{T}}{2} \int |\nabla(w - w_\Omega)|^2 + CAr^n.$$

Step 3: Conclusion. Plugging these into (5.1),

$$\int |\nabla(u - u_\Omega)|^2 + \mathfrak{T} \int |\nabla(w - w_\Omega)|^2 \leq C[Ar^n + C|B_r(x) \setminus \Omega|].$$

Recalling from Lemmas 3.2 and 3.3 that $A \leq C(R, v_{\max})$ and $|B_r(x) \setminus \Omega| \leq |B_r(x)| \leq Cr^n$ concludes the proof. \square

Lemma 5.2 immediately gives that u_Ω, w_Ω are Hölder continuous functions from [15], and also satisfy a Morrey-type estimate:

Corollary 5.3. *Let Ω be as in Lemma 5.2. Then $u_\Omega, \mathfrak{T}w_\Omega \in C^{0,\alpha}(Q_R)$ for any $\alpha < 1$, with*

$$(5.2) \quad [u_\Omega]_{C^{0,\alpha}} + \sqrt{\mathfrak{T}}[w_\Omega]_{C^{0,\alpha}} \leq C \sup_{x \in Q_R, r < r_0} r^{-n/2+1-\alpha} [\|\nabla u_\Omega\|_{L^2(B_r(x))} + \mathfrak{T}\|\nabla w_\Omega\|_{L^2(B_r(x))}] \leq C$$

for $C = C(R, v, v_{\max}, \eta, \alpha)$.

Below we show how to instead obtain the more precise estimate $u_\Omega, w_\Omega \in C^{0,1}$, which will follow easily from the growth estimate of Theorem 5.1. Before proving Theorem 5.1, let us recall some basic facts about Green's functions and prove a mean value-type inequality. Consider the (positive) Green's function $G(x, \cdot)$ for $B_r(x)$ (with pole at the center of the ball). For any continuous function ϕ which is smooth on a neighborhood of x and has $\Delta\phi$ represented by a finite Borel measure, we have that

$$\phi(x) + \int_{B_r(x)} G(x, y) d\Delta\phi(y) = \int_{\partial B_r(x)} g(\nabla G(x, y), \nu_y) \phi(y) d\mathcal{H}^{n-1}(y)$$

where ν_y is the outward unit normal to $\partial B_r(x)$. From, say, [26, Theorem 1.2.8], we have the following standard bounds on G :

$$(5.3) \quad -Cr^{1-n} \leq g(\nabla G(x, y), \nu_y) \leq -cr^{1-n} \quad y \in \partial B_r(x)$$

and

$$(5.4) \quad c(|y - x|^{2-n} - r^{2-n}) \leq G(x, y) \leq C(|y - x|^{2-n} - r^{2-n}) \quad y \in B_r(x) \setminus \{x\}$$

if $n \geq 3$, while

$$c(\log|y - x| - \log r) \leq G(x, y) \leq C(\log|y - x| - \log r)$$

instead when $n = 2$, with constants depending only on R (i.e. they are uniform in x and r).

Lemma 5.4 (Mean value-type inequality). *Let ϕ be a nonnegative continuous function that is either (1) smooth in a neighborhood of x or (2) satisfies the Morrey-type estimate (5.2) and has $\phi(x) = 0$. Then*

$$(5.5) \quad \frac{1}{r^{n-1}} \int_{\partial B_r(x)} \phi \leq C \left[\phi(x) + \int_0^r s^{1-n} |\Delta \phi|(B_s(x)) ds \right].$$

Proof. Let us first assume that $\phi \geq 0$ is smooth in a neighborhood of x . When $n \geq 3$, making use of the upper bound in (5.4), we have

$$\begin{aligned} \int_{B_r(x)} G(x, y) d\Delta \phi(y) &\leq \int_{B_r(x)} G(x, y) d|\Delta \phi|(y) \\ &\leq C \int_0^r \int_{\partial B_s(x)} [s^{2-n} - r^{2-n}] d|\Delta \phi| ds \\ &= C \int_0^r [s^{2-n} - r^{2-n}] \int_{\partial B_s(x)} d|\Delta \phi| ds = C(n-2) \int_0^r s^{1-n} \int_{B_s(x)} d|\Delta \phi| ds. \end{aligned}$$

The last step was an integration by parts in one variable, using that $s^{2-n} \int_{B_s(x)} |\Delta \phi| \rightarrow 0$ as $s \rightarrow 0$ (recall that $\Delta \phi$ is smooth near 0). From (5.3), we have

$$\frac{1}{r^{n-1}} \int_{\partial B_r(x)} \phi - C\phi(x) \leq C \left[-\phi(x) - \int_{\partial B_r(x)} g_y(\nabla G(x, y), \nu_y) \phi(y) d\mathcal{H}^{n-1}(y) \right] = C \int_{B_r(x)} G(x, y) d\Delta \phi(y).$$

Together these two estimates give (5.5) for $n \geq 3$. A similar computation gives an identical estimate when $n = 2$.

Now, suppose that instead of ϕ being smooth in a neighborhood of x , we simply know that ϕ satisfies (5.2) and $\phi(x) = 0$. For fixed $t \in (0, r/2)$, let η_t be a cutoff function which vanishes on $B_t(x)$, is 1 outside of $B_{2t}(x)$, and has $|\nabla \eta_t| \leq Ct^{-1}$, $|D^2 \eta_t| \leq Ct^{-2}$. In this way, $\phi \eta_t$ is smooth in a neighborhood of x and we may apply (5.5) to $\phi \eta_t$.

$$\Delta(\phi \eta_t) = \eta_t \Delta \phi + 2g(\nabla \phi, \nabla \eta_t) + \phi \Delta \eta_t$$

as distributions (with the second two terms absolutely continuous), so

$$\begin{aligned} |\Delta(\phi \eta_t)|(B_s(x)) &\leq |\Delta \phi|(B_s(x) \setminus B_t(x)) + \int_{B_{2t}(x)} 2g(\nabla \phi, \nabla \eta_t) + \phi \Delta \eta_t \\ &\leq |\Delta \phi|(B_s(x)) + Ct^{n/2+\alpha-1} t^{n/2-1} + Ct^\alpha t^{n-2} \\ &\leq |\Delta \phi|(B_s(x)) + Ct^{n-2+\alpha} \end{aligned}$$

for $s \geq t$, and 0 otherwise. Then (5.5) applies to $\eta_t \phi$, to give

$$\begin{aligned} \frac{1}{r^{n-1}} \int_{\partial B_r(x)} \phi &\leq C \left[\eta_t \phi(x) + \int_t^r s^{1-n} |\Delta \phi|(B_s(x)) + Cs^{1-n} t^{n-2+\alpha} ds \right] \\ &\leq C \left[0 + \int_t^r s^{1-n} |\Delta \phi|(B_s(x)) ds + Ct^\alpha \right]. \end{aligned}$$

Sending $t \rightarrow 0$ completes the proof. \square

We are now ready to prove Theorem 5.1.

Proof of Theorem 5.1. Step 1: Setup. Fix a point $x \in \partial \Omega$ and an $r < r_0$, with $r_0 \leq \text{inj}_M/4$ to be chosen below. We will show that

$$S = S(x, r) := \sup\{u(y) : y \in \partial B_r(x), d(y, \partial \Omega) = r\} \leq C_M r,$$

where u is either u_Ω or $\sqrt{\Sigma} w_\Omega$. As we know that u_Ω, w_Ω are bounded, the same inequality is automatic for $r \in [r_0, 1]$ with constant C/r_0 . This will imply the conclusion, as for every $y \in \Omega$ with $d(y, \partial \Omega) \leq 1$, there is an $x \in \partial \Omega$ and r with $r = d(y, \partial \Omega) = d(y, x)$, so $u \leq S(x, r)$.

If $B_{2r}(x)$ is not fully contained in Q_R , we have that $u_\Omega \leq Cw_\Omega \leq Cw_{Q_R}$ on Ω (using Lemma 3.4), and from elliptic regularity and the smoothness of Q_R , $w_{Q_R}(y) \leq C(R)d(y, \partial Q_R)$. This gives

$$S \leq C \sup_{\partial B_r(x)} w_{Q_R} \leq Cr,$$

and we are done. Assume, then, that $B_{2r}(x) \subseteq Q_R$.

Step 2: Initial S bound. In this step, we bound S from above in terms of the integral of u over spheres $\partial B_s(x)$ for $s \in (0, 1)$ using a multiple of the Green's function as a barrier. More specifically, let $z \in \partial B_r(x) \cap \{y : d(y, \partial\Omega) = r\}$ be a point such that $u(z) = S$. By definition $B_r(z) \subseteq \Omega$, and so from Lemma 3.1 we have $|\Delta u| \leq C$ on this ball. Applying the Harnack inequality [24, Theorems 8.17, 8.18],

$$S \leq \sup_{B_{r/2}(z)} u \leq C \inf_{B_{r/2}(z)} u + Cr^2.$$

If $S \leq r$, we are done; if not, choose r_0 small enough in terms of C so that $Cr^2 \leq \frac{1}{2}r \leq \frac{1}{2}S$, and thus

$$(5.6) \quad \inf_{B_{r/2}(z)} u \geq \frac{S}{C}.$$

In order to propagate this bound to balls centered at x (at least in an integral sense), consider the barrier $v(y) = \frac{S}{C_*} r^{n-2} G(z, y)$, where $G(z, y)$ is the Green's function for $B_r(z)$. On $\partial B_r(z)$, we have $v = 0 \leq u$, while $v \leq u$ on $\partial B_{r/2}(z)$ so long as C_* is chosen large relative to the constants in (5.6) and (5.4). On the annular region between them, $\Delta v = 0$ and $\Delta u \leq 0$. It follows from the comparison principle that $v \leq u$ on $B_r(z) \setminus B_{r/2}(z)$. In particular, in conjunction with (5.6), this means that for any $t \in (\frac{r}{2}, r)$ we have

$$\inf_{B_t(z)} u \geq cS \left(\frac{r^{n-2}}{t^{n-2}} - 1 \right) \geq cS \frac{r-t}{r}.$$

For any $s \in (0, r)$, choose $t \in (r/2, r)$ so that $r-t = s/2$. One may check that the set $B_t(z) \cap \partial B_s(x)$ has measure of this set is bounded from below by cs^{n-1} . So, integrating over this set and applying the previous estimate, we have

$$(5.7) \quad \int_{\partial B_s(x)} u \geq \int_{\partial B_s(x) \cap B_t(z)} u \geq cs^{n-1} \frac{s}{r} S$$

for all $s \in (0, r)$. This is the first of two estimates which will be combined to bound S .

Ultimately, in Step 4, we will use the mean value-type inequality of Lemma 5.4 in order to bound the left-hand side of (5.7). So, it is essential to estimate the $|\Delta|$ measure of balls $B_s(x)$.

Step 3: Estimate for $|\Delta|$. The second estimate involves computing the total variation of Δu when viewed as a measure, using Lemma 5.2. From Lemma 3.1, we know that $\Delta u = \mu - f$ in the sense of distributions, where μ is a nonnegative Borel measure supported on $\partial\Omega$ and f is either $\lambda_1(\Omega)u_\Omega$ or $\mathfrak{I}1_\Omega$, and in each case is uniformly bounded. Fix $s \leq r$ and let h be the harmonic replacement of u on $B_s(x)$: from Lemma 5.2,

$$\int_{B_s(x)} |\nabla(u-h)|^2 \leq Cs^n.$$

Using that $\Delta h = 0$ on $B_s(x)$ and $u = h$ on ∂B_s ,

$$\int_{B_s(x)} (h-u)d\Delta u = - \int_{B_s(x)} g(\nabla u, \nabla(h-u)) = \int |\nabla(u-h)|^2 \leq Cs^n.$$

The leftmost identity used that Δu represents the Laplacian of u in the sense of distributions. Decomposing this further,

$$\int_{B_s(x)} (h-u)d\Delta u = \int_{B_s(x) \cap \partial\Omega} (h-u)d\mu - \int_{B_s(x) \cap \Omega} (h-u)f.$$

The second term is controlled by Cs^n , as both h, u are bounded. In the first term, $u = 0$, so

$$(5.8) \quad \int_{B_s(x) \cap \partial\Omega} h d\mu \leq Cs^n.$$

We thus need to bound h from below on this set. Applying the Harnack inequality followed by (5.5) to h gives that

$$C \inf_{B_{s/2}(x)} h \geq h(x) \geq cs^{1-n} \int_{\partial B_s(x)} h.$$

Since $h = u$ on $\partial B_s(x)$, we find from (5.7) that $\inf_{B_{s/2}(x)} h \geq c \frac{s}{r} S$. Combining this with (5.8), we get

$$c \frac{s}{r} S \mu(B_{s/2}) \leq \mu(B_{s/2}) \inf_{B_{s/2}(x)} h \leq \int_{B_s(x) \cap \partial \Omega} h d\mu \leq Cs^n.$$

So, we arrive at our second main estimate:

$$(5.9) \quad |\Delta u|(B_{s/2}) \leq \frac{Cr}{S} s^{n-1} + Cs^n \leq \frac{Cr}{S} s^{n-1},$$

where the very last step used that $s \leq r$ and $S \leq C$ is bounded.

Step 4: Conclusion. By Corollary 5.3, and the fact that $u(x) = 0$, we may apply Lemma 5.4 to u . So, using (5.9), Lemma 5.4 tells us that

$$\frac{1}{r^{n-1}} \int_{\partial B_{r/4}(x)} u \leq C \int_0^{r/4} \frac{r}{S} ds = C \frac{r^2}{S}.$$

Applying (5.7) to the left-hand side gives $S^2 \leq Cr^2$. This bounds S , completing the proof. \square

By applying standard elliptic estimates on $B_d(x)$ for any $x \in \Omega$, $d = \min\{r_0, d(x, \partial \Omega)/2\}$ and using that $u_\Omega + \sqrt{\mathfrak{I}} w_\Omega \leq CdUP(\Omega)$ on this ball, we obtain the following gradient estimate as well:

Corollary 5.5. *Under the assumptions of Theorem 5.1,*

$$|\nabla u_\Omega| + \sqrt{\mathfrak{I}} |\nabla w_\Omega| \leq C(UP(\Omega), R, v, v_{max}).$$

6. EXISTENCE OF MINIMIZERS

We are now in a position to show that minimizers to the main functional \mathcal{F}_τ defined in (2.8) exist. Theorem 6.2 below establishes the existence of minimizers of \mathcal{F}_τ among sets in $\mathcal{H}_{R, v_{max}}$ for any Riemannian manifold. In the case when M/G_0 is compact for a subgroup G_0 of the isometry group for which the functional is invariant, Theorem 6.3 shows that the parameter R may be taken to ∞ and global minimizers exist. Throughout the section, we fix $R > 0$ and $0 < v < v_{max}$, and assume that $\eta < \eta_0(R, v, v_{max})$ and $\mathfrak{I} < \mathfrak{I}_0(R, v, v_{max}, \eta)$ are fixed to be small enough that all results in Sections 3, 4, and 5 apply.

In the following lemma, we show that any set Ω with low base energy is contained in an outward minimizer whose energy does not exceed the energy of Ω .

Lemma 6.1. *There are $\delta = \delta(R, v, v_{max}, \eta) > 0$ and $\tau_0 = \tau_0(R, v, v_{max}, \mathfrak{I}, \eta)$ such that if $\tau < \tau_0$ and $\Omega \in \mathcal{H}$ with $\mathcal{E}(\Omega) \leq \mathcal{E}_{min} + \delta$, then there exists a $U \in \mathcal{H}$ with $\Omega \subseteq U$, $\mathcal{F}_\tau(U) \leq \mathcal{F}_\tau(\Omega)$, and U an outward minimizer.*

Proof. Given any bounded open set E , let

$$m_E = \inf\{\mathcal{F}_\tau(V) : V \in \mathcal{H}, E \subseteq V\}.$$

Clearly $m_\Omega \leq \mathcal{F}_\tau(\Omega)$, and as $m_\Omega \geq \mathcal{E}_{min} > -\infty$ from Lemma 3.10 this is bounded from below. If $m_\Omega = \mathcal{F}_\tau(\Omega)$, then Ω is itself an outward minimizer and we are done by letting $U = \Omega$. If not, we construct the set U in the following way. Let $U_1 \in \mathcal{H}$ be a set with $\Omega \subseteq U_1$ which has $\mathcal{F}_\tau(U_1) < (m_\Omega + \mathcal{F}_\tau(\Omega))/2$. Now repeat this construction with U_1 in place of Ω , producing a nested sequence (finite or infinite) $U_k \in \mathcal{H}$ with $\Omega \subseteq U_k \subseteq U_{k+1}$, and $\mathcal{F}_\tau(U_k) < (m_{U_{k-1}} + \mathcal{F}_\tau(U_{k-1}))/2$.

The sequence is finite if $m_{U_k} = \mathcal{F}_\tau(U_k)$ for the final k , meaning it is also an outward minimizer. In this case, U_k satisfies the conclusions of this lemma and we are done by taking $U = U_k$. Now assume the sequence is infinite. Observe that $m_{U_k} \geq m_{U_{k-1}}$ from the fact that $U_{k-1} \subseteq U_k$. Set $m = \lim_k m_{U_k}$. On the other hand, $\mathcal{F}_\tau(U_k)$ is decreasing and

$$\mathcal{F}_\tau(U_k) - m_{U_k} < \frac{m_{U_{k-1}} + \mathcal{F}_\tau(U_{k-1})}{2} - m_{U_{k-1}} = \frac{\mathcal{F}_\tau(U_{k-1}) - m_{U_{k-1}}}{2},$$

meaning $m = \lim_k \mathcal{F}_\tau(U_k)$.

Let $U = \cup_k U_k$; this is an open set. Note that $\Omega \subseteq U \subseteq Q_R$ and $|U| = \lim |U_k| \leq v_{\max}$ from the monotone convergence theorem. We also have $m_U \geq m_{U_k}$ for each k , so $m_U \geq m$. Our goal is to show that $\mathcal{F}_\tau(U) = m$; if we do so, then clearly U is an outward minimizer and satisfies the assumptions.

First, $\mathcal{E}(U) \leq \liminf_k \mathcal{E}(U_k)$. Indeed, this follows by extracting a subsequence of first eigenfunctions u_{U_k} and torsion functions w_{U_k} for U_k which converge weakly in $H_0^1(U)$ to some u, w which may be used as competitors for the eigenvalue and torsional rigidity problems on U (together with the previously observed fact that $|U \setminus U_k| \rightarrow 0$):

$$\lambda_1(U) \leq \frac{\int |\nabla u|^2}{\int u^2} \leq \liminf_k \frac{\int |\nabla u_{U_k}|^2}{\int u_{U_k}^2} = \liminf_k \lambda_1(U_k),$$

and similarly $\text{tor}(U) \leq \liminf_k \text{tor}(U_k)$ and $f_{v,\eta}(U) = \lim_k f_{v,\eta}(U_k)$. This implies that

$$(6.1) \quad \mathcal{E}(U) \leq \lim \mathcal{F}_\tau(U_k) \leq \mathcal{F}_\tau(\Omega) \leq \mathcal{E}(\Omega) + \tau_0 \leq \mathcal{E}_{\min} + \delta + \tau_0.$$

Set $\alpha = \liminf_k \mathcal{E}(U_k) - \mathcal{E}(U) \geq 0$, noting that

$$\liminf_k (\lambda_1(U_k) + \text{tor}(U_k) - \lambda_1(U) - \text{tor}(U)) \leq \frac{\alpha}{\mathfrak{T}}.$$

Consider now the remaining term $\mathfrak{h}(U)$. From Proposition 3.9 applied to U and U_k (using Lemma 3.14 and (6.1), choosing δ and τ_0 small enough), we have that

$$\int |u_{U_k} - u_U| \leq C_4(\text{tor}(U_k) - \text{tor}(U) + \lambda_1(U_k) - \lambda_1(U)).$$

Therefore, from property (N3) of \mathfrak{h} , we have that

$$|\mathfrak{h}(U) - \mathfrak{h}(U_k)| \leq |U \setminus U_k| + C_4(\text{tor}(U_k) - \text{tor}(U) + \lambda_1(U_k) - \lambda_1(U)) \leq \frac{C_4\alpha}{\mathfrak{T}} + o_k(1).$$

In particular, if $\tau_0 C_4 < \mathfrak{T}/2$, this leads to

$$\begin{aligned} m \leq m_U \leq \mathcal{F}_\tau(U) &= \mathcal{E}(U) + \tau \mathfrak{h}(U_k) + \tau |\mathfrak{h}(U) - \mathfrak{h}(U_k)| \\ &\leq \mathcal{E}(U) + \tau \mathfrak{h}(U_k) + \tau \frac{C_4\alpha}{\mathfrak{T}} + o_k(1) \\ &= \mathcal{E}(U_k) - \alpha + \tau \mathfrak{h}(U_k) + \frac{\alpha}{2} + o_k(1) \\ &= \mathcal{F}_\tau(U_k) - \alpha + \frac{\alpha}{2} + o_k(1) = m - \frac{\alpha}{2} + o_k(1). \end{aligned}$$

Taking the limit, we see that $\alpha = 0$, $\mathcal{F}_\tau(U) = m$, and the conclusion follows. \square

We now show that a minimizer of \mathcal{F}_τ exists among sets in $\mathcal{H} = \mathcal{H}_{R,v_{\max}}$. The proof makes use of Lemma 6.1 to replace any minimizing sequence with a minimizing sequence of *outer minimizers*, whose eigenfunctions and torsions functions we know to be uniformly Lipschitz from the results of Section 5.

Theorem 6.2. *There is a $\tau_0 = \tau_0(R, v, v_{\max}, \mathfrak{T}, \eta)$ such that if $\tau < \tau_0$, there exists an $\Omega \subseteq \mathcal{H}$ that is a minimizer of \mathcal{F}_τ among sets in $\mathcal{H} = \mathcal{H}_{R,v_{\max}}$.*

Proof. Let $\alpha = \inf\{\mathcal{F}_\tau(E) : E \in \mathcal{H}\}$, and Ω_j a minimizing sequence for \mathcal{F}_τ . Note that $\mathcal{E}(\Omega) \leq \mathcal{F}_\tau(\Omega) \leq \mathcal{E}(\Omega) + \tau$ for any set Ω , so $\alpha \leq \mathcal{E}_{\min} + \tau$ and $\mathcal{E}(\Omega_j) \leq \mathcal{E}_{\min} + 2\tau$ for all j large enough. Apply Lemma 6.1 to replace each Ω_j with an outward minimizer with smaller $\mathcal{F}_\tau(\Omega_j)$, choosing $2\tau_0 < \delta_0$ there. Let $u_{\Omega_j}, w_{\Omega_j}$ be the first eigenfunctions and torsion functions respectively, and pass to subsequences with $u_{\Omega_j} \rightarrow u$ and $w_{\Omega_j} \rightarrow w$ weakly in $H_0^1(Q_R)$ and strongly in $L^2(Q_R)$. Our approach here is different from the previous lemma: we show that u is the unique first eigenfunction on its positivity set: $u = u_{\{u>0\}}$.

Apply Corollary 5.5 to see that the u_{Ω_j} and w_{Ω_j} are equicontinuous, and so converge uniformly to u and w (which are Lipschitz functions) on Q_R . This also implies that $\Omega = \{u > 0\}$ is, by definition, an open set. Let us show that Ω is minimizer.

Let $\lambda = \lim \lambda_1(\Omega_j)$ (pass to a subsequence if needed). From the uniform convergence of $u_{\Omega_j} \rightarrow u$, we have that for every $x \in \Omega$, there is a ball $B_r(x)$ such that $B_r(x) \subseteq \Omega_j$ for every j large enough. On this ball, the u_{Ω_j} converge in smooth topology to u from standard elliptic estimates, and so $-\Delta u(x) = \lambda u(x)$ passes to the limit. In particular, this implies that u_Ω is an eigenfunction for Ω with eigenvalue λ .

From the weak convergence of $u_{\Omega_j} \rightarrow u$ and $w_{\Omega_j} \rightarrow w$ in H_0^1 , we have that

$$\lambda_1(\Omega) + \mathfrak{T}\text{tor}(\Omega) \leq \frac{\int |\nabla u|^2}{\int u^2} + \mathfrak{T} \int \frac{|\nabla w|^2}{2} - w \leq \liminf_j \lambda_1(\Omega_j) + \mathfrak{T}\text{tor}(\Omega_j).$$

From Fatou's lemma, $|\Omega| \leq \liminf_j |\Omega_j|$, so $\mathcal{E}(\Omega) \leq \liminf_j \mathcal{E}(\Omega_j)$. We also have that $\Omega \subseteq Q_R$ and $|\Omega| \leq v_{\max}$, so $\Omega \in \mathcal{H}$ is an admissible competitor.

Using that $\mathcal{E}(\Omega) \leq \liminf \mathcal{E}(\Omega_j) \leq \mathcal{E}_{\min} + 2\tau$, we may apply Lemma 3.14 to deduce that $\lambda_2(\Omega) > \lambda_1(\Omega) + c(R, v, v_{\max})$. As $\mathcal{E}_{\min} \leq \mathcal{E}(\Omega)$ as well, we must have that

$$0 \leq \lim \lambda_1(\Omega_j) - \lambda_1(\Omega) \leq 2\tau_0,$$

and so $|\lambda_1(\Omega) - \lambda| \leq 2\tau_0$. Ensuring that τ_0 is small enough in terms of R, v, v_{\max} guarantees that $\lambda < \lambda_2(\Omega)$. In particular, this means there is a unique eigenfunction on Ω with eigenvalue at most λ , and so $u_\Omega = u$. In particular, this means that $u_{\Omega_j} \rightarrow u_\Omega$ uniformly.

Set $\beta = \liminf_j \mathcal{E}(\Omega_j) - \mathcal{E}(\Omega) \geq 0$, which has $\beta \geq \eta[\liminf_j |\Omega_j| - |\Omega|]$. From Fatou's lemma, $|\Omega| \leq \liminf_j |\Omega_j \cap \Omega|$, so $\liminf_j |\Omega \setminus \Omega_j| = 0$. This gives that

$$\liminf_j |\Omega \Delta \Omega_j| = \liminf_j |\Omega_j \setminus \Omega| = \liminf_j (|\Omega_j| - |\Omega|) \leq \frac{\beta}{\eta}.$$

It follows from assumption (N3) of Definition 2.11 on \mathfrak{h} that $|\mathfrak{h}(\Omega) - \mathfrak{h}(\Omega_j)| \leq o_j(1) + |\Omega \Delta \Omega_j| \leq o_j(1) + \frac{\beta}{\eta}$. We can now estimate the energy of Ω in the following way:

$$\begin{aligned} \alpha \leq \mathcal{F}_\tau(\Omega) &\leq \mathcal{E}(\Omega) + \tau \mathfrak{h}(\Omega_j) + o_j(1) + \tau \frac{\beta}{\eta} \\ &\leq \mathcal{E}(\Omega_j) - \beta + \tau \mathfrak{h}(\Omega_j) + o_j(1) + \tau \frac{\beta}{\eta} \\ &= \mathcal{F}_\tau(\Omega_j) - \beta + o_j(1) + \tau \frac{\beta}{\eta} = \alpha - \beta + o_j(1) + \tau \frac{\beta}{\eta}. \end{aligned}$$

If $\tau < \eta$, this implies that $\beta = 0$, $\mathcal{F}_\tau(\Omega) = \alpha$, and Ω is a minimizer. \square

Next, we prove that when M/G_0 is compact, a minimizer of \mathcal{F}_τ among all open bounded sets with $|\Omega| \leq v_{\max}$ exists. The main idea is to show that the minimizers Ω_R from the previous theorem have uniformly bounded diameter, and thus Ω_R is a global minimizer for R sufficiently large.

Theorem 6.3. *Assume M/G_0 is compact. Then exist constants $S = S(v, v_{\max}, \mathfrak{T}, \eta) > 0$ and $\tau_0 = \tau_0(v, v_{\max}, \mathfrak{T}, \eta)$ such that if $\tau < \tau_0$, there exists an open, bounded Ω with $|\Omega| \leq v_{\max}$ which minimizes \mathcal{F}_τ over all such sets. For every such minimizer Ω of \mathcal{F}_τ , there is an $e \in G_0$ and an such that $e(\Omega) \subseteq Q_S$.*

Recall that G_0 is a (possibly empty) subgroup of isometries under which \mathfrak{h} , and hence \mathcal{F}_τ , is invariant. The assumption here implies that M/G is compact.

Proof. For each $R \gg 1$, consider the minimizer Ω_R obtained from Theorem 6.2 of \mathcal{F}_τ over \mathcal{H}_R . Set

$$\alpha_R = \mathcal{F}_\tau(\Omega_R) = \inf\{\mathcal{F}_\tau(\Omega) : \Omega \in \mathcal{H}_R\};$$

note that α_R is nonincreasing in R and from Lemma 3.15, we have that $\alpha_R \geq \mathcal{E}(\Omega_R) \geq \mathcal{E}_{\min}(\infty) > -\infty$. We claim that so long as τ_0 and \mathfrak{T}_0 are chosen small enough, Ω_R has bounded diameter uniformly in R . Indeed, fix r small and apply the lower density estimate of Corollary 4.4 to Ω_R : for every $x \in \Omega_R$,

$$|\Omega_R \cap B_r(x)| \geq c.$$

As $|\Omega_R| \leq v_{\max}$, we this implies that Ω_R may be covered by a bounded number K of balls $B_r(x_k)$, with $x_k \in \Omega_R$. Assume that the union of these balls is disconnected: this implies that Ω itself cannot have

a connected component of measure greater than $|\Omega| - c$. Apply Lemma 3.16 with $\varepsilon = c/2$: this gives a contradiction to the third conclusion there. We infer that all the balls $B_r(x_k)$ have connected union of diameter at most $2Kr$, as promised.

It follows that there is an S and $e_R \in G_0$ such that $e_R(\Omega_R) \subseteq Q_S$ for all R (this uses only the diameter bound and that $\cup_{e \in G_0} e(Q_S) = M$ for some S , as in the proof of Lemma 3.15). Since \mathcal{F}_τ is invariant under isometries in G_0 , $\mathcal{F}_\tau(\Omega_R) = \mathcal{F}_\tau(e_R(\Omega_R))$ and $e_R(\Omega_R)$ is also a minimizer. This implies that for $R \geq S$,

$$\alpha_S \geq \alpha_R = \mathcal{F}_\tau(\Omega_R) = \mathcal{F}_\tau(e_R(\Omega_R)) \geq \alpha_S.$$

In particular, α_R is independent of R for $R \geq S$, and Ω_S minimizes \mathcal{F}_τ over all open bounded sets with $|\Omega| \leq v_{\max}$. The second conclusion follows from the fact that any open, bounded set Ω with $\mathcal{F}_\tau(\Omega) = \alpha_S$ lies in Q_R for a sufficiently large R , and so the above argument applies to Ω . \square

A consequence of Theorem 6.3 is that all constants pertaining to minimizers of \mathcal{F}_τ may be taken independently of R if M/G_0 is compact. This allows us to avoid tracking the dependence on R below, instead fixing a sufficiently large R and looking at minimizers in Q_R .

7. MEASURE-THEORETIC ESTIMATES

Having now established the existence of minimizers of the main energy, we move toward understanding some initial measure-theoretic properties of these minimizers. Ultimately, the results of this section will be used in Section 8 to derive the Euler-Lagrange equation satisfied by Ω .

In Section 7.1, we first prove that Ω is an NTA domain. This allows us to utilize an inhomogeneous boundary Harnack principle for NTA domains recently shown in [4] in Section 7.2 in order to prove some fine estimates for the Green's function. This allows us to show that the first eigenfunction of a minimizer grows at least linearly from the boundary in Proposition 7.8. Proposition 7.9 contains some finer Green's function estimates that will be crucial in deriving a useful form of the Euler-Lagrange equation in the following section. With these results in hand, in Section 7.3 we can recover some basic measure theoretic properties of Ω and understand the nontangential limits of $|\nabla u_\Omega|$ and $|\nabla w_\Omega|$ on the reduced boundary $\partial^*\Omega$, similarly to other recent approaches in vectorial free boundary problems [14, 32].

We assume throughout the section that R and $v < v_{\max}$ are fixed, and that $\eta < \eta_0(R, v, v_{\max})$, $\mathfrak{T} < \mathfrak{T}_0(R, v, v_{\max}, \eta)$, and $\tau < \tau_0(R, v, v_{\max}, \mathfrak{T}, \eta)$ are small enough that all results in earlier sections apply. We recall that the growth quantities $\text{DO}(\Omega)$ and $\text{UP}(\Omega)$ are defined in (2.14) and (2.13) respectively.

7.1. Density estimates and the NTA property. Let us summarize some consequences of the upper and lower bounds of Sections 4 and 5. The proofs are mostly standard and may be carried out in local coordinates, so we provide only brief sketches and references. First, we have uniform upper and lower bounds on the volume density of Ω and perimeter density of $\partial\Omega$.

Lemma 7.1. *Let $\Omega \in \mathcal{H}$ be a minimizer of \mathcal{F}_τ . Then there are $r_0, C > 0$ depending only on R, v, v_{\max}, η such that for any $x \in \partial\Omega$ and $r < r_0$,*

$$(7.1) \quad \frac{1}{C} < \frac{|B_r(x) \cap \Omega|}{|B_r(x)|} < 1 - \frac{1}{C}$$

and

$$(7.2) \quad \frac{1}{C} < \frac{\mathcal{H}^{n-1}(\partial\Omega \cap B_r(x))}{r^{n-1}} < C.$$

Sketch of proof. First, from Theorem 4.1 and Theorem 5.1, we have that $c \leq \text{DO}(\Omega) \leq \text{UP}(\Omega) \leq C$. To prove the lower bound in (7.1), let $y \in B_{r/2}(x) \cap \Omega$ be a point with $u_\Omega(y) + \sqrt{\mathfrak{T}}w_\Omega(y) \geq \text{DO}(\Omega)r/2$. By Corollary 5.5 we have $|\nabla(u_\Omega(y) + \sqrt{\mathfrak{T}}w_\Omega)| \leq C(\text{UP}(\Omega))$, and so we must have that $d(y, \partial\Omega) \geq \frac{\text{DO}(\Omega)r}{2C(\text{UP}(\Omega))} \geq cr$. So long as r_0 is small enough, $|B_{cr}(y)| \geq c'r^n$, $|B_r(x)| \leq Cr^n$, and so

$$\frac{|B_r(x) \cap \Omega|}{|B_r(x)|} \geq \frac{|B_{cr}(y)|}{|B_r(x)|} \geq \frac{1}{C}.$$

For the upper bound in (7.1), first note that if $B_{r/2}(x) \cap (M \setminus Q_R)$ is nonempty, then $|B_r(x) \setminus Q_R| \geq c(R)r^n$ using the smoothness of ∂Q_R , and this implies the estimate. If this is not the case, then $B_{r/2}(x) \subseteq Q_R$. We apply Lemma 5.2 to learn that if u, w are harmonic replacements for u_Ω, w_Ω respectively on $B_{r/2}(x)$, then

$$\int |\nabla(u_\Omega - u)|^2 + \mathfrak{I}|\nabla(w_\Omega - w)|^2 \leq C \left[|B_{r/2}(x) \cap \Omega| + r^n \sup_{B_{r/2}(x)} \mathfrak{I}(w_\Omega^2 + u_\Omega^2) \right].$$

Now, as $x \in \partial\Omega$, we have that $\sup_{B_{r/2}(x)} \mathfrak{I}w_\Omega^2 + u_\Omega^2 \leq C(\text{UP}(\Omega))r^2$. On the other hand, there must be a point $y \in B_{r/4}(x)$ with $\sqrt{\mathfrak{I}}w_\Omega(y) + u_\Omega(y) \geq \text{DO}(\Omega)r/4$. Recalling that $\Delta(\sqrt{\mathfrak{I}}w_\Omega(y) + u_\Omega(y)) \geq -C(R, v_{\max})$, applying [24, Theorem 8.16] to $u_\Omega - u, w_\Omega - w$ gives

$$\sup_{B_{r/2}(x)} \sqrt{\mathfrak{I}}[w_\Omega - w] + [u_\Omega(y) - u] \leq Cr^2.$$

At y then, $u(y) + \sqrt{\mathfrak{I}}w(y) \geq cr - Cr^2 \geq cr$ as long as r_0 is taken small enough. Using the Harnack inequality on $B_{r/2}(x)$, this implies $u + \sqrt{\mathfrak{I}}w \geq cr$ on $B_{r/4}(x)$. By contrast, we know that on $B_{\varepsilon r}(x)$, $u_\Omega + \sqrt{\mathfrak{I}}w_\Omega \leq \varepsilon r \text{UP}(\Omega) \leq \frac{cr}{2}$ as long as we choose ε small enough. Integrating and using the Poincaré inequality,

$$\begin{aligned} cr^{n+2} &\leq \int_{B_{\varepsilon r}(x)} |u_\Omega - u|^2 + \mathfrak{I}|w_\Omega - w|^2 \\ &\leq \int_{B_{r/2}(x)} |u_\Omega - u|^2 + \mathfrak{I}|w_\Omega - w|^2 \\ &\leq Cr^2 \int_{B_{r/2}(x)} |\nabla(u_\Omega - u)|^2 + \mathfrak{I}|\nabla(w_\Omega - w)|^2 \leq Cr^2 [|B_{r/2}(x) \cap \Omega| + r^{n+2}]. \end{aligned}$$

As long as r_0 is small enough, the last term may be reabsorbed on the left, giving $|B_r(x) \cap \Omega| \geq cr^n \geq c|B_r(x)|$.

For (7.2), the upper bound may be obtained as in [30, Lemma 2.4] or [32]. The lower bound follows from applying the relative isoperimetric inequality [23, 4.5.2(2)]. \square

The volume density bounds may be automatically improved slightly to give clean balls, or corkscrew points, instead.

Definition 7.2. *Let $\Omega \subseteq M$ be an open set. We say that Ω satisfies the inner (resp. outer) clean ball condition with constant K at $x \in \partial\Omega$ if for any $r < 1$, there is a point y such that $B_{r/K}(y) \subseteq B_r(x) \cap \Omega$ (resp. $B_{r/K}(y) \subseteq B_r(x) \setminus \Omega$). We say that Ω satisfies the inner (resp. outer) clean ball condition with constant K if it satisfies it at every $x \in \partial\Omega$.*

Corollary 7.3. *Let Ω be as in Lemma 7.1. Then there exists an $\alpha = \alpha(R, v, v_{\max}, \eta) > 0$ such that for any $x \in \partial\Omega$ and $r < r_0$ there are two balls $B_{\alpha r}(y) \subseteq (B_r(x) \cap \Omega)$ and $B_{\alpha r}(z) \subseteq (B_r(x) \setminus \Omega)$. In particular, Ω satisfies the inner and outer clean ball condition with constant depending only on r_0, α .*

Proof. For α small, take the collection $\{B_{\alpha r}(y) : y \in B_{r/2}(x) \setminus \Omega\}$ and from it pick a finite-overlapping subcover $\{B_{\alpha r}(y_k)\}_{k=1}^K$ of $B_{r/2}(x) \setminus \Omega$. Then if every one of $B_{\alpha r/2}(y_k)$ intersects $\partial\Omega$ at a point x_k , we have from Lemma 7.1 that

$$\mathcal{H}^{n-1}(B_{\alpha r}(y_k) \cap \partial\Omega) \geq \mathcal{H}^{n-1}(B_{\alpha r/2}(x_k) \cap \partial\Omega) \geq c(\alpha r)^{n-1}.$$

As $|B_{r/2}(x) \setminus \Omega| \geq c|B_{r/2}(x)| \geq cr^n$ and the $B_{\alpha r}(y_k)$ cover this set, we must have that $K(\alpha r)^n \geq cr^n$, or $K \geq c\alpha^{-n}$. Summing over k and using the finite overlapping property and the upper bound in (7.2),

$$\begin{aligned} Cr^{n-1} &\geq \mathcal{H}^{n-1}(\partial\Omega \cap B_r(x)) \\ &\geq c \sum_{k=1}^K \mathcal{H}^{n-1}(B_{\alpha r}(y_k) \cap \partial\Omega) \geq cK(\alpha r)^{n-1} \geq \frac{c}{\alpha} r^{n-1}. \end{aligned}$$

If α is taken small enough, this is a contradiction, so at least one of $B_{\alpha r/2}(y_k) \subseteq B_r(x) \setminus \Omega$. The inner clean ball may be found similarly. \square

Definition 7.4. Let $\Omega \subseteq M$ be an open set. We say that Ω satisfies the Harnack chain condition with constant K at $x, y \in \Omega$ if there is a curve $\gamma : [0, 1] \rightarrow \Omega$ with $\gamma(0) = x$, $\gamma(1) = y$, $l(\gamma([0, 1])) \leq Kd(x, y)$, and

$$d(\gamma(t), \partial\Omega) \geq \frac{1}{K} \min \{l(\gamma([0, t])), l(\gamma([t, 1]))\}.$$

Here $l(\gamma([a, b])) = \int_a^b |\dot{\gamma}| dt$ denotes length. We say Ω satisfies the Harnack chain condition if it satisfies it at every $x, y \in \Omega$. We say Ω is nontangentially accessible (NTA) with constant K if it satisfies the inner and outer clean ball conditions and the Harnack chain condition with constant K .

The next lemma follows from known results on Bernoulli-type free boundary problems [2], but we present a proof in Appendix B as we are unaware of a version with coefficients and nonzero right-hand side treated in the literature.

Lemma 7.5. Let $\Omega \in \mathcal{H}$ be a minimizer of \mathcal{F}_τ . Then there are $r_0, K > 0$ depending only on R, v, v_{\max} , and η such that for any $z \in \partial\Omega$ and any $x, y \in \Omega \cap B_{r_0}(z)$, Ω satisfies the Harnack chain condition at x and y with constant K .

From Lemma 7.5, we find that a minimizer Ω is an NTA domain.

Corollary 7.6. Let $\Omega \in \mathcal{H}$ be a minimizer of \mathcal{F}_τ . Then Ω is an NTA domain, with constant depending only on R, v, v_{\max} , and η .

Proof. First, observe that Ω is connected so long as $\eta_0, \mathfrak{F}_0, \tau_0, r$ are taken small enough. Indeed, we have from Lemma 3.14 that each connected component U of Ω besides one must have $|U| < \varepsilon$. On the other hand, Lemma 7.5 implies that $B_{r_0}(z) \cap \Omega$ lies within a single connected component of Ω for every $z \in \partial\Omega$, which when combined with Lemma 7.1 gives that $|U| \geq |\Omega \cap B_{r_0}(z)| \geq cr_0^n$; if ε is small enough then Ω has only one connected component.

Consider the set $U_r = \{x \in \Omega : d(x, \partial\Omega) > r\}$ for r small and fixed. We claim that any two points $x, y \in U_r$ may be connected by a path γ which has length bounded by $C(R, v, v_{\max}, \eta, r)$ and stays a distance $1/C$ away from $\partial\Omega$. To see this, cover Ω by a finitely-overlapping collection $\{B_r(z_k)\}_{k=1}^K$ with $z_k \in \Omega$. As $|\Omega| \leq C$, we have that $K \leq C$ from the finite-overlapping property and Lemma 7.1. For any two balls $B_r(z_k), B_r(z_j)$ with nontrivial intersection, we have two possibilities: either both $B_r(z_k), B_r(z_j) \subseteq U_r$, or at least one of them (say $B_r(z_k)$) intersects ∂U_r . In the first case, any pair of points $x \in B_r(z_k), y \in B_r(z_j)$ may be connected by a curve of length $2r$ which stays r away from $\partial\Omega$ as it stays inside $B_r(z_k) \cup B_r(z_j)$. In the second case, we have that as long as r is small enough, $B_r(z_k) \cup B_r(z_j) \subseteq B_{r_0}(x)$ for some $x \in \partial\Omega$. Applying Lemma 7.5 gives that for any $x \in B_r(z_k) \cap U_r$ and any $y \in B_r(z_j) \cap U_r$, x and y may be connected by a curve of length Cr staying r/C away from $\partial\Omega$.

Take a graph with vertices $\{B_r(z_k)\}$ and edges between $B_r(z_k), B_r(z_j)$ which have nontrivial intersection. As these balls cover Ω and Ω is connected, the graph must also be connected, and so any two balls may be connected by a chain of distinct, pairwise overlapping balls. For any $x \in B_r(z_k)$ and $y \in B_r(z_j)$ with $x, y \in U_r$, find such a path of pairwise intersection balls $B_r(z_{i_m})$, $m = 1, \dots, J$, with $k = i_1$ and $j = i_J$. We have shown that z_{i_m} and $z_{i_{m+1}}$, as well as x and z_{i_1} , z_{i_J} and y , may be connected by curves of length Cr and staying a distance C/r from $\partial\Omega$. By concatenating these curves (and using that $J \leq K$ is bounded), we see that x and y may be connected by a curve of uniformly bounded length remaining a distance r/C from the boundary, as promised.

This, together with Lemma 7.5, shows that Ω satisfies the Harnack chain condition. Combining with Corollary 7.3 implies Ω is an NTA domain. \square

7.2. Estimates on the Green's function. In Lemma 3.4, we saw from a basic maximum principle argument that $u_\Omega \leq Cw_\Omega$ for essentially arbitrary domains. The opposite inequality is far more subtle, and will in general fail even on polygonal domains in \mathbb{R}^n (in particular, it does not follow from the NTA property of Corollary 7.6). Nonetheless, we show in Proposition 7.8 below that it is, in fact, valid for *minimizing* Ω .

Let $G_\Omega(x, y)$ be the positive Green's function for Ω with pole at x . From [26, Theorem 1.2.8], we have that

$$(7.3) \quad \begin{aligned} G_\Omega(x, y) &\leq Cd^{2-n}(x, y) \quad \forall x, y \in Q_R \\ G_\Omega(x, y) &\geq cd^{2-n}(x, y) \quad \forall x, y \in \Omega, \quad d(x, y) \leq \frac{1}{2}d(x, \partial\Omega), \end{aligned}$$

where the constants depends only on R . If $n = 2$ the same is valid with $-\log d(x, y)$ in place of $d^{2-n}(x, y)$; we will only consider the case of $n > 2$ below, but all estimates remain valid if $n = 2$ after similar modification.

Lemma 7.7. *Let $\Omega \in \mathcal{H}$ be a minimizer of \mathcal{F}_τ . Then for every $c_0 > 0$ there is a C_0 depending only on $R, v, v_{\max}, \eta, c_0$ such that $w_\Omega(y) \leq C_0 G_\Omega(x, y)$ for any $y \in \Omega$ and any $x \in \Omega$ with $d(x, \partial\Omega) > c_0$.*

Proof. Choose $3r_0 \leq \min\{c_0, \text{inj}_M\}$. Let K be the NTA constant of Ω from Corollary 7.6. Let us first consider the case that y is such that $d(y, \partial\Omega) \geq r_0/K$. From the Harnack chain property of Corollary 7.6 we may find a sequence of finitely many balls (the number depending only on r_0 and K) $B_{c r_0/2}(y_k) \subseteq \Omega$ so that $B_{c r_0/2}(y_k) \cap B_{c r_0/2}(y_{k+1})$ is nonempty and $y_1 = y$ and $y_J = x$. Applying the Harnack inequality to $G_\Omega(x, \cdot)$ on each ball, we get

$$\sup_{z \in B_{c r_0/2}(y_k)} G_\Omega(x, z) \leq C \inf_{z \in B_{c r_0/2}(y_k)} G_\Omega(x, z)$$

for $k < J$, while the assumption that $d(x, \partial\Omega) > c_0$ guarantees that $G_\Omega(x, z) \geq c r_0^{2-n}$ for $z \in B_{c r_0/2}(x)$ from the Green's function lower bound in (7.3). Together these guarantee that for any such point y we have

$$(7.4) \quad G_\Omega(x, y) \geq c(r_0).$$

On the other hand, from Lemma 3.3 $w_\Omega(y) \leq C$ at y , while for any $z \in \Omega$ applying Theorem 4.1 and Lemma 3.4 gives

$$(7.5) \quad w_\Omega(z) \geq c[u_\Omega(z) + \sqrt{\mathfrak{X}}w_\Omega(z)] \geq c\text{DO}(\Omega)d(z, \partial\Omega).$$

Now apply [4, Theorem 2.2] on $B_{r_0}(z)$ for any $z \in \partial\Omega$, with $u_1 = w_\Omega/w_\Omega(y)$ and $u_2 = G_\Omega(x, \cdot)/G_\Omega(x, y)$, where $y \in B_{r_0}(z)$ is a point with $d(y, \partial\Omega) \geq r_0/K$ (such a point y exists from the inner clean ball property). In the theorem, we set $U = \Omega$, Q the collection of NTA domains, $\beta = 1$, $\zeta = 2$, $x^0 = y$, $V = \{u \in C(\bar{\Omega}) \cap C^2(\Omega) : u \geq 0, |\Delta u| \leq C\}$ and H the solution to the Dirichlet problem for Δ . The assumptions (P1-P7) follow from standard elliptic estimates and (P8) from [26, Lemma 1.3.7]. Using [4, Remark 2.4], we only need to verify the growth bound (2.1) for w_Ω , and it follows from (7.5). We conclude that

$$(7.6) \quad w_\Omega(a) \leq CG_\Omega(x, a)$$

for all $a \in B_{r_0}(z)$. Applying at any z , (7.6) remains valid for each $a \in \Omega$ with $d(a, \partial\Omega) < r_0$. For a with $d(y, \partial\Omega) \geq r_0$, (7.6) follows from (7.4) and that $|w_\Omega| \leq C$ instead. \square

We use the Green's function as a barrier for u from below and apply Lemma 7.7 to show that the torsion function bounds the eigenfunction from below:

Proposition 7.8. *Let $\Omega \in \mathcal{H}$ be a minimizer of \mathcal{F}_τ . Then $w_\Omega \leq C(R, v, v_{\max}, \eta)u_\Omega$ on Ω , and in particular*

$$\sup_{B_r(x)} u_\Omega \geq \frac{1}{C}\text{DO}(\Omega)r$$

for all $x \in \partial\Omega$ and $r \in (0, 1)$.

Proof. We have that $|\Omega| \leq v_{\max}$ and $\int u_\Omega^2 = 1$. This means there must be a point $x \in \Omega$ with $u_\Omega(x) \geq c$, and by Theorem 5.1, we have $d(x, \partial\Omega) \geq c/\text{UP}(\Omega) \geq c$. From the Lipschitz estimate of Corollary 5.5, we still have $u \geq c/2$ on a ball $B_r(x) \subseteq \Omega$, where $r > 0$ depends on c and $\text{UP}(\Omega)$, and thus on R, v, v_{\max} , and η . Using the Green's function upper bound in (7.3), we have $G_\Omega(x, y) \leq cr^{2-n} \leq c$ on $\partial B_r(x)$. As $\Delta u_\Omega \leq 0$ on Ω , we may use $cG_\Omega(x, \cdot)$ as a barrier from below for u_Ω on the set $\Omega \setminus \bar{B}_r(x)$; the comparison principle implies that

$$u_\Omega(y) \geq cG_\Omega(x, y) \geq cw_\Omega(y),$$

with the last inequality from Lemma 7.7. On $B_r(x)$, on the other hand, $u_\Omega \geq c/2$ while $w_\Omega \leq C$, so the same inequality follows. This completes the proof. \square

Next, we prove sharp bounds on G_Ω near $\partial\Omega$. To state these in a more useful fashion, recall that the *harmonic measure* on Ω , which we denote by ω_x , is defined as follows: for any $f \in C(\partial\Omega)$, let $\bar{f} \in C(\bar{\Omega})$ be the solution to the Dirichlet problem

$$\begin{cases} \Delta \bar{f} = 0 & \text{on } \Omega \\ \bar{f} = f & \text{on } \partial\Omega. \end{cases}$$

Then $f \mapsto \bar{f}(x)$ is a linear functional from $C(\partial\Omega)$ to \mathbb{R} , and from the maximum principle it has norm bounded by 1 and positive (i.e. $f \geq 0 \implies \bar{f}(x) \geq 0$). From the Riesz representation theorem, there is a positive Borel measure ω_x on $\partial\Omega$ such that

$$\int_{\partial\Omega} f d\omega_x = \bar{f}(x).$$

This is the harmonic measure; we clearly also have $\omega_x(\partial\Omega) = 1$.

Proposition 7.9. *Let $\Omega \in \mathcal{H}$ be a minimizer of \mathcal{F}_τ . Then the Green's function G_Ω satisfies*

$$(7.7) \quad c \frac{d(x, \partial\Omega)d(y, \partial\Omega)}{d^n(x, y)} \leq G_\Omega(x, y) \leq C \frac{d(x, \partial\Omega)d(y, \partial\Omega)}{d^n(x, y)}.$$

Moreover, for any $x \in \partial\Omega$, and any $y \in \Omega$ with $d(y, x) > 4r$, we have

$$(7.8) \quad c \frac{d(y, \partial\Omega)}{d^n(x, y)} \leq \frac{\omega_y(B_r(x))}{r^{n-1}} \leq C \frac{d(y, \partial\Omega)}{d^n(x, y)}$$

All constants depend only on v, v_{\max}, R , and η .

Proof. For any point $a \in W$, set $d_a = d(a, \partial\Omega)$. Take two points $x, y \in \Omega$. If $d(x, y) \leq \frac{1}{2}d_y$, then (7.7) follows directly from (7.3) (using the diameter bound for Ω in the lower bound) and there is nothing more to show. We now consider the case of $d(x, y) \geq \frac{1}{2} \max\{d_x, d_y\}$, breaking up the estimate into several cases depending on the locations of x and y . The notation $s \approx t$ below stands for $cs \leq t \leq Cs$ with constants depending only on v, v_{\max}, R, η .

Case 1: $2d(x, y) \geq d_y \geq r_0$. First, assume that $d_y \geq r_0$ for some fixed r_0 to be chosen below. We have from Lemma 7.7 that

$$cd_x \leq w_\Omega(x) \leq C(r_0)G_\Omega(y, x)$$

for any $x \in \Omega$. On the other hand, from (7.3) we have $G_\Omega(y, x) \leq C(r_0)$ for $x \in \partial B_{r_0/2}(y)$, while $u_\Omega(x) \geq cd_x$ there from Proposition 7.8. So, we may use a multiple of $G_\Omega(y, \cdot)$ as a lower barrier for u_Ω on the set $\Omega \setminus B_{r_0/2}(y)$: as $\Delta u_\Omega \leq 0$, from the comparison principle this leads to

$$G_\Omega(y, x) \leq Cu_\Omega(x) \leq CUP(\Omega)d_x$$

for $x \notin B_{r_0/2}(y)$. Together these two estimates show that in the case when $d_y \geq r_0$ and $d(x, y) \geq \frac{1}{2}d_y \geq \frac{r_0}{2}$, we have

$$(7.9) \quad G_\Omega(y, x) \approx d_x \approx \frac{d_x}{r_0^{n-1}} \approx \frac{d_x d_y}{d^n(x, y)}.$$

Case 2: $d_y \geq d(x, y)/4K$. Now fix r_0 small enough that $B_{2r_0}(y_0) \subseteq \Omega$ for some reference point y_0 , using the interior clean ball condition from Lemma 7.6. The next two cases we treat are when $d_y \geq \frac{1}{4K}d(x, y)$, with K the NTA constant of Ω . Take a point $a \in \partial B_{d_y/4}(y)$. The basic estimates (7.3) apply to this point to give $G_\Omega(y, a) \approx d^{2-n}(y, a) \approx d_y^{2-n}$. If $d_x \geq \frac{1}{4K}d_y$, then $d_x \approx d_y \approx d(x, y)$ are all comparable and we may use the Harnack chain condition to find a uniformly bounded chain of balls $B_r(x_i) \subseteq \Omega$ with $r \approx d_y$ connecting a and x ; applying the Harnack inequality finitely many times gives

$$(7.10) \quad G_\Omega(y, x) \approx G_\Omega(y, a) \approx d_y^{2-n} \approx \frac{d_y d_x}{d^n(x, y)}.$$

We may therefore assume that $d_x \leq \frac{1}{4K}d_y$. In this case, let $z \in \partial\Omega$ be a point with $d_x = d(x, z)$, and find an a' with $B_{d_y/4K}(a') \subseteq B_{d_y/4}(z) \cap \Omega$ using the interior clean ball property. Since $d_{a'} \geq d_y/4K$, we have $G_\Omega(y, a') \approx d_y^{2-n}$ from (7.10). On the other hand, from (7.9) we have $G_\Omega(y_0, a') \approx d_y$. Applying the

boundary Harnack principle [26, Lemma 1.3.7] to $G_\Omega(y, \cdot)$ and $G_\Omega(y_0, \cdot)$ on $B_{d_y/2}(z)$ (note that this excludes the pole at y), we have that

$$\frac{G_\Omega(y, x')}{G_\Omega(y_0, x')} \approx \frac{G_\Omega(y, a')}{G_\Omega(y_0, a')} \approx \frac{d_y^{2-n}}{d_y} \approx \frac{d_y}{d^n(x, y)}$$

for any $x' \in B_{d_y/4}(z)$ (with the last step using $d_y \approx d(x, y)$). In particular this is valid at $x' = x$, leading to

$$(7.11) \quad G_\Omega(y, x) \approx G_\Omega(y_0, x) \frac{d_y}{d^n(x, y)} \approx \frac{d_x d_y}{d^n(x, y)}$$

using (7.9) again.

Case 3: $\max\{d_y, d_x\} \leq d(x, y)/4K$. The only case remaining is when $\max\{d_y, d_x\} \leq \frac{1}{4K}d(x, y)$. Assume without loss of generality that $d_y \geq d_x$ and choose $z_x \in \partial\Omega$ with $d(z_x, x) = d_x$. Now use the clean ball property to find a point a_x with $B_{r/K}(a_x) \subseteq B_r(z_x) \cap \Omega$ for $r = \min\{d(x, y)/4, r_0\}$. Similarly construct z_y and a_y . At a_y , we have from (7.11) that

$$G_\Omega(y, a_y) \approx \frac{d_y d_{a_y}}{d^n(a_y, y)} \approx \frac{d_y}{r^{n-1}}.$$

Applying the Harnack inequality along a Harnack chain connecting a_x and a_y , this also gives $G_\Omega(y, a_x) \approx d_y/r^{n-1}$. We now use the boundary Harnack principle on $G_\Omega(y, \cdot)$ and $G_\Omega(y_0, \cdot)$ on the region $B_{2r}(z_x) \cap \Omega$ (which contains x) to give

$$\frac{G_\Omega(y, x)}{G_\Omega(y_0, x)} \approx \frac{G_\Omega(y, a_x)}{G_\Omega(y_0, a_x)} \approx \frac{d_y}{r^{n-1}} \frac{1}{r}.$$

In other words,

$$G_\Omega(y, x) \approx G_\Omega(y_0, x) \frac{d_y}{r^n} \approx \frac{d_x d_y}{r^n}.$$

Noting that $r \approx d(x, y)$ leads to the conclusion (7.7), which we have now established in all cases.

For (7.8), [26, Corollary 1.3.6] gives that at any $x \in \partial\Omega$, $r < r_0$, and $y \in \Omega \setminus B_{4r}(x)$, if z is a point with $B_{r/K}(z) \subseteq B_r(x) \cap \Omega$,

$$\omega_y(B_r(x)) \approx G_\Omega(z, y) r^{n-2}.$$

Applying (7.7) leads to

$$G_\Omega(z, y) \approx \frac{d_z d_y}{d^n(z, y)} \approx r \frac{d_y}{d^n(x, y)},$$

from which (7.8) follows. \square

Note that we used the upper and lower bounds $\text{UP}(\Omega), \text{DO}(\Omega)$ in a crucial way to get (7.9) and then (7.10) and the remaining estimates. For general NTA domains, the conclusions of this lemma need not hold.

A direct consequence of these estimates and Lemma 7.1 is that ω_y is absolutely continuous with respect to Hausdorff measure restricted to $\partial\Omega$, and the Radon-Nykodym derivative

$$(7.12) \quad K(x, y) = \frac{d\omega_y}{d\mathcal{H}^{n-1} \llcorner \partial\Omega}(x) \approx \frac{d(y, \partial\Omega)}{d^n(x, y)}$$

the *Poisson kernel*, is a bounded function of x satisfying similar estimates.

7.3. Blow-up analysis and the reduced boundary. Given a point $x \in \partial\Omega$ and vector $\nu \in T_x M$ with $|\nu| = 1$, let

$$B_{r,\nu}(x) = \exp_x \{v \in T_x M : |v| < r, g(v, \nu) < 0\} \subseteq B_r(x)$$

be a half-ball. For $r \leq \text{inj}_M$, define the function $l_{x,\nu} : B_r(x) \rightarrow \mathbb{R}$ be given by

$$(7.13) \quad l_{x,\nu}(\exp_x(v)) = g(v, -\nu)_+,$$

i.e. $l_{x,\nu}$ is a truncated linear function in normal coordinates.

Lemma 7.10. *Let Ω be a minimizer of \mathcal{F}_τ . For each $s \in (0, s_0)$, there are constants $\varepsilon(R, s, v, v_{\max}, \eta) > 0$ and $c = c(R, v, v_{\max}, \eta)$ such that if $x \in \partial\Omega$ has $|(B_r(x) \cap \Omega) \Delta B_{r,\nu}(x)| \leq \varepsilon r^n$ for some ν and $r < \varepsilon$, then:*

- (1) $B_{cr}(x) \cap \partial\Omega \subseteq \exp_x\{\nu \in T_x M : |\nu| < r, |g(\nu, \nu)| < sr\}$,
(2) $|u_\Omega(y) - \alpha l_{x,\nu}(y)| < sr$ for some $\alpha \in [c, 1/c]$ and for all $y \in B_{c\sqrt{sr}}(x)$,
(3) $|w_\Omega(y) - \beta l_{x,\nu}(y)| < sr$ for some $\beta \in [c, 1/c]$ and for all $y \in B_{c\sqrt{sr}}(x)$.

Proof. Conclusion (1) follows from standard geometric measure theory arguments using only Lemma 7.1: if there is a $y \in \partial\Omega \cap B_{cr}(x)$ with $y = \exp_x \nu$, $|g(\nu, \nu)| \geq sr$, then we may find clean balls $B_{csr}(y_1) \in \Omega \cap B_{sr/2}(y)$ and $B_{csr}(y_2) \in B_{sr/2}(y) \setminus \Omega$, with either both inside or both outside $B_{r,\nu}(x)$. This implies that $|(B_r(x) \cap \Omega) \Delta B_{r,\nu}(x)| \geq cs^n r^n$ (one of these two balls must be in this symmetric difference), and this is a contradiction if $\varepsilon \ll s^n$.

The other two conclusions follow from a compactness argument: assume, say, (2) is false for a given $s < s_0$, and take a sequence Ω_k of minimizers, points x_k in $\partial\Omega_k$, $r_k \rightarrow 0$, $\nu_k \in T_{x_k} M$ with norm 1, and $B_{r_k}(x_k)$ with

$$B_{r_k}(x_k) \cap \partial\Omega_k \subseteq \exp_x\{\nu \in T_{x_k} M : |\nu| < r_k, |g(\nu, \nu_k)| < s_k r_k\}$$

with $s_k \rightarrow 0$ (after applying (1) with s_k rather than with s , using that $\varepsilon_k \rightarrow 0$ to do so). Identify $T_{x_k} M$ with \mathbb{R}^n using an orthonormal basis with $e_n = \nu_k$, and define $\tilde{u}_k : B_1(0) \subseteq \mathbb{R}^n \rightarrow \mathbb{R}$ by

$$\tilde{u}_k(e) = \frac{u_k(r_k \exp_{x_k} e)}{r_k}.$$

Using Corollary 5.5, the functions \tilde{u}_k are uniformly Lipschitz functions defined on $B_1(0)$ and $\{\tilde{u}_k > 0\}$ converges to $B_{1,e_n}(0) := \{y \in B_1 : y \cdot e_n < 0\}$ in the Hausdorff topology. After passing to a subsequence, $\tilde{u}_k \rightarrow \tilde{u}_\infty$ uniformly on $B_1(0)$. Moreover, letting g_k denote the pullback of the metric g on $B_{r_k}(x_k)$ under the rescaled exponential map $e \in B_1(0) \mapsto \exp_{x_k}(r_k e)$, then $g_k \rightarrow g_\infty$ in C^2 topology where $g_\infty(e_i, e_j) \equiv \delta_{ij}$ the Euclidean metric.

Let Δ_k denote the Laplacian with respect to the metric g_k on $B_1(0)$. We have $\Delta_k \tilde{u}_k = -r_k \lambda_1(\Omega_k) \tilde{u}_k$ on $\{\tilde{u}_k > 0\}$; from elliptic estimates this means $\tilde{u}_k \rightarrow \tilde{u}_\infty$ in C^2 locally on $B_{1,e_n}(0)$. Rassing to the limit at any point in $B_{1,e_n}(0)$ gives $\Delta_\infty \tilde{u}_\infty = 0$. Therefore, \tilde{u}_∞ is harmonic on $B_{1,e_n}(0)$ and vanishes on the rest of $B_1(0)$ (using the uniform convergence again). In particular, using elliptic estimates means \tilde{u}_∞ is a piecewise smooth function on $B_1(0)$ with $|\nabla \tilde{u}_\infty| + |D^2 \tilde{u}_\infty| \leq C$ on $B_{1/2,e_n}(0)$.

We also have that $c\text{DO}(\Omega_k)r \leq \sup_{B_r(0)} \tilde{u}_k \leq C\text{UP}(\Omega_k)r$ for $r \in (0, 1)$ from Proposition 7.8, and the constants here are uniform in k . Passing to the limit, $cr \leq \sup_{B_r(0)} \tilde{u}_\infty \leq Cr$, which implies that $|\nabla \tilde{u}_\infty(0)| \geq c$ (from the $-e_n$ direction) and so $\tilde{u}_\infty(x) = \alpha(x \cdot -e_n)_+ + O(|x|^2)$ for some $c \leq \alpha \leq C$. On $B_{\kappa\sqrt{s}}(0)$, this implies that

$$\begin{aligned} |\tilde{u}_k(x) - \alpha(x \cdot -e_n)_+| &\leq |\tilde{u}_\infty(x) - \alpha(x \cdot -e_n)_+| + |\tilde{u}_k(x) - \tilde{u}_\infty(x)| \\ &\leq C|x|^2 + o_k(1) \leq C s \kappa^2 + o_k(1) < \frac{1}{2}s. \end{aligned}$$

in the last step we took κ small, and then k large. Changing back to the original variables,

$$\sup_{B_{\kappa\sqrt{sr}}(x_k)} |u_{\Omega_k} - \alpha l_{x_k, \nu_k}| < sr_k,$$

which is a contradiction to (2) failing for Ω_k at x_k . The argument for (3) is similar. \square

We say $x \in \partial^* \Omega$, the reduced boundary, if

$$\lim_{r \searrow 0} \frac{|(B_r(x) \cap \Omega) \Delta B_{r,\nu}(x)|}{|B_r(x)|} = 0$$

for some $\nu \in T_x M$ with $|\nu| = 1$; we will use the notation ν_x for this ν . Lemma 7.1 implies that $\mathcal{H}^{n-1}(\partial\Omega \setminus \partial^* \Omega) = 0$ (see [23, 4.5.11, 4.5.6]).

Corollary 7.11. *Let Ω be a minimizer of \mathcal{F}_τ , and $x \in \partial^* \Omega$. Then we have, for constants $0 < c \leq C < \infty$ depending only on R, v, v_{\max} , and η :*

(1) $\partial\Omega \cap B_r(x)$ converges to the approximate tangent plane $\partial B_{r,\nu_x}(x) \cap B_r(x)$ in the following Hausdorff sense:

$$\lim_{r \searrow 0} \frac{1}{r} \sup_{y \in \partial\Omega \cap B_r(x)} d(y, \partial B_{r,\nu_x}(x) \cap B_r(x)) = 0.$$

(2) $\mathcal{H}^{n-1}(\partial\Omega \cap B_r(x)) = \omega_{n-1} r^{n-1} + o(r^{n-1})$, where ω_{n-1} is the measure of the unit ball in \mathbb{R}^{n-1} .

(3) $u_\Omega(y) = |\nabla u_\Omega(x)| l_{x,\nu_x}(y) + o(d(x,y))$ for some number $|\nabla u_\Omega(x)| \in [c, C]$.

(4) $w_\Omega(y) = |\nabla w_\Omega(x)| l_{x,\nu_x}(y) + o(d(x,y))$ for some number $|\nabla w_\Omega(x)| \in [c, C]$.

The second conclusion can be found in [23, 4.5.6(2)]; the others follow directly from Lemma 7.10.

We say that a function $u \in C(\Omega)$ converges nontangentially to α at $x \in \partial\Omega$ if

$$\lim_{r \searrow 0} \sup \{ |u(y) - \alpha| : y \in B_r(x), d(y, \partial\Omega) > \varepsilon r \}$$

for all $\varepsilon > 0$. While the numbers $|\nabla u_\Omega(x)|$ in this corollary were abstract (not directly connected to ∇u_Ω proper), they may be reinterpreted as nontangential limits of $|\nabla u_\Omega|$:

Corollary 7.12. *Let Ω be a minimizer of \mathcal{F}_τ , $x \in \partial^*\Omega$, and $|\nabla u_\Omega(x)|, |\nabla w_\Omega(x)|$ as defined in Corollary 7.11. Then:*

(1) $-|\nabla u_\Omega(x)| \nu_x$ is the nontangential limit of ∇u_Ω at x .

(2) $-|\nabla w_\Omega(x)| \nu_x$ is the nontangential limit of ∇w_Ω at x .

Proof. We only prove (1), setting $\alpha = |\nabla u(x)|$. Using Corollary 7.11,

$$\lim_{r \searrow 0} \frac{\sup_{y \in B_r(x)} |u_\Omega(y) - \alpha l_{x,\nu_x}(y)|}{r} = 0.$$

Working in normal coordinates, at any $y \in B_{r,\nu_x}(x)$ the function $l_{x,\nu_x}(y)$ has

$$|\Delta l_{x,\nu_x}(y)| \leq C \|g_x - g_{euc}\|_{C^1} \leq Cr \quad |\nabla l_{x,\nu_x}(y) - \nabla l_{x,\nu_x}(x)| \leq Cr$$

where g_{euc} denotes the Euclidean metric. Fix $\varepsilon > 0$ and take $y \in B_r(x)$ with $d(x,y) \geq \varepsilon r$. Then applying elliptic estimates on $B_{r\varepsilon}(y)$,

$$\begin{aligned} \|\nabla(u_\Omega - \alpha l_{x,\nu_x})\|_{L^\infty(B_{\varepsilon r/2}(x))} &\leq C \left[r \|\Delta(u_\Omega - \alpha l_{x,\nu_x})\|_{L^\infty(B_{\varepsilon r}(x))} + \frac{1}{r} \|u_\Omega - \alpha l_{x,\nu_x}\|_{L^\infty(B_{\varepsilon r}(x))} \right] \\ &\leq Cr + o_r(1) = o_r(1), \end{aligned}$$

using that $|\Delta u_\Omega| \leq C$ in the last step. The constant here depends only on ε . It follows that

$$|\nabla u_\Omega(y) + \alpha \nu_x| \leq o_r(1) + |\alpha| |\nu_x + \nabla l_{x,\nu_x}(y)| = o_r(1) + |\alpha| |\nabla l_{x,\nu_x}(x) - \nabla l_{x,\nu_x}(y)| \leq o_r(1).$$

□

8. THE EULER-LAGRANGE EQUATION

Our next goal is to derive the Euler-Lagrange equation satisfied by minimizers by differentiating \mathcal{F}_τ with respect to smooth families of diffeomorphisms applied to Ω . There will essentially be three main steps to do this. First, in Section 8.1, we derive the distributional form of the Euler Lagrange equation. Computing the derivatives of the terms appearing in \mathcal{E} is more or less routine for smooth sets in Euclidean space, though we must take some care because minimizers are not necessarily smooth and we are working on a Riemannian manifold. For the distributional form of the Euler-Lagrange equation, the derivatives of the nonlinear term are essentially given in condition (N4) in the definition of admissible nonlinearity (Definition 2.11).

Next, in Section 8.2, we derive a pointwise form of the distributional Euler-Lagrange equation from Section 8.1. The basic idea is to plug into the distributional Euler-Lagrange equation a sequence of vector fields that approximate the outer unit normal to Ω at each point in the reduced boundary. Carrying this out for the terms coming from the base energy is not difficult, but it turns out to be fairly delicate for the term \mathfrak{h} for reasons discussed further below. This will require the optimal Green's function estimates from Proposition 7.9 of the previous section. Ultimately, we are able to arrive at the following pointwise form of

the free boundary condition (see Corollary 8.10 in Section 8.2). Then there is a constant A_0 such that for almost every $x \in \partial^* \Omega \cap Q_R$, we have the identity

$$(8.1) \quad -|\nabla u_\Omega(x)|^2 - \frac{\mathfrak{T}}{2} |\nabla w_\Omega(x)|^2 + \tau \left[b_\Omega(x) + \int_\Omega a_\Omega v_\Omega^x \right] = -A_0.$$

If $|\Omega| \neq v$, then $A_0 = f'_{v,\eta}(|\Omega|)$. Here, for each $x \in \partial\Omega \cap Q_R$, the function v_Ω^x is the solution of an auxiliary PDE.

It is not clear that the Euler-Lagrange equation (8.1) behaves as a Bernoulli-type problem or should yield any regularity. The most important part of this section, which is the content of Section 8.3, is to rewrite the Euler-Lagrange equation in a way which involves only $|\nabla u_\Omega|$, up to “lower order” terms. In doing so, the term $b_\Omega(x)$ is more or less innocuous and will be controlled by simply choosing the parameter τ to be sufficiently small. The $|\nabla w_\Omega|$ term is a priori problematic, but can be handled similarly to [14, 32]. The most challenging term to control is $\int_\Omega a_\Omega v_\Omega^x$, because the function v_Ω^x solves an auxiliary PDE that does not immediately lend itself to comparison with u_Ω . We will discuss how to overcome this difficulty at the beginning of Section 8.3 once we have seen introduced the equation for v_Ω^x . The main result of this section is the following free boundary condition.

Theorem 8.1. *There is a function $\rho \geq -C\tau$ with $\|\rho\|_{C^{0,\alpha}(\partial\Omega)} \leq C$, C, α depending only on v, v_{\max}, η, R , such that*

$$|\nabla u_\Omega(x)|^2 (1 + \rho(x)) = A_0$$

for \mathcal{H}^{n-1} -a.e. $x \in \partial^* \Omega \cap Q_R$. The constant A_0 has $A_0 \in [1/C, C]$.

As always, we assume that R is fixed, $\eta < \eta_0(R, v, v_{\max})$, $\mathfrak{T} < \mathfrak{T}_0(R, v, v_{\max}, \eta)$, and $\tau < \tau_0(R, v, v_{\max}, \mathfrak{T}, \eta)$ are small enough that all results in earlier sections apply. We recall that a_Ω and b_Ω are functions given in the definition of admissible nonlinearity in Definition 2.11.

8.1. The first variation along a vector field. The main goal of this subsection is to derive the distributional form of the Euler-Lagrange equation in Lemma 8.5. We start with some auxiliary lemmas quantifying how eigenvalues and eigenfunctions vary under change of domains. We let $\nu = \nu_\Omega$ denote the outer unit normal of Ω . If we knew that Ω had smooth boundary, then the following lemma would follow easily from the standard Hadamard variational formula; see [?]. Since thus far we only know that Ω satisfies some very basic measure theoretic properties, the proof requires some more care.

Lemma 8.2. *Let Ω be a minimizer of \mathcal{F}_τ . Let $\phi_t : M \rightarrow M$ be a one-parameter family of smooth diffeomorphisms with $\phi_0(x) = x$ and $\partial_t \phi_t|_{t=0} = T$ for a vector field T on M . Set $\Omega_t = \phi_t(\Omega)$. Then:*

- (1) $\limsup_{t \rightarrow 0} \frac{1}{|t|} (\lambda_1(\Omega_t) - \lambda_1(\Omega) + t \int_{\partial^* \Omega} |\nabla u_\Omega|^2 g(T, \nu) d\mathcal{H}^{n-1}) \leq 0$
- (2) $\limsup_{t \rightarrow 0} \frac{1}{|t|} (\text{tor}(\Omega_t) - \text{tor}(\Omega) + t \frac{1}{2} \int_{\partial^* \Omega} |\nabla w_\Omega|^2 g(T, \nu) d\mathcal{H}^{n-1}) \leq 0$
- (3) $\lim_{t \rightarrow 0} \frac{1}{t} (|\Omega_t| - |\Omega| - t \int_{\partial^* \Omega} g(T, \nu) d\mathcal{H}^{n-1}) = 0$
- (4) $\lambda_2(\Omega_t) \geq \lambda_1(\Omega_t) + c(R, v, v_{\max})$ for all $|t|$ small enough.
- (5) $\lim_{t \rightarrow s} \frac{1}{|t-s|} |\lambda_1(\Omega_t) - \lambda_1(\Omega_s) + (t-s)f_s^\phi| = 0$, where $f_s^\phi : (-c, c) \rightarrow (0, \infty)$ is a continuous function of s .

In particular, $\lambda_1(\Omega_t)$ is differentiable and $\partial_t \lambda_1(\Omega_t)|_{t=0} = - \int_{\partial^* \Omega} |\nabla u_\Omega|^2 g(T, \nu) d\mathcal{H}^{n-1}$.

Proof. Step 1: Some general expressions. Take any $v \in L^2(\Omega)$, and let $v_t = v \circ \phi_t^{-1} \in L^2(\Omega_t)$. Then we may compute

$$(8.2) \quad \int_{\Omega_t} v_t^2 = \int_\Omega v^2 |\det d\phi_t| = \int_\Omega v^2 (1 + t \operatorname{div} T + O(t^2)),$$

where the O depends only on $\|\phi_t\|_{C^1}$. If instead we consider a $v \in H^1(\Omega)$,

$$\int_{\Omega_t} |\nabla v_t|^2 = \int_\Omega |d\phi_t^{-1}(\nabla v)|^2 |\det d\phi_t| = \int_\Omega [|\nabla v|^2 (1 + t \operatorname{div} T) - 2tg(\nabla v, \nabla_{\nabla v} T)] + O(t^2) \int_\Omega |\nabla v|^2.$$

If v is smooth on Ω , bounded, and $-\Delta v = f \in L^2(\Omega)$, this may be further rewritten using the identity

$$\operatorname{div} (|\nabla v|^2 T - 2g(\nabla v, T)\nabla v) = |\nabla v|^2 \operatorname{div} T + 2g(\nabla v, T)f - 2g(\nabla v, \nabla_{\nabla v} T) \in L^2(\Omega)$$

to give

$$\int_{\Omega_t} |\nabla v_t|^2 = \int_{\Omega} |\nabla v|^2 + t \int_{\Omega} -2g(\nabla v, T)f + \operatorname{div} (|\nabla v|^2 T - 2g(\nabla v, T)\nabla v) + O(t^2) \int_{\Omega} |\nabla v|^2.$$

Applying the divergence theorem of [16] to the second term and assuming ∇v has nontangential limits at \mathcal{H}^{n-1} -a.e. point of $\partial\Omega$ lets us rewrite

$$\int_{\Omega_t} |\nabla v_t|^2 = \int_{\Omega} |\nabla v|^2 - t \int_{\Omega} 2g(\nabla v, T)f + t \int_{\partial^* \Omega} |\nabla v|^2 g(T, \nu) - 2g(\nabla v, T)g(\nabla v, \nu) d\mathcal{H}^{n-1} + O(t^2) \int_{\Omega} |\nabla v|^2.$$

Step 2: Proofs of (1)-(4). We now verify the conclusions of the lemma. For (1), take $v = u_{\Omega}$, and use v_t as a competitor for the definition of $\lambda_1(\Omega_t)$. Then $v_t \in H_0^1(\Omega_t)$, and

$$\int_{\Omega_t} v_t^2 = \int_{\Omega} u_{\Omega}^2 (1 + t \operatorname{div} T + O(t^2)) = 1 - 2t \int_{\Omega} u_{\Omega} g(\nabla u_{\Omega}, T) + O(t^2),$$

integrating by parts as $u_{\Omega} \in H_0^1(\Omega)$. On the other hand, ∇u_{Ω} has nontangential limits a.e. (from Corollary 7.12), and so

$$\int_{\Omega_t} |\nabla v_t|^2 = \lambda_1(\Omega) - t \int_{\Omega} 2g(\nabla u_{\Omega}, T)\lambda_1(\Omega)u_{\Omega} + t \int_{\partial^* \Omega} |\nabla u_{\Omega}|^2 g(T, \nu) - 2g(\nabla u_{\Omega}, T)g(\nabla u_{\Omega}, \nu) d\mathcal{H}^{n-1} + O(t^2)$$

Using the formula from Corollary 7.12 on the boundary gives $g(\nabla u_{\Omega}, T)g(\nabla u_{\Omega}, \nu) = |\nabla u_{\Omega}|^2 g(T, \nu)$, so this leads to

$$\lambda_1(\Omega_t) \leq \frac{\int |\nabla v_t|^2}{\int u_t^2} \leq \lambda_1(\Omega) - t \int_{\partial^* \Omega} |\nabla u_{\Omega}|^2 g(T, \nu) d\mathcal{H}^{n-1} + O(t^2),$$

as the terms of the form $t \int_{\Omega} u_{\Omega} g(\nabla u_{\Omega}, T)$ cancel to order t . This proves (1).

The proof of (2) is analogous: proceeding in the same way using $v = w_{\Omega}$, the terms $t \int_{\Omega} g(\nabla u_{\Omega}, T)$ in the expansion of $\operatorname{tor}(\Omega_t)$ again cancel and we find

$$\operatorname{tor}(\Omega_t) \leq \int \frac{1}{2} |\nabla v_t|^2 - v_t \leq \operatorname{tor}(\Omega) - \frac{t}{2} \int_{\partial^* \Omega} |\nabla w_{\Omega}|^2 g(T, \nu) d\mathcal{H}^{n-1} + O(t^2).$$

To prove (3), we use the constant function $v = 1$ in (8.2) and the divergence theorem to obtain

$$|\Omega_t| = |\Omega| + t \int_{\partial^* \Omega} g(T, \nu) d\mathcal{H}^{n-1} + O(t^2).$$

Using (1), (2), and (3), we see that $\mathcal{E}(\Omega_t) \leq \mathcal{E}(\Omega) + C|t| \leq \mathcal{E}_{\min} + \delta$ if $|t|$ is small enough. From Lemma 3.14, this implies that $\lambda_2(\Omega_t) > \lambda_1(\Omega) + c$, giving (4).

Step 3: Proof of (5). Finally, for (5) we proceed as for (1), except with Ω_s in place of Ω , using $u_{s,t} = u_{\Omega_s} \circ \phi_s \circ \phi_t^{-1} : \Omega_t \rightarrow \mathbb{R}$ as a competitor. These are nonnegative and have

$$(8.3) \quad \int_{\Omega_t} u_{s,t}^2 = 1 + O(|t-s|)$$

from (8.2). We do not try to justify the application of the divergence theorem in this case, instead only using

$$\begin{aligned} \lambda_1(\Omega_t) &\leq \frac{\int_{\Omega_t} |\nabla u_{s,t}|^2}{\int_{\Omega_t} u_{s,t}^2} \\ &= \lambda_1(\Omega_s) + (t-s) \int_{\Omega} |\nabla u_{\Omega_s}|^2 \operatorname{div} T_s - 2g(\nabla u_{\Omega_s}, \nabla_{\nabla u_{\Omega_s}} T_s) + 2u_{\Omega_s} g(\nabla u_{\Omega_s}, T_s) + O((t-s)^2) \\ &:= \lambda_1(\Omega_s) + (t-s) f_s^{\phi} + O((t-s)^2), \end{aligned}$$

where $T_s = \partial_t \phi_t|_{t=s}$. Clearly f_s^ϕ is bounded in terms of ϕ and $\lambda_1(\Omega_t)$, and $\lambda_1(\Omega_t)$ is bounded from (1). Applying with t and s reversed shows that

$$(8.4) \quad \lambda_1(\Omega_s) \leq \lambda_1(\Omega_t) + (s-t)f_t^\phi + O((t-s)^2),$$

and so $|\lambda_1(\Omega_s) - \lambda_1(\Omega_t)| \leq C|s-t|$. Together with (4), Lemma 3.8, and (8.3) this implies that

$$\begin{aligned} \|u_{s,t} - u_{\Omega_t}\|_{H_0^1(\Omega_t)}^2 &\leq \left\| \frac{u_{s,t}}{\sqrt{\int u_{s,t}^2}} - u_{\Omega_t} \right\|_{H_0^1(\Omega_t)}^2 + C|t-s| \\ &\leq C \left[\frac{\int_{\Omega_t} |\nabla u_{s,t}|^2}{\int_{\Omega_t} u_{s,t}^2} - \lambda_1(\Omega_t) \right] + C|t-s| \leq C|t-s|. \end{aligned}$$

On the other hand, as $u_{s,t}$ are defined by u_{Ω_s} composed with smooth diffeomorphisms approaching the identity, $\lim_{t \rightarrow 0} \|u_{s,t} - u_{\Omega_s}\|_{H^1(M)} = 0$ from the Vitali convergence theorem. Together, these give that $u_{\Omega_t} \rightarrow u_{\Omega_s}$ strongly in H^1 as $t \rightarrow s$, and it follows that f_t^ϕ is a continuous function of t . Combining with (8.4) leads to

$$\lambda_1(\Omega_s) \leq \lambda_1(\Omega_t) + (s-t)f_s^\phi + o(t-s).$$

This completes the proof of (5). The final statement of the lemma is immediate from (1) and (5). \square

We now turn to a finer analysis of how the eigenfunctions u_{Ω_t} vary under domain variation. We already saw in the proof of (5) that they vary continuously in t , in H^1 norm; however, we will need a much more careful estimate to handle the term \mathfrak{h} . Our first goal is a kind of $C_t^1 L_x^2$ estimate, which we break into two parts.

Lemma 8.3. *Let Ω be a minimizer of \mathcal{F}_τ . Let g_t be a smooth (in both x and t , uniformly on Ω), one-parameter family of metrics with $g_0 = g$, the original metric on M , and let Δ_t and m_t respectively be the Laplace-Beltrami operator and volume measure for the metric g_t . Let f_t be a C^1 function of t that is constant in x with $f_0 = \lambda_1(\Omega)$. Assume that $u_t \in H_0^1(\Omega)$ is the unique nonnegative function with $\int u_t^2 dm_t = 1$ and*

$$-\Delta_t u_t = f_t u_t,$$

and assume that no solutions $v \in H_0^1(\Omega)$ to

$$-\Delta_t v = \lambda f_t v$$

exist for any $\lambda \in (0, 1+c)$ (for some $c > 0$) apart from scalar multiples of u_t . Then u_t is continuously differentiable in t for $|t| < c$ and any $x \in \Omega$, the derivative v at $t=0$ lies in $C(\bar{\Omega}) \cap H_0^1(\Omega)$, and

$$\lim_{t \rightarrow 0} \frac{1}{t} \|u_t - u_0 - tv\|_{H_0^1(\Omega)} = 0.$$

Note that the existence of u_t is part of the lemma's hypotheses; we do not claim such a u_t exists.

Proof. Observe that $\|u_t\|_{H_0^1(\Omega)}$ is uniformly bounded in t , using u_t as a test function for itself:

$$(8.5) \quad \int |\nabla u_t|_t^2 dm_t = f_t \int u_t^2 dm_t \leq f_t.$$

For a $q \in H^{-1}(\Omega)$ consider the problem of finding a $p \in H_0^1(\Omega)$ which solves $-\Delta_t p = f_t p + q$, with $\int p u_t = 0$. Such a p exists if and only if $q(u_t) = 0$; if it exists it is unique and

$$(8.6) \quad \|p\|_{H^1(\Omega)} \leq C \|q\|_{H^{-1}(\Omega)}.$$

Indeed, the operator $-\Delta_t - f_t I : H_0^1(\Omega) \rightarrow H^{-1}(\Omega)$ is bounded and, by our assumptions, has a one-dimensional kernel spanned by u_t . From the Fredholm alternative and open mapping theorem (see [10]) that the range consists of all $q \in H^{-1}(\Omega)$ with $q(u_t) = 0$, and

$$\int \left| \nabla \left[p - u_t \int g_t(\nabla u_t, \nabla p) \right] \right|_{g_t}^2 dm_t \leq C \|q\|_{H^{-1}}^2.$$

Then (8.6) follows, as $\int g_t(\nabla u_t, \nabla p) = \int f_t u_t p = 0$ from the definition of u_t .

Insert u_t into the equation for u_s , to get (in weak form with $\phi \in H_0^1(\Omega)$ a test function)

$$\begin{aligned}
 \int g_s(\nabla u_t, \nabla \phi) dm_s &= \int g_s(\nabla u_t, \nabla \phi) dm_t + \int g_s(\nabla u_t, \nabla \phi) d(m_s - m_t) \\
 &= f_t \int u_t \phi dm_t + \int g_s(\nabla u_t, \nabla \phi) - g_t(\nabla u_t, \nabla \phi) dm_t + \int g_s(\nabla u_t, \nabla \phi) d(m_s - m_t) \\
 &= f_s \int u_t \phi dm_s + \int (f_t - f_s) u_t \phi + g_s(\nabla u_t, \nabla \phi) - g_t(\nabla u_t, \nabla \phi) dm_t \\
 &\quad + \int -f_s u_t \phi + g_s(\nabla u_t, \nabla \phi) d(m_s - m_t) \\
 &= f_s \int u_t \phi dm_s + q_s^t(\phi),
 \end{aligned}$$

where $q_s^t \in H^{-1}(\Omega)$ has

$$|q_s^t(\phi)| \leq \left[\|f_t - f_s\|_{L^\infty} + C \sup_{\Omega} |g_t - g_s| \right] \|u_t\|_{L^2} \|\phi\|_{L^2} + \sup_{\Omega} |g_t - g_s| \|u_t\|_{H_0^1(\Omega)} \|\phi\|_{H_0^1(\Omega)},$$

and in particular $\|q\|_{H^{-1}(\Omega)} \leq C|t - s|$. Here $|g_t - g_s| = |g_t - g_s|_g$ is with respect to the original metric. Subtracting u_s , we have (in weak form)

$$(8.7) \quad -\Delta_s(u_t - u_s) = f_s(u_t - u_s) + q_s^t.$$

By applying (8.6) to $u_t - u_s - u_s \int u_s(u_t - u_s) dm_s = u_t - au_s$, we see that

$$\|u_t - au_s\|_{H_0^1} \leq C \|q_s^t\|_{H^{-1}} \leq C|t - s|.$$

The constant a need not be 1, but we do know that as $u_t, u_s \geq 0$ by assumption, $a = \int u_t u_s dm_s \geq 0$. Using the Poincaré and triangle inequalities and the normalization on u_s, u_t ,

$$\begin{aligned}
 |1 - a| &= \left| \|u_t\|_{L^2(dm_t)} - \|au_s\|_{L^2(dm_s)} \right| \leq \left| \|u_t\|_{L^2(dm_s)} - \|au_s\|_{L^2(dm_s)} \right| + \left| \|u_t\|_{L^2(dm_s)} - \|u_t\|_{L^2(dm_t)} \right| \\
 &\leq \|u_t - au_s\|_{L^2(dm_s)} + C|t - s| \leq C|t - s|.
 \end{aligned}$$

Therefore, we see that

$$(8.8) \quad \|u_t - u_s\|_{H_0^1(\Omega)} \leq \|u_t - au_s\|_{H_0^1(\Omega)} + |1 - a| \|u_s\|_{H_0^1(\Omega)} \leq C|t - s|.$$

Set

$$q_t'(\phi) := - \int (\partial_t f_t) u_t \phi - (\partial_t g_t)(\nabla u_t, \nabla \phi) + m_t' [f_t u_t \phi - g_t(\nabla u_t, \nabla \phi)] dm_t,$$

where m_t' stands for the derivative of the volume form (i.e. $m_t(E) - m_s(E) = \int_s^t \int_E m_r' dm_r dr$). Then $|q_t'(\phi)| \leq C \|\phi\|_{H_0^1} \|u_t\|_{H_0^1}$, so $\|q_t'\|_{H^{-1}} \leq C$. On the other hand, directly estimating each term using the assumptions on f_t and g_t leads to

$$\begin{aligned}
 |q_s^t(w) - (s - t)q_t'(w)| &\leq C \left[\sup_{\Omega} |f_t - f_s| + (s - t) \partial_t f_t + \sup_{\Omega} |g_t - g_s| + (s - t)g_t' + \sup_{\Omega} |f_t - f_s|^2 + |g_t - g_s|^2 \right] \\
 &\quad \cdot \|u_t\|_{H_0^1(\Omega)} \|\phi\|_{H_0^1(\Omega)} \\
 &\leq o(|t - s|) \|u_t\|_{H_0^1(\Omega)} \|\phi\|_{H_0^1(\Omega)},
 \end{aligned}$$

so $\|q_s^t - (s - t)q_t'\|_{H^{-1}(\Omega)} = o(|t - s|)$. Using (8.8), $\|q_t' - q_s'\|_{H^{-1}(\Omega)} \rightarrow 0$ as $t \rightarrow s$, so this may be rewritten as

$$\|q_s^t - (s - t)q_s'\|_{H^{-1}(\Omega)} = o(|t - s|).$$

We also claim that $q_s'(u_s) = 0$. Indeed, as $-\Delta_s u_t = f_s u_t + q_s^t$ is a nontrivial solution, from the Fredholm alternative we must have $q_s^t(u_s) = 0$. But then

$$|t - s| |q_t'(u_s)| = o(|t - s|),$$

so taking $t \rightarrow s$ implies $q_s'(u_s) = 0$. Let $v_s^0 \in H_0^1(\Omega)$ be the unique solution to $-\Delta v_s^0 = f_s v_s^0 + q_s'$ with $\int u_s v_s^0 = 0$, using the Fredholm alternative from (8.6). Then together with (8.7), we have that

$$\|u_t - au_s - (s - t)v_s^0\|_{H_0^1} \leq C \|q_s^t - (s - t)q_s'\|_{H^{-1}(\Omega)} = o(|t - s|),$$

where $a = \int u_t u_s dm_s$ as before. Now set $v_s^1 = u_s \frac{1}{2} \int u_s^2 m'_s dm_s$, and estimate a again:

$$\begin{aligned} \left| \int u_t^2 dm_t - \int u_s^2 dm_s - (t-s) \int u_s^2 m'_s dm_s \right| &\leq \left| \int u_t^2 dm_t - \int u_s^2 dm_s - (t-s) \int u_t^2 m'_s dm_s \right| + C|t-s|^2 \\ &\leq C|t-s|^2, \end{aligned}$$

so taking square roots and recalling that $\int u_t^2 dm_t = 1$,

$$\begin{aligned} \left| \|u_t\|_{L^2(dm_t)} - \|u_t\|_{L^2(dm_s)} - (t-s) \frac{1}{2} \int u_s^2 m'_s dm_s \right| \\ \leq C|t-s|^2 + \frac{|2-1-\|u_t\|_{L^2(dm_s)}|}{1+\|u_t\|_{L^2(dm_s)}} |t-s| \int u_s^2 |m'_s| dm_s \leq C|t-s|^2. \end{aligned}$$

This may be used to estimate $1-a$ up to a correction:

$$\begin{aligned} \left| 1-a - (t-s) \frac{1}{2} \int u_s^2 m'_s dm_s \right| &\leq \left| \|u_t\|_{L^2(dm_s)} - a \|u_s\|_{L^2(dm_s)} \right| + C|t-s|^2 \\ &\leq \left| \|u_t\|_{L^2(dm_s)} - \|au_s - (s-t)v_s^0\|_{L^2(dm_s)} \right| + C|t-s|^2 \\ &\leq \|u_t - au_s - (s-t)v_s^0\|_{H_0^1} + C|t-s|^2 = o(|t-s|). \end{aligned}$$

The second step used that $\int u_s v_s^0 dm_s = 0$. Thus,

$$\begin{aligned} \|u_t - u_s - (s-t)[v_s^0 + v_s^1]\|_{H_0^1} &= \|u_t - au_s - (s-t)v_s^0\|_{H_0^1} + \left| 1-a - (t-s) \frac{1}{2} \int u_s^2 m'_s dm_s \right| \\ &= o(|t-s|). \end{aligned}$$

Set $v_s = v_s^0 + v_s^1$. This establishes the main conclusion of the lemma. It remains to verify that $v_0 \in C(\bar{\Omega})$ and that u_t is continuously differentiable on Ω .

These can be seen from the PDE for v_s^0 . As u_t is smooth on the interior of Ω , as long as $\phi \in H_0^1(U)$ for $U \subset\subset \Omega$, we have (integrating by parts as needed) that $|q'_s(\phi)| \leq C(U)\|\phi\|_{L^1}$. In other words, q'_s (and similarly q_s^t) may be represented by a bounded function. Applying elliptic regularity estimates to (8.7) gives that for $U' \subset\subset U$,

$$\|u_t - u_s - (s-t)v_s\|_{C^{1,\alpha}(U')} \leq C(U, U') [\|u_t - u_s - (s-t)v_s\|_{H^1(U)} + \|q_s^t - (s-t)q'_s\|_{L^\infty(U)}] = o(t-s).$$

In particular, this implies that $u_t(x), \nabla u_t$ are continuously differentiable in t , locally uniformly on Ω .

As for checking that $v_0 \in C(\bar{\Omega})$, note that q'_0 only depends on $u_0, \nabla u_0$, which are bounded uniformly on Ω (as $u_0 = u_\Omega$, and using Corollary 5.5). Thus from elliptic estimates on NTA domains (see e.g. [26, Theorem 1.2.8]), we have that for some $\alpha > 0$,

$$\|v_0^0\|_{C^{0,\alpha}(\bar{\Omega})} \leq C \left[\|v_0^0\|_{H_0^1(\Omega)} + \|q'_0\|_{L^\infty(\Omega)} \right] \leq C.$$

On the other hand, v_0^1 is a bounded multiple of u_0 , and therefore Lipschitz continuous. \square

Lemma 8.4. *Let Ω be a minimizer of \mathcal{F}_τ , $\phi_t : M \rightarrow M$ a one-parameter family of smooth diffeomorphisms with $\phi_0(x) = x$, $\partial_t \phi_t|_{t=0} = T$ a vector field on M , and $\Omega_t = \phi_t(\Omega)$. Then $u_{\Omega_t} : (-c, c) \rightarrow H^1(\Omega)$ is differentiable in t at $t=0$ with derivative $\dot{u}_\Omega = \dot{u}_\Omega^T \in L^\infty(M)$ vanishing outside Ω , in the following sense:*

$$\lim_{t \rightarrow 0} \frac{1}{t} \|u_{\Omega_t} - u_\Omega - t \dot{u}_\Omega\|_{L^2(M)} = 0.$$

The function \dot{u}_Ω satisfies

$$(8.9) \quad \begin{cases} -\Delta \dot{u}_\Omega = \lambda_1(\Omega) \dot{u}_\Omega - u_\Omega \int_{\partial^* \Omega} |\nabla u_\Omega|^2 g(T, \nu) d\mathcal{H}^{n-1} & \text{on } \Omega \\ \dot{u}_\Omega = -g(\nabla u_\Omega, T) & \text{on } \partial^* \Omega \\ \int_\Omega \dot{u}_\Omega u_\Omega = 0. \end{cases}$$

The boundary condition holds in the sense of nontangential limits at a.e. point on $\partial^* \Omega$.

We will write \dot{u}_Ω^T to denote the derivative found in Lemma 8.4 when we wish to emphasize its dependence on the vector field T , and will simply write \dot{u}_Ω otherwise in order to alleviate notation. We will discuss the specifics of the PDE satisfied by \dot{u}_Ω below, but note for now that from the absolute continuity of harmonic and Hausdorff measure from Proposition 7.9, the boundary condition is given on a sufficiently large portion of $\partial\Omega$ (i.e. there is a unique solution to this problem; see [26, Theorem 1.4.4]).

Proof. Set $u_t = u_{\Omega_t} \circ \phi_t : \Omega \rightarrow [0, \infty)$ to be the pullback of u_{Ω_t} . Then u_t satisfies a PDE on Ω ; namely, $-\Delta_t u_t = f_t u_t$, where Δ_t is the Laplace-Beltrami operator associated with the metric $g_t = \phi_t^* g$ and $f_t = \lambda_1(\Omega_t)$. From Lemma 8.2, part (5), f_t is a C^1 function. Moreover, u_t is the unique nonnegative solution to this equation in $H_0^1(\Omega)$ with $\|u_t\|_{L^2(\Omega, dm_t)} = 1$, while the same equation with f_t replaced by λf_t , $\lambda \in [0, 1 + c]$ admits no nontrivial solutions (this follows from Lemma 8.2, part (4)). Apply Lemma 8.3 to this to learn that u_t is continuously differentiable in t and to obtain a function $v \in H_0^1(\Omega) \cap C(\Omega)$ with

$$\lim_{t \rightarrow 0} \frac{1}{t} \|u_t - u_\Omega - tv\|_{H_0^1(\Omega)} = 0.$$

In particular, we have that

$$(8.10) \quad \|u_{\Omega_t} - u_\Omega \circ \phi_t^{-1}\|_{H_0^1(\Omega_t)} \leq C \|u_t - u_\Omega\|_{H_0^1(\Omega)} \leq C|t|.$$

Our goal now is to estimate $u_{\Omega_t} - u_t$. Consider, with t fixed, the function $\psi \circ \phi_s(x) : M \times [0, t] \rightarrow [0, \infty)$ for any $\psi \in H^1(M)$. This function lies in $H^1(M \times [0, t])$, with distributional derivative $d_{(x,s)}(\psi \circ \phi_s)(v, 1) = (d_{\phi_s(x)}\psi(d_x\phi_s(x)v), d_{\phi_s(x)}\psi(\partial_s\phi_s(x)))$; this may be verified by approximating by smooth functions. Plugging in u_{Ω_t} for ψ leads to the identity

$$u_{\Omega_t}(x) - u_t(x) = - \int_0^t \partial_s [u_{\Omega_t} \circ \phi_s](x) ds = - \int_0^t g(\nabla u_{\Omega_t}(\phi_s(x)), \partial_s \phi_s(x)) ds,$$

which is valid for almost every $x \in M$. Now, for $|s| \leq |t|$,

$$\begin{aligned} \|\nabla(u_{\Omega_t} \circ \phi_s) - \nabla u_\Omega\|_{L^2(M)} &\leq \|\nabla(u_{\Omega_t} \circ \phi_s) - \nabla(u_\Omega \circ \phi_t^{-1} \circ \phi_s)\|_{L^2(M)} + \|\nabla u_\Omega - \nabla(u_\Omega \circ \phi_t^{-1} \circ \phi_s)\|_{L^2(M)} \\ &\leq C|t| + o_t(1). \end{aligned}$$

We used (8.10) to bound the first term (after changing variables), while the second goes to 0 from the dominated convergence theorem (∇u_Ω is bounded while $\phi_t^{-1} \circ \phi_s(x) \rightarrow x$ pointwise). We also have that

$$|\partial_s \phi_s(x) - T| \leq C|s| \leq C|t|$$

as ϕ_s is smooth. Applying both of these,

$$\begin{aligned} \frac{1}{|t|} \|u_{\Omega_t} - u_t + tg(\nabla u_\Omega, T)\|_{L^2(M)} &\leq \frac{1}{|t|} \int_0^t \|g(\nabla u_{\Omega_t}(\phi_s(\cdot)), \partial_s \phi_s(\cdot)) - g(\nabla u_\Omega, T)\|_{L^2(M)} ds \\ &\leq C|t| + o_t(1) = o_t(1). \end{aligned}$$

Now, set $\dot{u}_\Omega = v - g(\nabla u_\Omega, T)$ (extending it by 0 outside of Ω). We have shown that

$$\frac{1}{|t|} \|u_{\Omega_t} - u_\Omega - t\dot{u}_\Omega\|_{L^2(M)} \leq \frac{1}{|t|} [\|u_t - u_\Omega - tv\|_{L^2(\Omega)} + \|u_{\Omega_t} - u_t + tg(\nabla u_\Omega, T)\|_{L^2(M)}] \rightarrow 0.$$

As both v and ∇u_Ω are bounded, so is \dot{u}_Ω . It remains to verify that \dot{u}_Ω solves the stated PDE.

First, recall that from Lemma 8.3 we know u_t (and its spatial derivatives) are continuously differentiable in t on Ω for small $|t|$. It follows that u_{Ω_t} is continuously differentiable in t on Ω_t as well, and $\partial_t u_{\Omega_t}|_{t=0} = \dot{u}_\Omega$ on Ω . Therefore, we may differentiate the PDE $-\Delta u_{\Omega_t} = \lambda_1(\Omega_t)u_{\Omega_t}$ pointwise on Ω at $t = 0$ to give

$$-\Delta \dot{u}_\Omega = \lambda_1(\Omega)\dot{u}_\Omega + (\partial_t \lambda_1(\Omega_t)|_{t=0})u_\Omega = \lambda_1(\Omega)\dot{u}_\Omega - u_\Omega \int_{\partial^* \Omega} |\nabla u_\Omega|^2 g(T, \nu) d\mathcal{H}^{n-1}$$

on Ω , where in the second equality we used the final statement of Lemma 8.2. The orthogonality condition may be inferred from the normalization of u_{Ω_t} : indeed,

$$1 = \int u_{\Omega_t}^2 = \int (u_\Omega + t\dot{u}_\Omega)^2 + o(t) = 1 + 2t \int u_\Omega \dot{u}_\Omega + o(t),$$

Rearranging and sending t to 0 gives $\int u_\Omega \dot{u}_\Omega = 0$. Finally, we have that $-g(\nabla u_\Omega, T)$ has nontangential limits a.e. on $\partial\Omega$ from Corollary 7.12, while $v \rightarrow 0$ on $\partial\Omega$ uniformly; this gives the boundary condition. \square

We are now in a position to prove the main result of this subsection, which is a “distributional” Euler-Lagrange equation for Ω .

Lemma 8.5 (Distributional Euler-Lagrange equation). *Let Ω be a minimizer of \mathcal{F}_τ . Fix $B_r(x_0) \subseteq Q_R$ and let $L := |B_r(x_0)|$. Let T be a smooth vector field on M with $|T| \leq 1$ that is compactly supported on $B_r(x_0)$.*

If T is volume preserving to first order in the sense that $\int_{\partial^\Omega} g(T, \nu_x) d\mathcal{H}^{n-1} = 0$, then*

$$\left| - \int_{\partial^*\Omega} |\nabla u_\Omega|^2 g(T, \nu) + \frac{\mathfrak{I}}{2} |\nabla w_\Omega|^2 g(T, \nu_x) d\mathcal{H}^{n-1} + \tau \left[\int_\Omega \dot{u}_\Omega a_\Omega + \int_{\partial^*\Omega} b_\Omega g(T, \nu) d\mathcal{H}^{n-1} \right] \right| \leq CL,$$

where a_Ω, b_Ω are the functions from property (N4) of \mathfrak{h} and $\dot{u}_\Omega = \dot{u}_\Omega^T$ is as in Lemma 8.4.

If $|\Omega| \neq v$, $L < ||\Omega| - v|$, and we do not assume $\int_{\partial^\Omega} g(T, \nu_x) d\mathcal{H}^{n-1} = 0$, we instead have*

$$\begin{aligned} & \left| - \int_{\partial^*\Omega} |\nabla u_\Omega|^2 g(T, \nu) + \frac{\mathfrak{I}}{2} |\nabla w_\Omega|^2 g(T, \nu) d\mathcal{H}^{n-1} \right. \\ & \quad \left. + f'_{v,\eta}(|\Omega|) \int_{\partial^*\Omega} g(T, \nu) d\mathcal{H}^{n-1} + \tau \left[\int_\Omega \dot{u}_\Omega a_\Omega + \int_{\partial^*\Omega} b_\Omega g(T, \nu) d\mathcal{H}^{n-1} \right] \right| \leq CL. \end{aligned}$$

The term containing \dot{u}_Ω may first appear to be lower-order since it is integrated on Ω rather than the boundary $\partial^*\Omega$. This is not the case: due to how $\dot{u}_\Omega = \dot{u}_\Omega^T$ depends on T , it will end up being of the same order as the others.

Proof of Lemma 8.5. Let $\phi_t(x)$ be the flow associated with T : $\partial_t \phi_t(x) = T(\phi_t(x))$; then ϕ_t is a smooth family of diffeomorphisms for all $|t| \leq c$ small, with $\phi_t(x) = x$ outside of $B_r(x_0)$. Setting $\Omega_t = \phi_t(\Omega)$, apply Lemma 8.2 to Ω and ϕ_t : this implies that $\lambda_2(\Omega_t) > \lambda_1(\Omega_t) + c$ and

$$\mathcal{E}(\Omega_t) \leq \mathcal{E}(\Omega) - t \int_{\partial^*\Omega} |\nabla u_\Omega|^2 g(T, \nu) + \frac{\mathfrak{I}}{2} |\nabla w_\Omega|^2 g(T, \nu) d\mathcal{H}^{n-1} + o(t).$$

Let us now consider $\mathfrak{h}(\Omega_t)$: using property (N4), we have that

$$\mathfrak{h}(\Omega_t) - \mathfrak{h}(\Omega) = \int (u_{\Omega_t} - u_\Omega) a_\Omega + \int_{\Omega_t} b_\Omega - \int_\Omega b_\Omega + o(t) + O(tr^n).$$

The second term here may be estimated as in Lemma 8.2 using (8.2) to give

$$\begin{aligned} \int_{\Omega_t} b_\Omega - \int_\Omega b_\Omega &= t \int_\Omega b_\Omega \operatorname{div} T + o(t) = t \int_{\partial^*\Omega} b_\Omega g(T, \nu) d\mathcal{H}^{n-1} - t \int_\Omega g(\nabla b_\Omega, T) + o(t) \\ &= t \int_{\partial^*\Omega} b_\Omega g(T, \nu) d\mathcal{H}^{n-1} + O(tr^n) + o(t), \end{aligned}$$

using that $\|b_\Omega\|_{C^{0,1}} \leq C$ by assumption. For the first term, we see that $\|u_{\Omega_t} - u_\Omega - t\dot{u}_\Omega\|_{L^2} = o(t)$ from Lemma 8.4, and so

$$\int (u_{\Omega_t} - u_\Omega) a_\Omega = t \int \dot{u}_\Omega a_\Omega + o(t).$$

Putting everything together and using that $\mathcal{F}_\tau(\Omega) \leq \mathcal{F}_\tau(\Omega_t)$,

$$0 \leq -t \int_{\partial^*\Omega} |\nabla u_\Omega|^2 g(T, \nu) + \frac{\mathfrak{I}}{2} |\nabla w_\Omega|^2 g(T, \nu) d\mathcal{H}^{n-1} + \tau \left[t \int_{\partial^*\Omega} b_\Omega g(T, \nu) d\mathcal{H}^{n-1} + t \int \dot{u}_\Omega a_\Omega \right] + o(t) + O(tr^n).$$

The conclusion follows from dividing by t and sending $t \rightarrow 0$.

If $|\Omega| \neq v$ and $||\Omega| - |\Omega_t|| < ||\Omega| - v|$ (which always holds if t is small enough) then $f_{v,\eta}$ is linear and we may estimate

$$f_{v,\eta}(|\Omega_t|) = f_{v,\eta}(|\Omega|) + f'_{v,\eta}(|\Omega|)[|\Omega_t| - |\Omega|] = f_{v,\eta}(|\Omega|) + t f'_{v,\eta}(|\Omega|) \int_{\partial^*\Omega} g(T, \nu) d\mathcal{H}^{n-1} + o(t).$$

We then proceed as previously. \square

8.2. The first variation at a point. Next, we wish to derive a pointwise form of the distributional Euler-Lagrange equation of Lemma 8.5 by sending T to $\delta_x \nu_x - \delta_y \nu_y$ at points $x, y \in \partial^* \Omega$. Here and in the sequel, we let $\nu_x = \nu_\Omega(x)$ be the outer unit normal of Ω at x . Passing to this limit is actually quite delicate, with the main challenge arising from making sure that the functions $\dot{u}_\Omega = \dot{u}_\Omega^T$ behave well under such an approximation procedure. Heuristically speaking, recalling the equation (8.9) satisfied by \dot{u}_Ω^T , the limit of the \dot{u}_Ω^T should solve an equation with a Dirac delta as boundary data. On the other hand, the boundary data in (8.9) is achieved only \mathcal{H}^{n-1} -a.e. and in a nontangential limit sense. To rigorously derive this limit, we will use the PDE (8.9) solved by \dot{u}_Ω by writing \dot{u}_Ω in terms of the Green's function and Poisson kernel.

It will be convenient to decouple the Poisson kernel part and the Green's function part, or in other words, the harmonic part and its remainder part, of \dot{u}_Ω in the following way. For any smooth vector field as in Lemma 8.5, let the harmonic part h^T of \dot{u}_Ω^T be given by

$$h^T(x) = \int_{\partial\Omega} \dot{u}_\Omega^T d\omega_x,$$

where ω_x is harmonic measure; i.e. it solves the PDE

$$\begin{cases} -\Delta h^T = 0 & \text{on } \Omega \\ h^T = \dot{u}_\Omega^T = -g(\nabla u_\Omega, T) & \text{on } \partial\Omega \end{cases}$$

with the boundary condition in the sense of nontangential limits at a.e. point. Then, let $q^T = \dot{u}_\Omega^T - h^T$ be the remainder part. Using the equation for \dot{u}_Ω^T from Lemma 8.4, we see that q^T solves the PDE

$$(8.11) \quad \begin{cases} -\Delta q^T = \lambda_1(\Omega)(q^T + h^T) - u_\Omega \int_{\partial^* \Omega} |\nabla u_\Omega|^2 g(T, \nu) d\mathcal{H}^{n-1} & \text{on } \Omega \\ q^T = 0 & \text{on } \partial\Omega \\ \int (q^T + h^T) u_\Omega = 0. \end{cases}$$

We construct our approximating vector fields in the following way. Given a point $x \in \partial^* \Omega$ and scale r sufficiently small, let $T_{x,r}$ be a vector field chosen so that $|T_{x,r}| \leq 1$, $\text{supp } T_{x,r} \subseteq B_{c(n)r}(x)$, and $\int_{\partial^* \Omega} g(T_{x,r}, \nu) d\mathcal{H}^{n-1} = \omega_{n-1} r^{n-1}$ where ω_{n-1} is the volume of the unit ball in \mathbb{R}^{n-1} ; this may always be found by starting with the vector field $\nu_x \kappa_r$ in normal coordinates near x , with κ_r a smooth cutoff function approximating $1_{B_r(x)}$, and then multiplying by a constant.

Below, recall the notation

$$K(x, y) = \frac{d\omega_y}{d\mathcal{H}^{n-1} \llcorner \partial\Omega}(x)$$

for the Poisson kernel, i.e. the Radon-Nikodym derivative of harmonic and surface measures, from (7.12).

First, in the following lemma we will show that the harmonic parts $h^{T_{x,r}}$, after suitable renormalization, converge to a multiple of the Poisson kernel as $r \rightarrow 0$. The main tool in the proof is the sharp estimates for the harmonic measure established in Proposition 7.9.

Lemma 8.6. *For every $x \in \partial^* \Omega$, if $\text{supp } T \subseteq B_r(x)$ then the function h^T satisfies*

$$|h^T(y)| \leq \frac{C}{d^{n-1}(y, x)} \int_{\partial^* \Omega} |T| d\mathcal{H}^{n-1}$$

whenever $d(x, y) \geq Cr$.

For \mathcal{H}^{n-1} -almost every $x \in \partial^* \Omega$, if $h^x(y) = |\nabla u_\Omega(x)| K(x, y)$ is a multiple of the Poisson kernel, then h^x is well-defined, satisfies

$$|h^x(y)| \leq \frac{C}{d^{n-1}(y, x)},$$

and

$$(8.12) \quad \lim_{r \searrow 0} \sup_{y \in \Omega \setminus B_{Cr}(x)} \left| h^x(y) - \frac{h^{T_{x,r}}(y)}{\omega_{n-1} r^{n-1}} \right| d^{n-1}(y, x) = 0.$$

Proof. The first conclusion follows from the harmonic measure estimate (7.8) in Proposition 7.9, along with the fact that $|\nabla u_\Omega| \leq C$ from Lemma 5.5. For the others, select x with the following properties: (1) $x \in \partial^*\Omega$, (2) x is a Lebesgue point of $|\nabla u|$ with respect to \mathcal{H}^{n-1} , (3) x is a Lebesgue point of $K(x, y)$ with respect to \mathcal{H}^{n-1} for a fixed $y \in \Omega$ with $B_{c_0}(y) \subseteq \Omega$. Then the Poisson kernel estimates of (7.12) following Proposition 7.9 directly give that for any $z \in \Omega$,

$$|h^x(z)| = |\nabla u(x)|K(x, z) \leq \frac{C}{d^{n-1}(z, x)}.$$

As a first step toward proving (8.12), we have that for this one fixed y ,

$$\begin{aligned} \left| h^x(y) - \frac{h^{T_{x,r}}(y)}{\omega_{n-1}r^{n-1}} \right| &= \left| |\nabla u(x)|K(x, y) - \frac{h^{T_{x,r}}(y)}{\omega_{n-1}r^{n-1}} \right| \\ &= \frac{1}{\omega_{n-1}r^{n-1}} \left| \int_{\partial^*\Omega \cap B_{cr}(x)} |\nabla u(x)|K(x, y)g(T_{x,r}(z), \nu_z) d\mathcal{H}^{n-1}(z) - h^{T_{x,r}}(y) \right| \\ &= \frac{1}{\omega_{n-1}r^{n-1}} \left| \int_{\partial^*\Omega \cap B_{cr}(x)} |\nabla u(x)|K(x, y)g(T_{x,r}(z), \nu_z) - \dot{u}_\Omega^{T_{x,r}}(z)K(z, y) d\mathcal{H}^{n-1}(z) \right| \\ &\leq \frac{C}{r^{n-1}} \int_{\partial^*\Omega \cap B_{cr}(x)} \left| |\nabla u(x)|K(x, y)g(T_{x,r}(z), \nu_z) + g(\nabla u_\Omega(z), T_{x,r}(z))K(z, y) \right| d\mathcal{H}^{n-1}(z) \\ &= \frac{C}{r^{n-1}} \int_{\partial^*\Omega \cap B_{cr}(x)} \left| |\nabla u(z)|K(x, y)g(T_{x,r}(z), \nu_z) + g(\nabla u_\Omega(z), T_{x,r}(z))K(z, y) \right| d\mathcal{H}^{n-1}(z) + o_r(1) \\ &= \frac{C}{r^{n-1}} \int_{\partial^*\Omega \cap B_{cr}(x)} \left| K(z, y)g(T_{x,r}(z), \nabla u_\Omega(z)) - g(\nabla u_\Omega(z), T_{x,r}(z))K(z, y) \right| d\mathcal{H}^{n-1}(z) + o_r(1) \\ &= o_r(1), \end{aligned}$$

using the definitions, both of the Lebesgue point assumptions, and the fact that $K(z, y), |\nabla u_\Omega(z)|, |T_{x,r}(z)|$ are all bounded.

Now, let us estimate the same quantity replacing y with any $z \in \Omega \setminus B_{Cr}(x)$. To this end, we use the triangle inequality and the definitions of h^x and $h^{T_{r,x}}$ to write

$$\left| h^x(z) - \frac{h^{T_{x,r}}(z)}{\omega_{n-1}r^{n-1}} \right| = \left| h^x(y) \frac{K(x, z)}{K(x, y)} - \frac{h^{T_{x,r}}(z)}{\omega_{n-1}r^{n-1}} \right| \leq I + II,$$

where we set

$$\begin{aligned} I &= \left| h^x(y) - \frac{h^{T_{x,r}}(y)}{\omega_{n-1}r^{n-1}} \right| \frac{K(x, z)}{K(x, y)} \\ II &= \frac{C}{r^{n-1}} \left| \int_{\partial^*\Omega \cap B_{cr}(x)} \dot{u}_\Omega^{T_{x,r}}(q) \left[K(q, z) - K(q, y) \frac{K(x, z)}{K(x, y)} \right] d\mathcal{H}^{n-1}(q) \right| \end{aligned}$$

Since $d(x, y) \geq c_0$, from (7.12) and Proposition 7.9 we have

$$\frac{K(x, z)}{K(x, y)} \leq CK(x, z) \leq \frac{C}{d^{n-1}(z, x)}$$

and therefore we see that $I \leq o_r(1)/d^{n-1}(z, x)$ using the first step above. Next, from [26, Corollary 1.3.20],

$$\left| \frac{d\omega_z}{d\omega_y}(x) - \frac{d\omega_z}{d\omega_y}(q) \right| \leq C \left| \frac{d\omega_z}{d\omega_y}(x) \right| d^\alpha(x, q) \leq C \frac{d^\alpha(x, q)}{d^{n-1}(z, x)}$$

for some $\alpha \in (0, 1)$, giving

$$\left| K(q, z) - K(q, y) \frac{K(x, z)}{K(x, y)} \right| \leq C \frac{d^\alpha(x, q)}{d^{n-1}(z, x)} K(q, y) \leq C \frac{d^\alpha(x, q)}{d^{n-1}(z, x)}$$

for any $q \in \partial\Omega$. We therefore see that $II \leq Cr^\alpha/d^{n-1}(z, x)$. This proves (8.12). \square

We now move toward estimating the remainder q^T and pass to a limit for the (renormalized) functions $q^{T_{x,r}}$ corresponding to the vector fields $T_{x,r}$ defined above. The remainder q^T satisfies a simpler boundary condition than h^T (recall (8.11)), and we will use the following auxiliary elliptic estimate to help control it.

Lemma 8.7. *Let p be a bounded function on Ω with $\int p u_\Omega = 0$. Then*

$$\begin{cases} -\Delta u = \lambda_1(\Omega)u + p & \text{on } \Omega \\ \int u u_\Omega = 0 \\ u = 0 & \text{on } \partial\Omega \end{cases}$$

admits a unique solution $u \in H_0^1(\Omega)$, which has the estimate

$$(8.13) \quad \|u\|_{L^P} \leq C(P)\|p\|_{L^1}$$

for any $P \in [1, \frac{n}{n-2})$.

The existence and uniqueness are immediate from the Fredholm alternative (using that $p \in L^\infty$), as in (8.6). The main point is the estimate in terms of only the L^1 norm of p .

Sketch of proof. The estimate

$$\|u\|_{L^\infty(\Omega)} \leq C[\|p\|_{L^{P'}(\Omega)} + \|u\|_{L^2(\Omega)}]$$

for any $P' > n/2$ may be found in [24, Theorem 8.15]. Combining with the Hilbert space estimate from the Fredholm alternative (as in (8.6))

$$\|u\|_{L^2} \leq \|u\|_{H_0^1(\Omega)} \leq C\|p\|_{L^2(\Omega)}$$

gives

$$\|u\|_{L^\infty(\Omega)} \leq C\|p\|_{L^{P'}(\Omega)}.$$

Then (8.13) follows from duality. \square

Corollary 8.8. *For almost every $x \in \partial^*\Omega$, the functions $q^{T_{x,r}}/\omega_{n-1}r^{n-1}$ converge in L^p topology (for any $p \in [1, \frac{n}{n-2})$) to a function q^x , which solves*

$$\begin{cases} -\Delta q^x = \lambda_1(\Omega)(q^x + h^x) - u_\Omega |\nabla u_\Omega(x)|^2 & \text{on } \Omega \\ \int_\Omega (q^x + h^x) u_\Omega = 0 \\ q^x = 0 & \text{on } \partial\Omega, \end{cases}$$

in the sense that q^x is the Green potential of the right-hand side: for a.e. $y \in \Omega$,

$$(8.14) \quad q^x(y) = \int_\Omega [\lambda_1(\Omega)(q^x(z) + h^x(z)) - u_\Omega(z)|\nabla u_\Omega(x)|^2] G_\Omega(z, y) dz.$$

It is straightforward to check that in fact q^x is continuous (up to the boundary) away from x , and satisfies the PDE and boundary condition classically. However, the statement here is actually stronger: it implies that at x , q^x still is 0 in the sense of that it does not form measure-valued boundary data there.

Proof. Recalling the equation (8.11) satisfied by q^T , we wish to further decompose q^T further as $q^T = q_1^T - u_\Omega \int h^T u_\Omega$ in order to apply the estimate of Lemma 8.7 above to the piece q_1^T , which is orthogonal to u_Ω by construction. Let us integrate by parts twice to compute

$$(8.15) \quad \begin{aligned} \lambda_1(\Omega) \int_\Omega h^T u_\Omega &= - \int_\Omega h^T \Delta u_\Omega \\ &= \int_\Omega g(\nabla h^T, \nabla u_\Omega) - \int_{\partial^*\Omega} h^T g(\nabla u_\Omega, \nu) d\mathcal{H}^{n-1} = - \int_{\partial^*\Omega} |\nabla u_\Omega|^2 g(T, \nu) d\mathcal{H}^{n-1}. \end{aligned}$$

We used the divergence theorem of [16], the fact that $\Delta h^T = 0$, and Corollary 7.12. In particular, if we write $h_1^T = h^T - u_\Omega \int h^T u_\Omega$ (so that $h_1^T + q_1^T = h^T + q^T$), we have that

$$\begin{cases} -\Delta q_1^T = \lambda_1(\Omega)(q_1^T + h_1^T) & \text{on } \Omega \\ \int q_1^T u_\Omega = 0 \\ q_1^T = 0 & \text{on } \partial\Omega \end{cases}$$

and $\int h_1^T u_\Omega = 0$.

We claim that $h_1^{T_{x,r}}/\omega_{n-1}r^{n-1}$ is a Cauchy sequence in L^1 (with respect to index r). To see this, apply Lemma 8.6, we know that $\lim_{r \searrow 0} \|\frac{h^{T_{x,r}}}{\omega_{n-1}r^{n-1}} - h^x\|_{L^1} = 0$, and that

$$\lim_{r \searrow 0} \int u_\Omega \frac{h^{T_{x,r}}}{\omega_{n-1}r^{n-1}} = \int h^x u_\Omega.$$

From (8.15) we see this implies that at each Lebesgue point x of $|\nabla u_\Omega|$ on $\partial^* \Omega$ we have

$$\int h^x u_\Omega = \lim_{r \searrow 0} -\frac{1}{\lambda_1(\Omega)\omega_{n-1}r^{n-1}} \int_{\partial^* \Omega} |\nabla u_\Omega|^2 g(T_{x,r}, \nu_x) d\mathcal{H}^{n-1} = -\frac{|\nabla u_\Omega(x)|^2}{\lambda_1(\Omega)}.$$

So, applying Lemma 8.7, we see that $q_1^{T_{x,r}}/\omega_{n-1}r^{n-1}$ is Cauchy in L^p , and so converges to some function q_1^x . Passing the orthogonality condition to the limit, $\int q_1^x u_\Omega = 0$. It then also follows that $q^{T_{x,r}}/\omega_{n-1}r^{n-1} \rightarrow q^x := q_1^x - u_\Omega \int h^x u_\Omega$.

We have verified the orthogonality condition $\int_\Omega (q^x + h^x) u_\Omega = 0$. The boundary condition and PDE may be checked by writing

$$q^{T_{x,r}}(y) = \int_\Omega \left[\lambda_1(\Omega) (q^{T_{x,r}}(z) + h^{T_{x,r}}(z)) - u_\Omega(z) \int_{\partial^* \Omega} |\nabla u_\Omega(y')|^2 g(T_{x,r}, \nu_{y'}) d\mathcal{H}^{n-1}(y') \right] G_\Omega(z, y) dz,$$

and passing both sides to the limit in L^p , which results in

$$q^x(y) = \int_\Omega [\lambda_1(\Omega)(q^x(z) + h^x(z)) - u_\Omega(z)|\nabla u_\Omega(x)|^2] G_\Omega(z, y) dz$$

for a.e. $y \in \Omega$. □

Remark 8.9. We may recover further estimates for q^x from this argument if we wish. For example, for any $\alpha, \beta \in \mathbb{R}$,

$$\|\alpha q^x + \beta q^y\|_{L^1} \leq C\|\alpha h^x + \beta h^y\|_{L^1} + C|\alpha|\nabla u_\Omega(x)|^2 + \beta|\nabla u_\Omega(y)|^2,$$

by taking the limit in the corresponding estimates for q^T with $T = \alpha T_{x,r} + \beta T_{y,r}$. If we choose $\alpha = \frac{1}{|\nabla u_\Omega(x)|^2}$ and $\beta = -\frac{1}{|\nabla u_\Omega(y)|^2}$, the second term vanishes and we get

$$\left\| \frac{q^x}{|\nabla u_\Omega(x)|^2} - \frac{q^y}{|\nabla u_\Omega(y)|^2} \right\|_{L^1} \leq C \left\| \frac{h^x}{|\nabla u_\Omega(x)|^2} - \frac{h^y}{|\nabla u_\Omega(y)|^2} \right\|_{L^1}.$$

These may also be derived directly from the potential identity (8.14).

Set $v_\Omega^x = q^x + h^x$. We are now in a position to derive a pointwise Euler-Lagrange equation, or free boundary condition, satisfied \mathcal{H}^{n-1} -a.e. along $\partial\Omega$.

Corollary 8.10 (Pointwise form of the Euler-Lagrange equation). *Let Ω be a minimizer. Then there is a constant A_0 such that for almost every $x \in \partial^* \Omega \cap Q_R$, we have the identity*

$$(8.16) \quad -|\nabla u_\Omega(x)|^2 - \frac{\mathfrak{T}}{2} |\nabla w_\Omega(x)|^2 + \tau \left[b_\Omega(x) + \int_\Omega a_\Omega v_\Omega^x \right] = -A_0.$$

Here $v_\Omega^x = q^x + h^x$ where q^x is as in (8.14) and $h^x = |\nabla u_\Omega(x)|K(x, \cdot)$. If $|\Omega| \neq v$, then $A_0 = f'_{v,\eta}(|\Omega|)$.

Proof of Corollary 8.10. Take x, y any two points in $\partial^* \Omega \cap Q_R$ to which both Lemma 8.6 and Corollary 8.8 apply. Using $T_r = (T_{x,r} - T_{y,r})/\omega_{n-1}r^{n-1}$ for T in Lemma 8.5, we see that

$$\left| \int_{\partial^* \Omega} \left[-|\nabla u_\Omega(q)|^2 - \frac{\mathfrak{T}}{2} |\nabla w_\Omega(q)|^2 + \tau b_\Omega(q) \right] g(T_r(q), \nu_q) d\mathcal{H}^{n-1}(q) + \tau \int_\Omega a_\Omega \dot{u}_\Omega^{T_r} \right| \leq Cr \rightarrow 0.$$

Restricting further to only those x, y which are Lebesgue points of both $|\nabla u_\Omega|$ and $|\nabla w_\Omega|$ with respect to $\mathcal{H}^{n-1} \llcorner \partial\Omega$, we see that the first term converges to

$$|\nabla u_\Omega(y)|^2 - |\nabla u_\Omega(x)|^2 + \frac{\mathfrak{T}}{2} [|\nabla w_\Omega(y)|^2 - |\nabla w_\Omega(x)|^2] + \tau [b_\Omega(x) - b_\Omega(y)].$$

From Lemma 8.6, we have that $h^{T_r} \rightarrow h^y - h^x$ in $L^1(\Omega)$, while from Corollary 8.8, $q^{T_r} \rightarrow q^y - q^x$ in $L^1(\Omega)$. Passing to the limit then gives

$$|\nabla u_\Omega(y)|^2 - |\nabla u_\Omega(x)|^2 + \frac{\tau}{2} [|\nabla w_\Omega(y)|^2 - |\nabla w_\Omega(x)|^2] + \tau [b_\Omega(x) - b_\Omega(y)] + \tau \int_\Omega a_\Omega [v_\Omega^x - v_\Omega^y] = 0.$$

This implies the quantity in (8.16) is independent of x .

If $|\Omega| \neq v$, we may instead just use $T_r = T_{x,r}/\omega_{n-1}r^{n-1}$; proceeding similarly gives that

$$-|\nabla u_\Omega(x)|^2 - \frac{\tau}{2} |\nabla w_\Omega(x)|^2 + f'_{v,\eta}(|\Omega|) + \tau \left[b_\Omega(x) + \int_\Omega a_\Omega v_\Omega^x \right] = 0.$$

This completes the proof. \square

8.3. The Euler-Lagrange equation rewritten. As we discussed at the beginning of the section, the pointwise form of the Euler-Lagrange equation in Corollary 8.10 is not particularly useful from the perspective of regularity theory. The goal of this subsection is to prove Theorem 8.1, that is, to rewrite (8.16) in a way which involves only $|\nabla u_\Omega|$, up to “lower order” terms.

The most difficult term in (8.16) to control is $\int_\Omega a_\Omega v_\Omega^x$, so let us give a brief heuristic idea of how this is done. Recall that $v_\Omega^x = h^x + q^x$; we will consider these two terms separately. The term $h^x = |\nabla u_\Omega|K(x, \cdot)$ is strictly harder to control, so let us focus on the term $\int_\Omega a_\Omega h^x$. In order to compare this term to $|\nabla u_\Omega|^2$, we reinterpret the equation solved by h^x using that the Poisson kernel is the (normal) derivative of the Green’s function. So,

$$\int_\Omega h^x(y) dm(y) \leq |\nabla u_\Omega(x)| \left| \nabla_x \int_\Omega G_\Omega(x, y) a_\Omega(y) dm(y) \right| = |\nabla u_\Omega| |\nabla p(x)|,$$

where $p(x)$ is the potential solving $\Delta p = a_\Omega$ with $p = 0$ on $\partial\Omega$. From here, this term can be handled in a similar way to the term $|\nabla w_\Omega|^2$, using several applications of the inhomogeneous boundary Harnack principle.

Let us now move toward proving Theorem 8.1. Before proving the main theorem, we will need several lemmas that will allow us to control the other terms in (8.16). The first lemma shows that the ratio between the first eigenfunction and a function p that solves an equation is Hölder continuous up to the boundary. One application of this lemma will be to $p = w_\Omega$ in the proof of Theorem 8.1. It will also be used to help control the bad term $\int_\Omega a_\Omega v_\Omega^x$ in the Euler-Lagrange equation (8.16).

Lemma 8.11. *Let p be a continuous function with $p = 0$ on $\partial\Omega$ and $-\Delta p = f$, $\|f\|_{L^\infty} \leq 1$. Then*

$$(8.17) \quad \left\| \frac{p}{u_\Omega} \right\|_{C^{0,\alpha}(\bar{\Omega})} \leq C$$

for some C, α depending only on v, v_{max}, R, η . Moreover, ∇p admits nontangential limits at almost every point x of $\partial^*\Omega$, $|g(\nabla p(x), \nu_x)| = |\nabla p(x)|$, and

$$\frac{|\nabla p(x)|}{|\nabla u_\Omega(x)|} = \lim_{y \rightarrow x \text{ n.t.}} \left| \frac{p(y)}{u_\Omega(x)} \right|.$$

In particular, $|\nabla p(x)| = |\nabla u_\Omega(x)| \rho_p(x)$ along $\partial\Omega$, where $\|\rho_p(x)\|_{C^{0,\alpha}(\partial\Omega)} \leq C$.

When one instead has the ratio of two harmonic functions, this fact follows from a standard argument iteratively applying the classical boundary Harnack inequality. The argument here is essentially similar, but we include the details since our functions of interest solve equations with a right-hand side and require the form of the boundary Harnack estimate established in [4].

Proof. First, let p^+ solve

$$\begin{cases} -\Delta p^+ = f_+ & \text{on } \Omega \\ p^+ = 0 & \text{on } \partial\Omega. \end{cases}$$

We will show that p^+ satisfies (8.17); then by also applying to $-p^-$, we obtain the statement for any p . Note that $0 \leq p^+ \leq w_\Omega$ from the maximum principle, so in particular from Proposition 7.8, $p/u_\Omega \leq C$.

Fix $r < r_0$ small and $z \in \partial\Omega$, and let

$$M(r) = \sup_{\{y \in B_r(z) \cap \Omega : d(y,z) > tr\}} \frac{p^+}{u_\Omega} \quad m(r) = \inf_{\{y \in B_r(z) \cap \Omega : d(y,z) > tr\}} \frac{p^+}{u_\Omega}$$

where t is small and fixed. We claim that there are $c, \varepsilon > 0$ and a $\theta < 1$ such that

$$M(cr) - m(cr) \leq \theta[M(r) - m(r)] + r^\varepsilon.$$

This claim implies (8.17): indeed, a classic iteration argument then gives that $M(r) - m(r) \leq Cr^{\varepsilon'}$, and then applying this along a Harnack chain connecting any two points will give the conclusion.

Our claim is immediate if $M(r) \leq r^\varepsilon$. If not, then there is a point $y \in B_r(z) \cap \Omega$ with $d(y, z) \geq tr$ and $p^+(y) \geq r^\varepsilon u_\Omega(y) \geq cr^{\varepsilon+1}$. From the Harnack inequality, $p^+ \geq cr^{\varepsilon+1} - Cr^2 \geq cr^{\varepsilon+1}$ on $B_{tr/2}(y)$ as long as r_0 was taken small enough. Then we have

$$p^+(y') \geq cr^{1+\varepsilon} r^{n-2} G_\Omega(y, y') \geq cr^{n-1+\varepsilon} G_\Omega(y, y')$$

for y' along $\partial B_{tr/2}(y)$, and by comparison principle on the entire compliment of this ball; we used (7.3) here. Applying the Green's function bounds from Proposition 7.9 gives that for $x \in B_{cr}(z)$,

$$p^+(x) \geq cr^{n-1+\varepsilon} G(y, x) \geq cr^\varepsilon d(x, \partial\Omega) \geq cd^{1+\varepsilon}(x, \partial\Omega).$$

We now apply the inhomogeneous boundary Harnack principle [4, Theorem 1.3] (working in normal coordinates and scaling suitably) on this region $B_{cr}(z)$ to

$$M(r) - \frac{p^+}{u_\Omega} = \frac{M(r)u_\Omega + p^+}{u_\Omega}$$

and

$$\frac{p^+}{u_\Omega} - m(r) = \frac{p^+ - m(r)u_\Omega}{u_\Omega}.$$

Both of these are quotients of positive functions vanishing on $\partial\Omega$ and having bounded Laplace-Beltrami operator, and they both satisfy the growth condition with exponent $\beta = 1 + \varepsilon$. This implies that

$$M(r) - m(cr) \leq C[M(r) - M(cr)] \quad M(cr) - m(r) \leq C[m(cr) - m(r)].$$

These may be rewritten to give $M(cr) - m(cr) \leq \frac{C-1}{C+1}[M(r) - m(r)]$, proving our claim.

Now we consider the remaining conclusions. As $|p(x)| \leq Cu_\Omega(x) \leq Cd(x, \partial\Omega)$, it follows from applying elliptic estimates on $B_{d(x, \partial\Omega)/2}(x)$ that $|\nabla p| \leq C$. Let $\rho_p = p/u_\Omega \in C^{0,\alpha}$; then near any $x \in \partial^*\Omega$ we have from Lemma 7.10 that

$$p(y) = \rho_p(y)u_\Omega(y) = \rho_p(x)[|\nabla u_\Omega(x)|l_{x,\nu_x} + o(|x-y|)].$$

Here l_{x,ν_x} is a truncated linear function in normal coordinates as defined in (7.13). Applying elliptic estimates as in Corollary 7.12, we obtain that $\nabla p(y) \rightarrow -\rho_p(x)|\nabla u_\Omega(x)|\nu_x$ nontangentially and the remaining conclusions follow. \square

We focus our attention toward the term $\int_\Omega a_\Omega v_\Omega^x$. Following the discussion above, a key point will be to rewrite the equation solved by v_Ω^x using that the Poisson kernel is the derivative of the Green's function:

Lemma 8.12. *The derivative of the Green's function $\nabla_x G_\Omega(x, y)$ admits nontangential limits as $x \rightarrow z \in \partial\Omega$ for every y at \mathcal{H}^{n-1} -a.e. $z \in \partial^*\Omega$. Moreover, $\nabla_x G_\Omega(x, y) = -\nu_x |\nabla_x G_\Omega(x, y)|$ and $K(x, y) = |\nabla_x G_\Omega(x, y)|$ a.e. on $\partial^*\Omega$.*

Sketch of proof. The existence of nontangential limits and the expression $\nabla_x G_\Omega(x, y) = -\nu_x |\nabla_x G_\Omega(x, y)|$ may be obtained as in Lemma 8.11. The last point follows from the integration by parts formula

$$\int_{\partial^*\Omega} g(\nabla_x G_\Omega(x, y), \nu_x) \phi(x) d\mathcal{H}^{n-1}(x) = -\phi(y) - \int_\Omega G_\Omega(x, y) \Delta \phi(x) dx$$

for any ϕ smooth (this may be justified using the divergence theorem of [16]). This gives a representation formula for every harmonic function on Ω with continuous boundary conditions, so it follows from the definition of harmonic measure that $K(x, y) = |\nabla_x G_\Omega(x, y)|$ at \mathcal{H}^{n-1} -a.e. x . \square

The two lemmas will be used to control the term $\int_{\Omega} a_{\Omega} v_{\Omega}^x$ in (8.16). Recall from the previous section that $v_{\Omega}^x = h^x + q^x$. We focus on the term h^x first.

Lemma 8.13. *Let $f : \Omega \rightarrow \mathbb{R}$ be a function with $|f| \leq 1$. Then at a.e. $x \in \partial\Omega$,*

$$\int_{\Omega} h^x f = |\nabla u_{\Omega}(x)|^2 \rho_f(x),$$

where $\|\rho_f\|_{C^{0,\alpha}} \leq C$ and C, α depend only on v, v_{max}, η, R . As a consequence,

$$\left\| \frac{h^x}{|\nabla u_{\Omega}(x)|^2} - \frac{h^y}{|\nabla u_{\Omega}(y)|^2} \right\|_{L^1} \leq C d^{\alpha}(x, y).$$

Proof. From the definition of h^x and Lemma 8.12, at a.e. $x \in \partial^* \Omega$ we have

$$\int_{\Omega} h^x f = |\nabla u_{\Omega}(x)| \int_{\Omega} K(x, y) f(y) dy = -|\nabla u_{\Omega}(x)| \int_{\Omega} g(\nabla_x G_{\Omega}(x, y), \nu_x) f(y) dy.$$

Let $x_k \in \Omega$ be a sequence of points converging to x nontangentially (i.e. $\varepsilon d(x_k, x) \leq d(x_k, \partial\Omega)$): then $\nabla_x G_{\Omega}(x_k, y)$ converges to $\nabla_x G_{\Omega}(x, y)$ for every y . From [26, Lemma 1.2.8(iv)], $\sup_x \|\nabla_x G_{\Omega}(x, y)\|_{L^P(dy)} \leq C$ for $P \in [1, \frac{n}{n-1})$. Working in normal coordinates and applying the Vitali convergence theorem, this means that

$$\int_{\Omega} g(\nabla_x G_{\Omega}(x, y), \nu_x) f(y) dy = \lim_k \int_{\Omega} g(\nabla_x G_{\Omega}(x_k, y), \nu_x) f(y) dy.$$

Set $p(z) = \int_{\Omega} G_{\Omega}(z, y) f(y) dy$, which solves

$$\begin{cases} -\Delta p = f & \text{on } \Omega \\ p = 0 & \text{on } \partial\Omega. \end{cases}$$

From Lemma 8.11, ∇p admits nontangential limits \mathcal{H}^{n-1} -a.e. on $\partial\Omega$ and $|\nabla p(x)| = |\nabla u_{\Omega}(x)| \rho_f(x)$, with $\|\rho_f\|_{C^{0,\alpha}} \leq C$. In particular,

$$g(\nabla p(x), \nu_x) = \lim_k g(\nabla p(x_k), \nu_x) = \lim_k g(\nabla_x \int_{\Omega} G_{\Omega}(x_k, y) f(y) dy, \nu_x).$$

We may pass the derivative under the integral sign (again using that $\sup_x \|\nabla_x G_{\Omega}(x, y)\|_{L^P(dy)} \leq C$ to justify this), to obtain

$$|\nabla u_{\Omega}(x)| \rho_f(x) = \lim_k \left| \int_{\Omega} g(\nabla_x G_{\Omega}(x_k, y), \nu_x) f(y) dy \right| = \frac{1}{|\nabla u_{\Omega}(x)|} \left| \int_{\Omega} h^x f \right|.$$

The first conclusion of the lemma follows using ρ_f sign $\int_{\Omega} h^x f$ as the ρ_f . The L^1 estimate is now immediate from duality. \square

Next, the term involving q^x is much easier to handle.

Lemma 8.14. *Let $f : \Omega \rightarrow \mathbb{R}$ be a function with $|f| \leq 1$. Then at a.e. $x \in \partial\Omega$,*

$$\int_{\Omega} q^x f = |\nabla u(x)|^2 \rho_f(x),$$

where $\|\rho_f\|_{C^{0,\alpha}} \leq C$ and C, α depend only on R, v, v_{max}, η .

Proof. We apply Lemma 8.13 to Remark 8.9 to obtain

$$\left\| \frac{q^x}{|\nabla u(x)|^2} - \frac{q^y}{|\nabla u(y)|^2} \right\|_{L^1} \leq C \left\| \frac{h^x}{|\nabla u(x)|^2} - \frac{h^y}{|\nabla u(y)|^2} \right\|_{L^1} \leq C d^{\alpha}(x, y).$$

\square

We are now in a position to prove Theorem 8.1.

Proof of Theorem 8.1. Apply Lemmas 8.13 and 8.14 to h^x, q^x with $f = a_\Omega$ to obtain that for almost every $x \in \partial^* \mathcal{H}^{n-1}$,

$$\int_{\Omega} a_\Omega(x) v_\Omega^x = |\nabla u(x)|^2 \rho_1(x),$$

where $\|\rho_1\|_{C^{0,\alpha}(\partial\Omega)} \leq C$. Then apply Lemma 8.11 to $p = w_\Omega$ to give

$$|\nabla w_\Omega(x)|^2 = |\nabla u_\Omega(x)|^2 \rho_2(x)$$

where $\|\rho_2\|_{C^{0,\alpha}(\partial\Omega)} \leq C$. Inserting both into Corollary 8.10 gives

$$|\nabla u_\Omega(x)|^2 \left(-1 - \frac{\tau}{2} \rho_2(x) + \tau \rho_1(x) \right) = -A_0 - \tau b_\Omega(x).$$

So long as τ_0 is small enough, $-1 - \frac{\tau}{2} \rho_2(x) + \tau \rho_1(x) \leq -\frac{1}{2}$. From Lemma 7.10, we have that $|\nabla u_\Omega| \in [1/C, C]$ a.e., so A_0 must also be bounded above and below. Choosing τ_0 smaller as necessary, we may rewrite

$$A_0 = |\nabla u_\Omega(x)|^2 \frac{1 + \frac{\tau}{2} \rho_2(x) - \tau \rho_1(x)}{1 - \frac{\tau}{A_0} b_\Omega(x)} = |\nabla u_\Omega(x)|^2 (1 + \rho(x)),$$

where $\|\rho\|_{C^{0,\alpha}(\partial\Omega)} \leq C$ and $-C\tau \leq \rho$. \square

9. VISCOSITY FORM OF THE EULER-LAGRANGE EQUATION AND $C^{1,\alpha}$ ESTIMATES

Thus far, we have derived the Euler-Lagrange equation satisfied by minimizers of the main energy in a pointwise (\mathcal{H}^{n-1} a.e.) sense and we established Theorem 8.1 to express the Euler-Lagrange equation in a more useful form. In this section, we reformulate the (useful form of the) Euler-Lagrange equation in the viscosity sense. This will allow us to apply known regularity results for one-phase free boundary problems. In particular, when minimizers of the base energy are sufficiently regular (for instance, this is the case when M is a simply connected space form), we find in Theorem 9.4 below that any minimizer of the main energy is a small $C^{1,\alpha}$ perturbation a minimizer of the base energy. We additionally use the viscosity form of the Euler-Lagrange equation to guarantee that minimizers satisfy the volume constraint $|\Omega| = v$.

The viscosity form of the Euler-Lagrange follows directly from the pointwise version and generic free boundary arguments, and we only briefly sketch the proof. Details may be found in [30].

Lemma 9.1. *Theorem 8.1 holds in the viscosity sense: let Ω be a minimizer of \mathcal{F}_τ , $x \in \partial\Omega \cap Q_R$, and ϕ a smooth function on $B_r(x)$. Assume that $\phi_+ \leq (\geq) u_\Omega$ on $B_r(x)$ and $\phi(x) = 0$. Then*

$$|\nabla \phi(x)|^2 (1 + \rho(x)) \leq (\geq) A_0.$$

Sketch of proof. We work in normal coordinates around x . Let $u_r(y) = u_\Omega(ry)/r$ be rescalings, which converge (along a subsequence) locally uniformly on \mathbb{R}^n to a function u . Similarly rescaling $\phi_r(y) = \frac{\phi(ry)}{r}$, this converges to the linear function $\phi_\infty(y) = \nabla \phi(0) \cdot y$. We have that $-\Delta u_r \rightarrow -\Delta u$ in the sense of distributions, so in particular as measures; it follows that u is harmonic on $\{u > 0\}$. As $u_r \geq (\leq) (\phi_r)_+$, this passes to the limit to give $u \geq (\leq) \phi_\infty$. In particular, $\{u > 0\}$ contains (is contained in) the half-space $\{y : \nabla \phi(0) \cdot y > 0\}$. We also recover that $cs \leq \sup_{B_s(0)} u \leq Cs$ in the limit.

Possibly choosing a further subsequence, it may be shown that u is a half-linear function $\alpha(y \cdot \nu)_+$, with $\nu = \nabla \phi / |\nabla \phi|$. See [13, Lemma 11.17] or [30, Lemma 3.6] for details.

As measures, we have that $-\Delta u_r = |\nabla u_r| d\mathcal{H}^{n-1} \llcorner \partial\Omega + o_r(1)$, and $|\nabla u_r(y)| = \sqrt{A_0/(1 + \rho(ry))}$ a.e. on $\partial\Omega$. Passing to the limit in the sense of distributions shows that

$$\int_{\partial^* \Omega/r} \psi |\nabla u_r(y)| d\mathcal{H}^{n-1} = - \int \psi d\Delta u_r + o_1(r) \rightarrow - \int \psi d\Delta u = \int_{y \cdot \nu = 0} \alpha \psi d\mathcal{H}^{n-1}$$

for any smooth ψ . The left-most integral converges to

$$\sqrt{A_0/(1 + \rho(0))} \lim_r \int_{\partial^* \Omega/r} \psi d\mathcal{H}^{n-1},$$

and as $\Omega/r \rightarrow \{y \cdot \nu > 0\}$ locally as sets of finite perimeter (possibly taking a further subsequence), this limit is

$$\sqrt{A_0/(1 + \rho(0))} \int_{y \cdot \nu = 0} \psi d\mathcal{H}^{n-1}.$$

Therefore $\alpha = A_0/(1 + \rho(0))$, and the conclusion follows from $\alpha \geq (\leq) |\nabla \phi(0)|$. \square

The following is an application of a regularity theorem, first shown in [5] in simpler settings and generalized by De Silva [20] to equations and free boundary conditions which include the one satisfied by u_Ω here. It guarantees that at sufficiently flat points of the boundary $\partial\Omega$, it may be represented as a $C^{1,\alpha}$ graph over a tangent plane.

Lemma 9.2. *There is a $\delta = \delta(v, v_{max}, \eta, R)$ such that the following holds. Fix any $x \in \partial\Omega \cap Q_R$ and $r < \delta$ with $B_r(x) \subseteq Q_R$. If $|\Omega \Delta B_{r,\nu}(x)| < \delta |B_r(x)|$ then we may parametrize the entirety of $\partial\Omega \cap B_{r/2}(x)$ as a $C^{1,\alpha}$ normal graph over the lateral boundary $\partial B_{r,\nu}(x) \cap B_r(x)$, with $C^{1,\alpha}$ norm bounded in terms of only v, v_{max}, η, R .*

Proof. Choosing δ small, Lemma 7.10 implies that

$$B_{cr}(x) \cap \partial\Omega \subseteq \exp_x \{ \nu \in T_x M : |\nu| < r, |g(\nu, \nu)| < \delta' r \}$$

and

$$\sup_{B_{cr}(x)} |u_\Omega - l_{x,\nu}| \leq \delta' cr$$

for a small δ' . Applying Lemma 9.1, we see that Ω and u_Ω is a viscosity solution to the free boundary problem

$$\begin{cases} -\Delta u_\Omega = \lambda_1(\Omega) u_\Omega & \text{on } B_{cr}(x) \cap \Omega \\ u_\Omega > 0 & \text{on } B_{cr}(x) \cap \Omega \\ u_\Omega = 0 & \text{on } B_{cr}(x) \cap \partial\Omega \\ |\nabla u_\Omega| = f(x) & \text{on } B_{cr}(x) \cap \partial\Omega, \end{cases}$$

where $0 < c \leq f \leq C < \infty$ and $\|f\|_{C^{0,\alpha}(\partial\Omega)} \leq C$. Applying the main theorem of [20], it follows that $B_{cr/2}(x) \cap \partial\Omega$ may be parametrized, in normal coordinates, as a $C^{1,\alpha}$ graph over $\{e \cdot \nu = 0\}$, with bounded $C^{1,\alpha'}$ norm for some $\alpha' > 0$. The conclusion now follows from a standard covering argument. \square

Remark 9.3. The previous lemma actually implies that if, working in normal coordinates around x , $\partial\Omega = \{(x', f(x')) : |x'| < \frac{r}{2}\}$ on $B_{r/2}(x)$ with $\|f\|_{C^{1,\alpha}} \leq C$, then f is small. In particular, Lemma 7.10 gives that $|f| = C\delta^{1/n}$, while $[\nabla f]_{C^{0,\alpha'}} \leq Cr^{\alpha-\alpha'} \leq C\delta^{\alpha-\alpha'}$ for $\alpha' < \alpha$. This means that $\|\nabla f\|_{C^{0,\alpha'}} \leq C\delta^\zeta$ for some $\zeta > 0$.

In the next theorem, we say that a collection of open sets, in this case \mathcal{M} , is *uniformly* $C^{k,\alpha}$ if there is a constant C such that for every $x \in \partial U$ for $U \in \mathcal{M}$, $B_{1/C}(x)$ admits normal coordinates, and on this ball ∂U may be expressed as a graph with $C^{k,\alpha}$ norm at most C over a hyperplane in normal coordinates.

Theorem 9.4. *Assume that \mathcal{M} is uniformly $C^{2,\alpha}$. Then for every $r_1 > 0$ there is an $\tau_1 = \tau_1(v, v_{max}, \eta, R, r_1) > 0$ such that if Ω minimizes \mathcal{F}_τ and $\tau < \tau_1$, then $\partial\Omega \cap \{x \in Q_R : d(x, \partial Q_R) > r_1\}$ may be parametrized as a $C^{1,\alpha}$ normal graph (with $C^{1,\alpha}$ norm bounded by r_1) over ∂U for some $U \in \mathcal{M}$.*

If it is known that $\Omega \subset\subset Q_R$ (uniformly), then one may take a fixed $r_1 < d(\Omega, M \setminus Q_R)$ and obtain that the entire boundary $\partial\Omega$ may be parametrized in this manner. In particular, this always applies to the case of M/G_0 compact, using Theorem 6.3. The assumption that $\partial U \in C^{2,\alpha}$ is not needed to guarantee that $\partial\Omega \in C^{1,\alpha}$ (even a flatness condition at every point would suffice), but it is necessary to express $\partial\Omega$ as a graph over ∂U .

Proof. Let $S = \{x \in Q_R : d(x, \partial Q_R) > r_1\}$, and let δ be the parameter from Lemma 9.2. First, observe that under the assumptions on \mathcal{M} , there is an $0 < r' < \delta$ such that for every $x \in \partial U \cap S$ for some $U \in \mathcal{M}$, $|B_{r',\nu_x}(x) \Delta U| < |B_{r'}(x)| \frac{\delta}{4}$ and $B_{r',\nu_x}(x) \setminus B_{r'}(x)$ may be represented as a $C^{1,\alpha}$ normal graph over

∂U with $C^{1,\alpha}$ norm bounded by $1/r'$. Indeed, this follows from the regularity assumption and the tubular neighborhood theorem.

Choose τ_1 small enough that $\mathcal{E}(\Omega) \leq \mathcal{E}_{min} + \tau \leq \mathcal{E}_{min} + \tau_1$ and Lemma 3.14 imply that $|\Omega \Delta U| \leq \frac{\delta}{4} \inf_{x \in Q_R} |B_{r'}(x)|$ for some $U \in \mathcal{M}$. Up to further decreasing τ_1 , we may apply Lemma 7.1 to take Hausdorff distance between ∂U and $\partial \Omega$ smaller than $\kappa r_1/2$, where κ is a fixed constant to be specified below. Then for any $x \in \partial \Omega \cap S$, let $y \in \partial U \cap S$ with $d(x, y) < \kappa r'$, and compute in normal coordinates around y :

$$\begin{aligned} |\Omega \Delta B_{r', \nu_y}(x)| &\leq |U \Delta \Omega| + |B_{r', \nu_y}(y) \Delta U| + |B_{r', \nu_y}(x) \Delta B_{r', \nu_y}(y)| \\ &\leq \frac{\delta}{4} |B_{r'}(x)| + \frac{\delta}{4} |B_{r'}(x)| + |B_{r', \nu_y}(x) \Delta B_{r', \nu_y}(y)| \leq \left[\frac{\delta}{2} + C\kappa \right] |B_{r'}(x)| \end{aligned}$$

where C is a constant depending only on the metric. By choosing κ small enough depending on C and δ , the right-hand side is bounded above by $\delta |B_{r'}(x)|$, so the hypotheses of Lemma 9.2 are satisfied at x at scale r' in direction ν_y .

Apply Lemma 9.2 (and Remark 9.3) on $B_r(x)$ to obtain that $\partial \Omega \cap B_{r'/2}$ may be expressed as a $C^{1,\alpha}$ graph over $B_{r', \nu_x}(x) \setminus B_{r'}(x)$ with $C^{1,\alpha}$ norm bounded by cr_1 for a small c . We may then parametrize $\partial \Omega$ over ∂U instead by composing with the parametrization of $B_{r', \nu_x}(x) \setminus B_{r'}(x)$ over ∂U to obtain the conclusion. \square

Another application of our Euler-Lagrange equations is the following fact about the volume $|\Omega|$:

Proposition 9.5. *There is a $\eta_*(v, v_{max}, R) > 0$, $\mathfrak{T}_*(v, v_{max}, R, \eta) > 0$, and $\tau_*(v, v_{max}, R, \eta, \mathfrak{T}) > 0$ such that for any $\eta < \eta_*$, $\mathfrak{T} < \mathfrak{T}_*$, and $\tau < \tau_*$, every minimizer Ω of \mathcal{F}_τ has $|\Omega| = v$. If M/G_0 is compact, all constants may be taken independent of R .*

Proof. If M/G is compact, we fix η small and select $R = R(\eta)$ large enough that $\Omega \subset \subset Q_R$; this is always possible from Theorem 6.3. All constants will then be independent of this $R(\eta)$, as we will verify below.

If $|\Omega| \neq v$, we apply Corollary 8.10 to give that $A_0 = f'_{v,\eta}(|\Omega|)$ and

$$|\nabla u_\Omega(x)|^2 + \frac{\mathfrak{T}}{2} |\nabla w_\Omega(x)|^2 - \tau \left[b_\Omega(x) + \int_\Omega a_\Omega v_\Omega^x \right] = f'_{v,\eta}(|\Omega|).$$

From Theorem 8.1, we have that $|\nabla w_\Omega(x)|^2 \leq C_0 |\nabla u_\Omega(x)|^2$ and $|\int_\Omega a_\Omega v_\Omega^x| \leq C_0 |\nabla u_\Omega(x)|^2$, where this constant C_0 depends on η . Now select \mathfrak{T}_* small enough (in terms of η) so that $\frac{\mathfrak{T}}{2} |\nabla w_\Omega(x)|^2 \leq \frac{1}{10} |\nabla u_\Omega(x)|^2$. Then select τ_* small enough in terms of η so that $\tau |\int_\Omega a_\Omega v_\Omega^x| \leq \frac{1}{10} |\nabla u_\Omega(x)|^2$, and also so that $\tau |b_\Omega| \leq \eta/2$. This gives

$$\frac{1}{3} f'_{v,\eta}(|\Omega|) \leq |\nabla u_\Omega|^2 \leq 3 f'_{v,\eta}(|\Omega|)$$

along $\partial^* \Omega \cap Q_R$.

Consider the integration by parts formula

$$\frac{1}{3} \sqrt{f'_{v,\eta}(|\Omega|)} \mathcal{H}^{n-1}(\partial \Omega \cap Q_R) \leq - \int_{\partial^* \Omega} g(\nabla u_\Omega(x), \nu_x) d\mathcal{H}^{n-1} = \int_\Omega -\Delta u_\Omega \leq C(R, M) |\Omega|.$$

We have that

$$(9.1) \quad \inf \left\{ \frac{\mathcal{H}^{n-1}(\partial \Omega \cap Q_R)}{|\Omega|} : \Omega \in \mathcal{H} \right\} > c = c(R, v_{max}) > 0$$

from the relative isoperimetric inequality on Q_R , [23, 4.5.2(2)], so $f'_{v,\eta}(|\Omega|) \leq C(R)$. Choosing η_* so that $\eta_* < \frac{1}{C(R)}$, we see that $f_{v,\eta}(|\Omega|) < \frac{1}{\eta_*}$ and $|\Omega| \leq v$. In the case of M/G_0 compact, $\partial \Omega \cap Q_R = \partial \Omega$, we instead use the isoperimetric inequality: for example, from [25, Theorem 3.2] we have $|\Omega|^{\frac{n-1}{n}} \leq C(\mathcal{H}^{n-1}(\partial \Omega) + |\Omega|)$ with a constant depending only on M , and this bounds (9.1) from below by $c(M)v_{max}^{-1/n}$.

It follows that if $|\Omega| \neq v$, then along $\partial^* \Omega \cap Q_R$ the derivative $|\nabla u_\Omega| \leq 9\eta_*$ is small. On $\partial \Omega \cap \partial Q_R$, an elementary comparison argument with w_{Q_R} shows that $|\nabla u_\Omega| \leq C(R)$. From Bochner's identity we have that $\Delta(|\nabla u_\Omega|^2 + Su_\Omega^2) \geq -2\lambda_1(\Omega) Su_\Omega^2$ for S taken large enough in terms of the Ricci curvature of (M, g) . Using Lemma 3.2 and the maximum principle, we see that $|\nabla u_\Omega| \leq C(R)$ on Ω (the point being that this bound is independent of η); the constant may be taken independent of R if M/G_0 is compact.

Now, $|\Omega| < v$ and $\int u_\Omega^2 = 1$, so there must be a point $x \in \Omega$ with $u_\Omega(x) \geq \frac{1}{2}\sqrt{v}$. Using the gradient estimate above, there is a ball $B_r(x) \subseteq \Omega$, with r depending only on R and v . We construct a sequence of balls $B_r(x_k) \subseteq \Omega$ as follows: if there is an $x \in B_{r/2}(x_j)$ for some $j < k$ with $B_r(x) \subseteq \Omega$ and $z \in \partial B_r(x) \cap \partial\Omega \cap Q_R$, set $x_k = x$ and stop. If not, choose any ball $B_r(x)$ with $d(x, x_j) \geq r/2$ and equality for one j . If this is impossible, then $\partial\Omega = \partial Q_R$, which contradicts our standing assumption that $v_{\max} < |Q_R|$. This process must terminate after $C(r)v_{\max}$ steps, and so we have a chain of balls culminating in $B_r(x_j)$ which has $z \in \partial B_r(x_j) \cap \partial\Omega \cap Q_R$. Applying the Harnack inequality along this chain, we have $u_\Omega \geq c(r, J)$ on $B_{r/2}(x_j)$. Letting G' be the Green's function for $B_r(x_j)$, we have that $C(r, J)u_\Omega(y) \geq G'(x_j, y)$ on $B_r(x_j) \setminus B_{r/2}(x_j)$ from the comparison principle. From standard estimates on the Green's function (see [26, Theorem 1.2.8]),

$$C(r, J)u_\Omega(y) \geq G'(x_j, y) \geq c(r, R)d(z, y).$$

From the viscosity form of the Euler-Lagrange equation, Lemma 9.1, we then must have that

$$c(r, R) \leq |\nabla G'(x_j, z)| \leq \sqrt{3f'_{v,\eta}(|\Omega|)} \leq 3\eta_*.$$

If η_* is chosen small in terms of r, R , this gives a contradiction. If M/G_0 is compact, it is easy to see that all constants here may be taken uniformly in R . \square

10. HIGHER REGULARITY

For the purpose of deriving stability estimates, it will be helpful to have $C^{2,\alpha}$ regularity for $\partial\Omega$. Higher regularity depends on the regularity of the functions a_Ω, b_Ω which appear in property (N4) of \mathfrak{h} . In particular, in the case of a_Ω being a multiple of u_U for a smooth domain U , a_Ω is at best Lipschitz continuous on Ω (unless $\Omega = U$). Under this assumption we will have $\partial\Omega$ locally $C^{2,\alpha}$ at flat points for any $\alpha < 1$ and not better. Our approach is based on a boundary Harnack estimate of De Silva and Savin.

Proposition 10.1 ([22], Theorem 2.4). *Let $f \in C(\Omega)$ satisfy $\|-\Delta f\|_{C^{0,\alpha}(\Omega)} \leq 1$, $f = 0$ on $\partial\Omega$, and Ω be a minimizer of \mathcal{F}_τ . Assume that for some $x_0 \in \partial\Omega$, $r < r_0$, and $\delta > 0$ small enough we have $B_r(x_0) \subseteq Q_R$ and $\partial\Omega \cap B_{r/2}(x_0)$ a $C^{1,\alpha}$ graph over $\{e : \nu \cdot e = 0\}$ (in normal coordinates around x_0) with $C^{1,\alpha}$ norm bounded by C_α . Then*

$$\left\| \frac{f}{u_\Omega} \right\|_{C^{1,\alpha}(\bar{\Omega} \cap B_{r/4}(x_0))} \leq C$$

and

$$\left\| \frac{g(\nabla f, \nu_x)}{|\nabla u_\Omega|} \right\|_{C^{1,\alpha}(\partial\Omega \cap B_{r/4}(x_0))} \leq C.$$

Here $\alpha \in (0, 1)$ and the constants depend only on $R, v, v_{\max}, \eta, \alpha, C_\alpha$.

Proof. From Theorem 2.4 in [22],

$$\left\| \frac{f}{u_\Omega} \right\|_{C^{1,\alpha}(\bar{\Omega} \cap B_{r/4}(x))} \leq C \left[\|f\|_{L^\infty(\bar{\Omega} \cap B_{r/4}(x))} + \|\Delta f\|_{C^{0,\alpha}} \right] \leq C,$$

using also the lower bound on u_Ω . The second conclusion follows as $\frac{f}{u_\Omega} \rightarrow \frac{g(\nabla f, \nu_x)}{g(\nabla u_\Omega, \nu_x)}$ nontangentially at every point of $\partial\Omega \cap B_{r/2}(x_0)$, and then passing to the limit. \square

Lemma 10.2. *Let Ω and x_0 be as in Proposition 10.1. Then*

$$\left\| \frac{1}{|\nabla u_\Omega|^2} \int_\Omega h^x f \right\|_{C^{1,\alpha}(\partial\Omega \cap B_{r/4}(x_0))} \leq C \|f\|_{C^{0,\alpha}(\Omega)}.$$

Proof. As in the proof of Lemma 8.13, we have that

$$\int_\Omega h^x f = |\nabla u_\Omega(x)| \int_\Omega K(x, y) f(y) dy = -|\nabla u_\Omega(x)| g\left(\nabla_x \int_\Omega G_\Omega(x, y) f(y) dy, \nu_x\right).$$

If we set $p(x) = \int_{\Omega} G_{\Omega}(x, y) f(y) dy$, this satisfies $-\Delta p = f$ on Ω and $p = 0$ on $\partial\Omega$; applying Proposition 10.1 gives

$$\left\| \frac{g(\nabla p, \nu)}{|\nabla u_{\Omega}|} \right\|_{C^{1,\alpha}(\partial\Omega \cap B_{r/4}(x_0))} \leq C \|f\|_{C^{0,\alpha}}.$$

The conclusion follows from

$$\frac{1}{|\nabla u_{\Omega}(x)|^2} \int_{\Omega} h^x f = \frac{-g(\nabla p(x) \cdot \nu)}{|\nabla u_{\Omega}(x)|}.$$

□

Lemma 10.3. *Let Ω and x_0 be as in Proposition 10.1. Then*

$$\left\| \frac{1}{|\nabla u_{\Omega}|^2} \int_{\Omega} q^x f \right\|_{C^{1,\alpha}(\partial\Omega \cap B_{r/4}(x_0))} \leq C \|f\|_{C^{0,\alpha}}.$$

The function q^x is actually somewhat better behaved than this, but we will not require an optimal estimate below.

Proof. Set $\bar{h}^x = \frac{h^x}{|\nabla u(x)|^2}$ and $\bar{q}^x = \frac{q^x}{|\nabla u(x)|^2}$. Rewriting the equation for q in terms of these, we have that

$$\begin{cases} -\Delta \bar{q}^x = \lambda_1(\Omega)[\bar{q}^x + \bar{h}^x] - u_{\Omega} & \text{on } \Omega \\ \int \bar{q}^x u_x = -\int \bar{h}^x u_x = -\frac{1}{\lambda_1(\Omega)} \\ \bar{q}^x = 0 & \text{on } \partial\Omega. \end{cases}$$

Let us decompose $f = f_1 + f_2$, where $f_2 = u_{\Omega} \int u_{\Omega} f$, and then solve for the potential function

$$\begin{cases} -\Delta p = \lambda_1(\Omega)p + f_1 & \text{on } \Omega \\ \int p u_{\Omega} = 0 \\ p = 0 & \text{on } \partial\Omega. \end{cases}$$

This admits a unique solution via the Fredholm alternative (using $\int f_1 u_{\Omega} = 0$) which satisfies $\|p\|_{L^2} \leq C \|f_1\|_{L^2} \leq C \|f\|_{C^{0,\alpha}}$. From elliptic regularity, we have that

$$\|p\|_{C^{0,\alpha}(\Omega)} \leq C[\|f_1\|_{L^{\infty}} + \|p\|_{L^2}] \leq C \|f\|_{C^{0,\alpha}}.$$

Now, multiply the equation for \bar{q}^x by p and integrate:

$$\int \lambda_1(\Omega) \bar{h}^x p = \int p(-\Delta - \lambda_1(\Omega)) \bar{q}^x = \int f_1 \bar{q}^x.$$

The other term on the left vanishes as $\int p u_{\Omega} = 0$, while the simplification on the right is from Green's identity and the equation for p .

On the other hand,

$$\int f_2 \bar{q}^x = \left(\int u_{\Omega} f \right) \left(\int u_{\Omega} \bar{q}^x \right) = -\frac{1}{\lambda_1(\Omega)} \int u_{\Omega} f.$$

Combining, we have that

$$\int f \bar{q}^x = \lambda_1(\Omega) \int \bar{h}^x p - \frac{1}{\lambda_1(\Omega)} \int u_{\Omega} f,$$

which is bounded in $C^{1,\alpha}$ by Lemma 10.2 applied to p . □

Corollary 10.4. *Let Ω be a minimizer of \mathcal{F}_{τ} . Assume that for some $x_0 \in \partial\Omega$, $r < r_0$, and $\delta > 0$ small enough we have $B_r(x_0) \subseteq Q_R$ and $|\Omega \Delta B_{r,\nu}(x_0)| < \delta |B_r(x)|$. Then for each $\alpha \in (0, 1)$, $\partial\Omega \cap B_{r/64}(x_0)$ is given by a $C^{2,\alpha}$ normal graph over $\partial B_{r,\nu}(x_0) \cap B_r(x_0)$, with the $C^{2,\alpha}$ norm bounded by a constant depending only on R, v, v_{max}, η , and α .*

Proof. From Lemma 9.2, $\partial\Omega \cap B_{r/2}(x_0)$ is a C^{1,α_0} graph over $\{e : e \cdot \nu = 0\}$ in normal coordinates, where $\alpha_0 > 0$ is the fixed exponent from that lemma. Applying Lemmas 10.2 and 10.3 with $f = a_\Omega \in C^{0,1}$, we have that

$$\left\| \frac{v_\Omega^x}{|\nabla u(y)|^2} \right\|_{C^{1,\alpha_0}(\partial\Omega \cap B_{r/4}(x_0))} \leq C.$$

Applying Proposition 10.1 to w_Ω gives that

$$\left\| \frac{|\nabla w_\Omega|^2}{|\nabla u_\Omega|^2} \right\|_{C^{1,\alpha_0}(\partial\Omega \cap B_{r/4}(x_0))} \leq C.$$

We also have $\|b_\Omega\|_{C^{1,1}} \leq 1$ by assumption (N4). Proceeding as in Theorem 8.1, we may write the Euler-Lagrange equation for u as

$$|\nabla u(x)|^2 = \rho(x),$$

where $c \leq \rho \leq C$ and $\rho \in C^{1,\alpha_0}(\partial\Omega \cap B_{r/4}(x_0))$. Now apply [30, Section 9] or [21] to obtain that $\partial\Omega$ may be represented as a C^{2,α_0} graph on $\partial\Omega \cap B_{r/8}(x_0)$.

Now repeat this argument on $B_{r/8}(x)$, except using that $\partial\Omega$ is $C^{2,\alpha_0} \subseteq C^{1,\alpha}$ rather than C^{1,α_0} to recover the conclusion as stated. \square

Combining this with Theorem 9.4 gives that all of $\partial\Omega \subset\subset Q_R$ is $C^{2,\alpha}$, so long as all sets in \mathcal{M} are smooth:

Corollary 10.5. *Let Ω be as in Theorem 9.4 and assume that \mathcal{M} is uniformly C^4 . Then $\partial\Omega \cap \{x \in Q_R : d(x, \partial Q_R) > r_1\}$ may be parametrized as a $C^{2,\alpha}$ normal graph over ∂U for any $\alpha < 1$, with $C^{2,\alpha}$ norm bounded by r_1 .*

APPENDIX A. ADMISSIBLE NONLINEARITIES

In Section 2.2, we gave several examples of nonlinearities \mathfrak{h} used in the definition of the main functional \mathcal{F}_τ . In this appendix, we verify that these examples are admissible nonlinearities in the sense of Definition 2.11. We also define the notion of a set center, which generalizes the notion of the barycenter of a set in Euclidean space.

To begin, it will be useful to show that a constant multiple of $d_*(\Omega, U)^2$ defined in (2.10) is itself an admissible nonlinearity. Given any bounded open set U with C^2 boundary, define the functional

$$(A.1) \quad \mathfrak{h}_U(\Omega) = \frac{d_*(\Omega, U)^2}{C_1} = \frac{1}{C_1} \left[\int_\Omega \psi_U - \int_U \psi_U + \int |u_\Omega - u_U|^2 \right]$$

on bounded open sets Ω with $|\Omega| \leq v_{\max}$ and with $\mathcal{E}(\Omega) \leq \mathcal{E}_{\min} + \tau_0$; recalling Remark 2.5, \mathfrak{h}_U is well-defined on this class of sets.

Lemma A.1. *For C_1 chosen sufficiently large depending on U and v_{\max} , the functional $\mathfrak{h}_U(\Omega)$ is an admissible nonlinearity with respect to the trivial subgroup of the isometry group.*

Proof. Choose $C_1 \geq v_{\max} |\max f| + 4$ to guarantee that $\mathfrak{h}_U(\Omega) \in [0, 1]$ for any Ω so (N1) is satisfied. To see the upper bound, recall that $\int u_\Omega^2 = \int u_U^2 = 1$; the lower bound follows because $d_*(U, \Omega) \geq 0$. Since G_0 is trivial, (N2) holds automatically. Checking (N3) is also straightforward:

$$(A.2) \quad |\mathfrak{h}_U(\Omega) - \mathfrak{h}_U(\Omega')| \leq \frac{1}{C_1} \left[|\Omega \Delta \Omega'| \max |f| + \int 2|u_{\Omega'} - 2u_\Omega| |u_U| \right],$$

so as long as $C_1 \geq 2 \max u_U$, which is bounded by Lemma 3.2, and so (N3) follows. For (N4), we set $\Omega_t = \phi_t(\Omega)$ and perform the same computation more carefully to give

$$(A.3) \quad \mathfrak{h}_U(\Omega_t) - \mathfrak{h}_U(\Omega) = \frac{1}{C_1} \left[\int_{\Omega_t} \psi_U - \int_\Omega \psi_U + 2 \int (u_\Omega - u_{\Omega_t}) u_U \right].$$

Set $a_\Omega = -\frac{2}{C_1} u_U$, $b_\Omega = \frac{1}{C_1} \psi_U$. After possibly choosing C_1 larger, we see that (N4) is satisfied, using elliptic regularity and the smoothness of U to bound ∇u_U . \square

Notice that in this example, while b_Ω here is smooth, a_Ω is not better than Lipschitz; is the limiting factor for the higher regularity discussed in Section 10.

Example 2.6 considered a nonlinearity \mathfrak{h} in the case when there is a unique (regular) minimizer U of the base energy. Up to multiplication by a constant $C = C(U)$, it follows that the nonlinearity in Example 2.6 is admissible with respect to the trivial subgroup of the isometry group for any choice of $c \in (0, 1]$ in (2.9). Indeed, the functional $\mathfrak{h} = \mathfrak{h}_c$ of Example 2.6 (normalized by a constant $C = C(U)$ independent of c) is given by $\phi_c \circ \mathfrak{h}_U$ where $\phi_c : [0, 1] \rightarrow [0, 1]$ is a smooth function with $|\phi'_c| \leq 1$ for all $c \in (0, 1]$, and is thus admissible by Remark 2.12.

We now move toward verifying that the nonlinearity of Example 2.10 is admissible. To this end, let us introduce the notion of a set center.

Definition A.2 (Set centers). *Let $G_0 \leq G$ be a closed subgroup of isometries of M , $\varepsilon > 0$, and U a fixed bounded open set. We say that a mapping from the class of bounded, open, nonempty subsets E of M with*

$$(A.4) \quad \inf_{e \in G_0} |E \Delta e(U)| \leq \varepsilon$$

to points x_E in a complete Riemannian manifold (N, g_N) is a set center adapted to U if it satisfies the following properties:

(C1) *If $E, E' \subseteq Q_R$, then $d_N(x_E, x_{E'}) \leq C(R)|E \Delta E'|$.*

(C2) *For every E , there is an $e \in G_0$ such that $x_E = x_{e(U)}$.*

(C3) *For any $e, e' \in G_0$, $|e(U) \Delta e'(U)| \leq C(E)d_N(x_{e(E)}, x_{e'(E)})$.*

(C4) *For $x \in M$ and $r \leq r_0$, let ϕ_t be a one-parameter family of diffeomorphisms with $\phi_0(x) = x$ and $|\partial_t \phi_t| \leq 1$ such that $\{\phi_t(x) \neq x\} \subseteq B_r(x_0)$. If the sets $\phi_t(\Omega)$ satisfy (A.4) and $\mathcal{H}^{n-1}(\partial\Omega) < \infty$, then $t \mapsto x_{\phi_t(\Omega)}$ defines a C^1 curve in N and for any tangent vector $v \in T_{x_\Omega}N$,*

$$(A.5) \quad \limsup \frac{1}{|t|} \left| t g_N \left(x'_{\phi_t(\Omega)}|_{t=0}, v \right) - \int_{\phi_t(\Omega)} a_\Omega^v + \int_\Omega a_\Omega^v \right| \leq Cr^n$$

for a function a_Ω^v independent of ϕ_t and depending linearly on v with $\|a_\Omega^v\|_{C^2(Q_R)} \leq C(R)|v|$.

Given a bounded open set E we let $U_E = e(U)$ denote image under the unique isometry $e \in G_0$ of U such that E and U_E have the same set center.

Property (C1) implies x_E is a kind of Lipschitz mapping from sets to the space of centers, while (C3) suggests it acts like a projection onto the images of U under G_0 action. Property (C4) implies that it is a C^1 -regular projection in an appropriate sense: to understand (A.5), note that

$$\int_{\phi_t(\Omega)} a_\Omega^v - \int_\Omega a_\Omega^v = t \int_{\partial^* \Omega} a_\Omega^v g(\phi'_0, \nu_x) d\mathcal{H}^{n-1} - t \int_\Omega g(\phi'_0, \nabla a_\Omega^v) + o(t).$$

The second term is bounded by Cr^nt , while the first captures the ‘‘leading-order’’ change in Ω under $t \mapsto \phi_t(\Omega)$. Thus (A.5) ensures that the v component of the derivative of $x_{\phi_t(\Omega)}$ exists and is proportional to $\int_{\partial^* \Omega} a_\Omega^v g(\phi'_0, \nu_x) d\mathcal{H}^{n-1}$ for some sufficiently smooth function a_Ω^v , up to error of size r^n . This type of error is lower-order when considering localized deformations near a point on $\partial\Omega$.

Note that a set center may map to points in a Riemannian manifold (N, g_N) that is not (M, g) . An example to keep in mind here is if (M, g) is a cylinder $S^{n-1} \times \mathbb{R}$ equipped with the standard product metric, a set center can be defined by choosing the Euclidean barycenter of a set when (M, g) is embedded into \mathbb{R}^{n+1} in the standard way.

Let us give some examples of set centers.

Example A.3. If G_0 is trivial, the constant mapping $x_E = x_0$ is a set center adapted to any U .

Example A.4. Suppose (M, g) is simply connected and has nonpositive sectional curvature. The barycenter $x \mapsto \operatorname{argmin}_x \int_E d^2(x, y) dm(y)$ exists and is unique. This is a set center adapted to any U so long as (1) G_0 acts transitively on M and (2) if $e \in G_0$ fixes x_U , it fixes U . In particular, this applies to \mathbb{R}^n or hyperbolic space, $G_0 = G$, and $U = B_r(x)$ a ball. It also applies to \mathbb{R}^n with U an arbitrary (nice) open set and G_0 the group of translations.

Example A.5. Let (M, g) be the round sphere, embedded in the standard way in \mathbb{R}^{n+1} . Let $G_0 = G$ be the full isometry group and $U = B_r(x)$ be a geodesic ball with $r < \pi$. For any open $E \subseteq S^n$, let $y_E = \int_E y d\mathcal{H}^n(y)$, where $y, y_E \in \mathbb{R}^{n+1}$ and this is a (vector-valued) surface integral. Note that $y_U \neq 0$ and that for any E, E' , we have

$$|y_E - y_{E'}| \leq \frac{1}{|E|} \int_{E \Delta E'} |y| + |S^n| \left| \frac{1}{|E|} - \frac{1}{|E'|} \right| \leq C|E \Delta E'|.$$

In particular, if $|E \Delta U| \leq \varepsilon$, we may ensure that $|y_E| \geq r > 0$. For any such E , the mapping $E \mapsto x_E = \frac{y_E}{|y_E|}$ is a set center. Indeed, then $|x_E - x_{E'}| \leq \frac{1}{r}|y_E - y_{E'}|$, so this satisfies (C1). It also satisfies (C2), as G acts transitively on S^n . Property (C3) can be checked using that $U = B_r(x)$ is invariant under rotation about its set center, and (C4) may be verified similarly to the proof of Proposition A.6 below.

Let us give a proof that the Euclidean barycenter is a set center in the sense of Definition A.2; the other examples described above can be checked similarly. First of all, the barycenter on Euclidean space is an example of a set center.

Proposition A.6. *Let (M, g) be Euclidean space, and let $G_0 = G$ and $U = B_1$. Then the barycenter $x_E = \int_E x \in \mathbb{R}^n$ is a set center satisfying (C1)-(C4). Here ε in (A.4) may be taken arbitrary.*

Proof. For (C1),

$$|x_E - x_{E'}| = \left| \frac{1}{|E|} \int_E x - \frac{1}{|E'|} \int_{E'} x \right| \leq \frac{1}{|E|} \int_{E \Delta E'} |x| + \left| \frac{|E'|}{|E|} - 1 \right| R \leq C|E' \Delta E|.$$

Then (C2), (C3) are immediate, and for (C4) we may compute $|\Omega_t| - |\Omega| = \mathcal{H}^{n-1}(\partial^* \Omega)t + o(t) = O(t)$ and

$$\int_{\phi_t(\Omega)} x - \int_{\Omega} x = t \int_{\partial^* \Omega} x \cdot \nu_x - t \int \phi'_t(0) + o(t) = O(t),$$

from which we see that

$$x_{\Omega_t} = \frac{1}{|\phi_t(\Omega)|} \int_{\phi_t(\Omega)} x = x_{\Omega} + \int_{\phi_t(\Omega)} \left[\frac{x}{|\Omega|} - \int_{\Omega} x \right] - \int_{\Omega} \left[\frac{x}{|\Omega|} - \int_{\Omega} x \right] + O(t^2).$$

This shows x_{Ω_t} is a C^1 curve, and $x'_{\Omega_t}|_{t=0}$ satisfies (A.5) with $\alpha_{\Omega}^v = \left[\frac{x}{|\Omega|} - \int_{\Omega} x \right] \cdot v$. \square

More general constructions can be carried out here, for example considering set projections onto \mathcal{M} in the case when \mathcal{M} admits “smooth” parameterizations by finite-dimensional manifolds, in which case the key point will always be verifying properties (N3) and (N4) using analogues of (C1) and (C4).

Now, given a bounded open set U with C^2 boundary, a subgroup of isometries $G_0 \leq G$, and a set center $\Omega \mapsto x_{\Omega}$, define the functional

$$(A.6) \quad \mathfrak{h}(\Omega) := \mathfrak{h}_{U_{\Omega}}(\Omega) \quad x_{U_{\Omega}} = x_{\Omega}$$

where \mathfrak{h}_U was defined in (A.1). By properties (C2) and (C3) of set centers, \mathfrak{h} is well-defined for all Ω with $|\Omega \Delta e(U)| \leq \varepsilon$ for some $e \in G_0$. From Lemma 3.10, this holds for all Ω with $\mathcal{E}(\Omega) \leq \mathcal{E}_{min} + \tau_0$.

Proposition A.7. *The functional $\mathfrak{h}(\Omega)$ in (A.6) is an admissible nonlinearity with respect to G_0 in the sense of Definition 2.11 for C_1 sufficiently large.*

Proof of Proposition A.7. Properties (N1) and (N2) are immediate from the construction. For (N3), take Ω, Ω' and apply property (C1) of set centers:

$$d_N(x_{\Omega}, x_{\Omega'}) \leq C(R)|\Omega \Delta \Omega'|.$$

This means there exists $e \in G_0$ with $e(U_{\Omega}) = U_{\Omega'}$ and $|U_{\Omega'} \Delta U_{\Omega}| \leq C|\Omega \Delta \Omega'|$. To simplify notation, we set $U = U_{\Omega}$ and $U' = U_{\Omega'}$.

First of all, after possibly composing e with an isometry which fixes U , we may assume $d(e, \text{id}) \leq \varepsilon_1$ small. Indeed, if this is false then there is a sequence $e_k \in G_0$ with $|U \Delta e_k(U)| \rightarrow 0$ but $d(e_k, e_0) > \varepsilon_1$ for

any e_0 fixing U . It is then straightforward to show (see [33], Theorem 3 and subsequent remarks) that the e_k have a subsequence converging to some e with $|U\Delta e(U)| = 0$ and therefore fixing U . This is a contradiction.

Let $S_U \subseteq \mathfrak{G} = T_{\text{id}}G_0$ be the tangent space to the subgroup of isometries which fix U , and H_U a subspace of \mathfrak{G} such that $H_U \oplus S_U = \mathfrak{G}$; set $V = \{v \in H_U : |v| = 1\}$. One may verify that that for any $U' = \exp_{\text{id}} w(U)$ for w with $|w| \leq \varepsilon_1$, U' may also be represented as $\exp_{\text{id}} tv(U)$ for $v \in V$ and $|t| \leq C\varepsilon_1$ (see e.g. [35, Theorem 3.58]). We may estimate the value of t more precisely:

$$(A.7) \quad |U\Delta U'| \leq C \sup_{x \in U} d(\exp_{\text{id}} tv(x), x) \leq C(U)t.$$

The reverse bound is also valid. Indeed, for any fixed $v \in V$, we have $\int_{\partial U} |g(v(x), \nu_x)| > 0$ where ν_x is the outward unit normal to ∂U : if $|g(v(x), \nu_x)| \equiv 0$, then the isometries generated by v fix U , contradicting that $v \in V \subseteq H_U \setminus \{0\}$. As V is compact, this gives that $\inf_{v \in V} \int_{\partial U} |g(v(x), \nu_x)| \geq c(U) > 0$. Then we may estimate

$$(A.8) \quad |U\Delta U'| = t \int_{\partial U} |g(v(x), \nu_x)| + O(t^2) \geq c(U)t.$$

This value t also controls the distance between U and U' in stronger topologies:

$$\sup_{x \in Q_R} |d(x, \partial U) - d(x, \partial U')| = \sup_{x \in Q_R} |d(x, \partial U) - d(\exp_{\text{id}} -tv(x), \partial U)| \leq Ct$$

and

$$\int |u_U - u_{U'}| \leq \sup_x |u_U(x) - u_U(\exp_{\text{id}} -tv(x))| \leq Ct \sup_x |\nabla u_U| \leq Ct.$$

In the line above, we recall from Lemma 3.14 that U has a *unique* nonnegative, normalized, first eigenfunction u_U , and so $u_{U'} = u_U \circ e^{-1}$. So, we find that

$$|\mathfrak{h}_U(\Omega) - \mathfrak{h}_{U'}(\Omega)| \leq \frac{1}{C_1} \left[\int_{\Omega} |\psi_U - \psi_{U'}| + 2 \int |u_{\Omega}| |u_U - u_{U'}| \right] \leq \frac{1}{C_1} Ct \leq \frac{C}{C_1} |\Omega \Delta \Omega'|.$$

Combining this with (A.2) of Lemma A.1 and choosing C_1 small enough gives (N3).

We now verify (N4). To begin, let $U_t = e_t(U)$ denote $U_{\phi_t(\Omega)}$, the unique isometry of U with the same set center as $\phi_t(\Omega)$. Our first claim is that after possibly modifying the e_t , they may be represented for t small as $e_t = \exp_{\text{id}} \gamma(t)$, where $\gamma(t)$ is a C^1 curve in $H_U \subseteq \mathfrak{G}$ with $\gamma(0) = 0$. To see this, consider the mapping $T : v \mapsto x_{\exp_{\text{id}} v(U)}$ from a small ball $B_{\varepsilon_1}(0) \subseteq \mathfrak{G}$ to N . By property (C4) of set centers, T has continuous directional derivatives on B_{ε_1} , and so it is C^1 . We also have that $d_0 T(S_U) = \{0\}$, since any curve of isometries fixing U has set center constant x_U and by (A.8) and property (C3), this is precisely the kernel of $d_0 T$. So, $T|_{H_U}$ is a C^1 diffeomorphism onto its image from the inverse function theorem.

Applying property (C4), the curve $t \mapsto x_{\phi_t(\Omega)}$ is C^1 , and there is a unique C^1 curve $\gamma : (-c, c) \rightarrow H_U \subseteq \mathfrak{G}$ such that $x_{\phi_t(\Omega)} = \exp_{\text{id}} \gamma(t)(U)$. By definition of U , $\gamma(0) = 0$. Furthermore, we have that

$$\gamma'(0) = d_0 T^{-1}|_{H_U} (x'_{\phi_t(\Omega)}|_{t=0}) = \lim_{t \rightarrow 0} \frac{1}{|t|} \int_{\phi_t(\Omega)} A_{\Omega} - \int_{\Omega} A_{\Omega} + O(r^n)$$

for some $A_{\Omega} : Q_R \rightarrow \mathfrak{G}$ with $\|A_{\Omega}\|_{C^2} \leq C$, i.e. for each x , $A_{\Omega}(x)$ is a vector field generating isometries.

Let us next estimate the ψ_U term:

$$\psi_{U_t}(x) = \psi_U(\exp_{\text{id}} -\gamma(t)(x)) = \psi_U(x) - tg(\nabla \psi_U(x), \gamma'(0)(x)) + o(t)|D^2 \psi_U|.$$

so

$$\int_{\Omega} [\psi_{U_t}(\Omega) - \psi_U(\Omega)] = - \int_{\Omega} tg_x(\nabla \psi_U(x), \gamma'(0)) + o(t) + O(r^n t).$$

In particular,

$$\limsup_{t \rightarrow 0} \frac{1}{|t|} \left| \int_{\Omega} \psi_{U_t}(\Omega) - \psi_U(\Omega) + \int_{\phi_t(\Omega)} \int_{\Omega} g(A_{\Omega}(y), \nabla \psi_U(x)) dx dy - \int_{\Omega} \int_{\Omega} g(A_{\Omega}(y), \nabla \psi_U(x)) dx dy \right| \leq Cr^n.$$

Here and in the remainder of the proof we write dx to denote $dm(x)$ to consolidate notation. The estimate on the eigenfunction term is similar, at least so long as we avoid the set $\{d(x, \partial U) < C|t|\}$ for C large enough

that $\exp_{\text{id}} - \gamma(s)(x) \neq \partial U$ for $|s| \leq |t|$, the issue being that u_U is not smooth near ∂U . Away from this set, we have:

$$u_{U_t}(x) = u_U(x) - tg(\nabla u_U(x), \gamma'(0)) + O(t^2 \sup_U |D^2 u_U|).$$

Note that while u_U is not smooth, $D^2 u_U$ is uniformly bounded on U . Therefore,

$$\begin{aligned} \int |u_{U_t} - u_\Omega|^2 - \int |u_U - u_\Omega|^2 &= 2 \int u_\Omega (u_U - u_{U_t}) \\ &= 2t \int g(\nabla u_U(x), \gamma'(0)) + o(t) + O\left(\int_{\{d(x, \partial U) < C|t|\}} |u_U| + |u_{U_t}| + t|g(\nabla u_U(x), \gamma'(0))|\right). \end{aligned}$$

The rightmost term is actually $o(t)$ as well; $|\{d(x, \partial U) < C|t|\}| \rightarrow 0$, while also $|u_U| + |u_{U_t}| \leq Ct$ and $t|g(\nabla u_U(x), \gamma'(0))| \leq Ct$ over this region. So, rewriting the expansion above, we have

$$\limsup_{t \rightarrow 0} \frac{1}{|t|} \left| \int |u_{U_t} - u_\Omega|^2 - \int |u_U - u_\Omega|^2 - 2t \int g(\nabla u_U(x), \gamma'(0)) \right| = 0.$$

Moreover, the first-order term may be re-expressed in terms of A_Ω as well:

$$\limsup_{t \rightarrow 0} \frac{1}{|t|} \left| t \int g(\nabla u_U(x), \gamma'(0)) - \int_{\phi_t(\Omega)} \int_\Omega g_x(A_\Omega(y), \nabla u_U(x)) dx dy - \int_\Omega \int_\Omega g_x(A_\Omega(y), \nabla u_U(x)) dx dy \right| \leq Cr^n.$$

To summarize, setting

$$a_\Omega^0 = \int_\Omega g_x(A_\Omega(y), 2\nabla u_U(x) + \nabla \psi_U(x)) dx$$

we have that $\|a_\Omega^0\|_{C^2(Q_R)} \leq C(U)$, and

$$\limsup_{t \rightarrow 0} \frac{1}{|t|} \left| \mathfrak{h}_{U_t}(\Omega) - \mathfrak{h}_U(\Omega) + \int_{\phi_t(\Omega)} a_\Omega^0 - \int_\Omega a_\Omega^0 \right| \leq Cr^n.$$

Together with (A.3) this implies property (N4). This completes the proof. \square

As we saw following Lemma A.1 above, it follows immediately from Proposition A.7 and Remark 2.12 that the nonlinearity of Example 2.10 is an admissible nonlinearity with respect to G in the sense of Definition 2.11 for all $c \in (0, 1]$, after possibly normalizing by a constant $C = C(U)$.

APPENDIX B. DOMAINS ADMITTING LINEAR-GROWTH SOLUTIONS TO ELLIPTIC EQUATIONS ARE NTA

The purpose of this appendix is to show that if a domain Ω supports a nonnegative function u vanishing on $\partial\Omega$, growing like $d(x, \partial\Omega)$ from the boundary, and solving an elliptic equation with bounded right-hand side, then Ω is NTA (locally). Our first lemma follows the approach of [2], using a monotonicity formula argument, except that the u is not assumed to be a minimizing solution to a free boundary problem. In fact, it may solve an equation with right-hand side, and not even all the way up to the boundary (so long as we are only looking for Harnack chains between points a distance η from $\partial\Omega$).

Lemma B.1. *Let $\Omega \subseteq \mathbb{R}^n$ be open and u be a continuous function on $\bar{\Omega} \cap B_1$ (which is C^2 on Ω), with $u = 0$ on $\partial\Omega$ and $u > 0$ on Ω . Then for any $K, M, m > 0$, there are $\eta, \delta > 0$ such that if*

- (1) $0 \in \partial\Omega$.
- (2) $|\nabla u| \leq 1$ on $\Omega \cap B_1$.
- (3) For all $x \in \bar{\Omega} \cap B_1$ and $r < 1$, $\sup_{B_r(x)} u \geq mr$.
- (4) $|\Delta u| \leq \delta$ on $\{d(x, \Omega^c) > \delta\} \cap B_1$.
- (5) Ω satisfies the outer clean ball condition with constant K at 0.

then for any points $x_1, x_2 \in B_{M\eta}$ with $d(x_1, \Omega^c), d(x_2, \Omega^c) > \eta$, there exists a curve $\gamma \subseteq B_1 \cap \{d(x, \Omega^c) > c\eta\}$ of length at most C connecting x_1 and x_2 .

Proof. Set $v(x) = u(x) + \frac{\delta}{2n}|x|^2$; then $\Delta v \geq 0$ on $\{d(x, \Omega^c) > \delta\} \cap B_1$, and moreover as long as $\delta < 1$ assumption (2) gives $|\nabla v| \leq 2$. Consider the set $A = \{x \in \Omega : v(x) > c_0\eta\} \cap B_{1/2}$, noting that so long as we choose δ small enough in terms of η, m ,

$$\left\{d(x, \Omega^c) > \frac{1}{2}\eta\right\} \cap B_{1/2} \subseteq A \subseteq \{d(x, \Omega^c) > c_1\eta\},$$

for some c_0, c_1 depending only on n and m . Indeed, if $d(x, \Omega^c) > \frac{1}{2}\eta$, we apply assumption (3) on $B_{\frac{\eta}{8}}(x)$ to learn that $\sup_{B_{\frac{\eta}{8}}(x)} u \geq m\frac{\eta}{8}$. Then so long as $\delta < \eta/4$ (and hence (4) applies on $B_{\eta/4}(x)$) the Harnack inequality gives

$$m\frac{\eta}{8} \leq \sup_{B_{\frac{\eta}{8}}(x)} u \leq Cu(x).$$

Setting $c_0 = \frac{m}{16C}$ ensures that $u(x) \geq 2c_0\eta$, and as long as $\delta < c_0\eta$, this gives $v(x) > u(x) - \delta \geq c_0\eta$ and so $x \in A$. On the other hand, if $x \in A$ and $\delta < \frac{c_0}{2}\eta$, we have $u(x) > \frac{c_0}{2}\eta$. In light of (2), this guarantees that $d(x, \Omega^c) > \frac{c_0}{2}\eta$, and we set $c_1 = \frac{c_0}{2}$.

Let A_1, A_2 be the two connected components of A which contain x_1, x_2 respectively. We will show that $A_1 = A_2$. From a computation like the one just performed, we have that $v(x_i) > 2c_0\eta$. Set $d_i = d(x_i, A^c)$, and let $y_i \in \partial A_i$ such that $|y_i - x_i| = d_i$. From assumption (1) we must have $d_i < M\eta$. Our first task will be to show that $|\nabla v|$ is relatively large on a set near x_i .

Consider a line segment with endpoints x_1 and y_1 . Integrating along it,

$$\begin{aligned} c_0\eta < v(x_1) - v(y_1) &= \int_0^1 \nabla v(tx_1 + (1-t)y_1) \cdot (x_1 - y_1) dt \\ &\leq d_1 \int_0^1 |\nabla v(tx_1 + (1-t)y_1)| dt \leq M\eta \int_0^1 |\nabla v(tx_1 + (1-t)y_1)| dt. \end{aligned}$$

As $|\nabla v| \leq 2$ everywhere, the second-to-last step implies that $d_1 \geq \frac{c_0\eta}{2}$. We can say more: if $|\nabla v| < \frac{c_0}{8M}$ for all $t > \frac{c_0}{8M}$, we could estimate

$$M\eta \int_0^1 |\nabla v(tx_1 + (1-t)y_1)| dt \leq \left(1 - \frac{c_0}{8M}\right) \frac{c_0\eta}{8} + \frac{c_0\eta}{8} 2 < \frac{c_0\eta}{2},$$

and this is a contradiction. Hence there must be a $t > \frac{c_0}{8M}$ with $|\nabla v(tx_1 + (1-t)y_1)| \geq \frac{c_0}{8M}$; set $z_1 = tx_1 + (1-t)y_1$. Letting $c_2 = \frac{c_0}{8M} \cdot \frac{c_0}{8}$, we have that

$$|z_1 - y_1| = t|x_1 - y_1| = td_1 \geq \frac{c_0}{8M} \cdot \frac{c_0\eta}{2} = c_2\eta.$$

In particular $B_{c_2\eta}(z_1) \subseteq A_1 \subseteq \{d(x, \Omega^c) > \delta\}$ and we may apply an elliptic regularity estimate to v on this ball:

$$(c_2\eta)^{1+\alpha} [\nabla v]_{C^{0,\alpha}(B_{\frac{c_2\eta}{2}}(z_1))} \leq C \left[\operatorname{osc}_{B_{c_2\eta}(z_1)} v + \delta(c_2\eta)^2 \right] \leq C\eta,$$

so in particular

$$|\nabla v(x) - \nabla v(z_1)| \leq C \left(\frac{|x - z_1|}{\eta} \right)^\alpha < \frac{c_0}{16M}$$

so long as $|x - z_1| < c_3\eta$. For any such x , we have $|\nabla v(x)| \geq \frac{c_0}{16M}$. Integrating,

$$\frac{1}{(2M\eta)^2} \int_{B_{2M\eta} \cap A_1} |\nabla v|^2 |x|^{2-n} dx \geq \frac{1}{(2M\eta)^n} \int_{B_{c_3\eta}(z_1)} |\nabla v|^2 \geq \left(\frac{c_3}{2M} \right)^n |B_1| \frac{c_0}{16M} \geq c_4.$$

Set

$$J(r) = \frac{1}{r^4} \int_{B_r \cap A_1} |\nabla(v - c_0\eta)_+|^2 |x|^{2-n} dx \int_{B_r \cap A_2} |\nabla(v - c_0\eta)_+|^2 |x|^{2-n} dx.$$

We have just shown that $J(2M\eta) \geq c_4^2$ (as the computations apply equally well to x_2 and A_2). It is straightforward to check that $J(1/2) \leq C$ using just $|\nabla v| \leq 2$. From the enhanced monotonicity property of J [2, Lemma 4.4] and assumption (5), $r \mapsto \frac{J(r)}{r^\beta}$ is an increasing function for some $\beta > 0$ depending only on K . This means that $c_4^2 \leq J(2M\eta) \leq C(4M\eta)^\beta$. This is a contradiction if η is chosen small in terms of M and c_4 .

We have shown that $A_1 = A_2$. Now cover the closure of A_1 by balls of radius $r = \frac{c_1}{4}\eta$; it is always possible to do so by at most $N = C\eta^{-n}$ balls, and then to connect any two points in A_1 by a piecewise linear curve of length $CN\eta$ contained in the union of these balls. This curve has bounded length and stays a distance at least $c_1\eta/4$ away from $\partial\Omega$, concluding the proof. \square

If instead u solves an elliptic equation with Lipschitz coefficients, the previous lemma still applies after a suitable change of variables and scaling argument. Here we carefully exploit the form of assumption (4) above.

Lemma B.2. *Let $\Omega \subseteq \mathbb{R}^n$ be open and u be a continuous function on $\bar{\Omega} \cap B_1$ (which is C^2 on Ω), with $u = 0$ on $\partial\Omega$ and $u > 0$ on Ω . Then for any $\lambda, K, M, m > 0$, there is an $\eta_0 > 0$ such that if*

- (1) $0 \in \partial\Omega$
- (2) $|\nabla u| \leq 1$ on $\Omega \cap B_1$.
- (3) For all $x \in \bar{\Omega} \cap B_1$ and $r < 1$, $\sup_{B_r(x)} u \geq mr$.
- (4) $a_{ij}\partial_i\partial_j u = f$ on $\Omega \cap B_1$, with $\|f\|_{C^{0,1}} \leq 1$, where $\lambda|\xi|^2 \leq a_{ij}(x)\xi_i\xi_j \leq \lambda^{-1}|\xi|^2$ is a matrix-valued function with $[a_{ij}]_{C^{0,1}(B_1)} \leq 1$.
- (5) Ω satisfies the outer clean ball condition with constant K at 0.

then for any $\eta < \eta_0$ and points $x_1, x_2 \in B_{M\eta}$ with $d(x_1, \Omega^c), d(x_2, \Omega^c) > \eta$, there exists a curve $\gamma \subseteq B_1 \cap \{d(x, \Omega^c) > c\eta\}$ of length at most $C\eta$ connecting x_1 and x_2 .

Proof. First, take $x \in \Omega$ and $d = d(x, \Omega^c)$: then

$$d|D^2u(x)| \leq C \left[\sup_{B_{d/2}(x)} |\nabla u| + d^2[f]_{C^{0,1}} \right]$$

from standard elliptic estimates and assumption (4). If $B_{d/2}(x) \subseteq B_1$, combining with (2) gives $|D^2u| \leq Cd^{-1}$.

Fix the matrix $b_{ij} = a_{ij}(0)$, and compute

$$b_{ij}\partial_i\partial_j u = f + (a_{ij} - b_{ij})\partial_i\partial_j u.$$

On $B_r \cap \{d(x, \Omega^c) > \delta_0 r\}$, this has

$$|b_{ij}\partial_i\partial_j u| \leq 1 + \delta_0^{-1}.$$

There exists a linear map $L : \mathbb{R}^n \rightarrow \mathbb{R}^n$ such that $B_\kappa \subseteq L(B_1) \subseteq B_{\kappa^{-1}}$ and if $b_{ij}\partial_i\partial_j u = g$, then $\Delta(u \circ L) = g \circ L$, where $\kappa = \kappa(\lambda)$. Let v be given by

$$v(x) = \frac{u(rLx)}{\kappa^{-1}r},$$

where r will be chosen, but at least $r < \kappa$. Let $V = L^{-1}(\Omega)/r$. Then $|\nabla v| \leq 1$ on $B_1 \cap V$ and $0 \in \partial V$. Moreover, for every $x \in \bar{V} \cap B_1$ and $s < 1$,

$$\sup_{B_s(x)} v \geq m\kappa^2 s.$$

Finally,

$$|\Delta v(x)| \leq C\kappa r(1 + \delta_0^{-1})$$

on $\{d(x, V^c) > \delta_0/\kappa\} \cap B_1$.

Fix $M' = M/\kappa$ and $m' = \kappa^2 m$, and apply Lemma B.1 with parameters M', m' to v and V : then there exists an η', δ' , such that if

$$|\Delta v(x)| \leq \delta'$$

on $\{d(x, V^c) > \delta'\} \cap B_1$, and any $y_1, y_2 \in B_{M'\eta'}$ with $d(y_i, V^c) > \eta'$, there is a path connecting y_1 to y_2 of length at most C and staying $c\eta'$ away from V^c .

Select $\delta_0 = \kappa\delta'$, and $r = \kappa\eta/\eta'$. Choose η_0 so that $r < \kappa$ and

$$|\Delta v(x)| \leq C\kappa r(1 + \delta_0^{-1}) \leq C\kappa^2 \frac{\eta_0}{\eta'} (1 + \delta_0^{-1}) < \delta'$$

on $\{d(x, V^c) > \delta_0/\kappa\} \cap B_1$. Then set $y_i = rLx_i$: we have that $y_i \in B_{M'\eta'}$ and $d(y_i, V^c) \geq \eta'$, so the above applies to give a path γ' connecting y_1 to y_2 . Let $\gamma = r^{-1}L^{-1}\gamma'$; then γ connects x_1 and x_2 , has length at most $Cr \leq C\eta$, and stays a distance $cr\eta' \geq c\eta$ from $\partial\Omega$. \square

By iterating this lemma finitely many times, we can now show that Ω satisfies the Harnack chain condition.

Theorem B.3. *Let $\Omega \subseteq \mathbb{R}^n$ be an open set, and u be a continuous function on $\bar{\Omega} \cap B_1$ (which is C^2 on Ω), with $u = 0$ on $\partial\Omega$ and $u > 0$ on Ω . Assume that*

- (1) $0 \in \partial\Omega$.
- (2) $|\nabla u| \leq 1$ on Ω .
- (3) For all $x \in \bar{\Omega} \cap B_1$ and $r < 1$, $\sup_{B_{r(x)}} u \geq mr$.
- (4) $a_{ij}\partial_i\partial_j u = f$ on $\Omega \cap B_1$, with $\|f\|_{C^{0,1}} \leq 1$, where $\lambda|\xi|^2 \leq a_{ij}(x)\xi_i\xi_j \leq \lambda^{-1}|\xi|^2$ is a matrix-valued function with $[a_{ij}]_{C^{0,1}} \leq 1$.
- (5) Ω satisfies the inner and outer clean ball condition with constant K at every point in $\partial\Omega \cap B_1$.

Then there is a $\delta > 0$ such that Ω satisfies the Harnack chain condition with constant C at any points $x_1, x_2 \in B_\delta \cap \Omega$. The constants δ, C depend only on m, λ, K and n .

One may check that the *inner* clean ball condition follows from assumptions (2-3). The outer clean ball condition, however, does not.

Proof. Let y_i be a point in $\partial\Omega$ with $d_i = d(x_i, \partial\Omega) = |x_i - y_i|$. Without loss of generality, assume $d_1 \leq d_2$. We will first show that the conclusion holds if $|x_1 - x_2| \leq 4Kd_2$. Note that if $|x_1 - x_2| < \frac{1}{2}d_2$, we may connect x_1 and x_2 by a line segment and the conclusion is immediate.

Construct a sequence of points z_k as follows, using the clean inner ball property: $z_1 = x_1$, and then given z_k , $d(z_{k+1}, \partial\Omega) \geq 2d(z_k, \partial\Omega)$ while $|z_{k+1} - z_k| \leq 4Kd(z_k, \partial\Omega)$. We continue constructing such points until $d(z_k, \partial\Omega) \geq d_2$; let the final point be z_J . Each pair z_k, z_{k+1} lies inside $B_{CK\eta_k}(p_k)$, where $\eta_k = d(z_k, \partial\Omega)$ and $p_k \in \partial\Omega$ is a point with $\eta_k = |z_k - p_k|$. They also have $d(z_k, \partial\Omega), d(z_{k+1}, \partial\Omega) \geq \eta_k$. Applying Lemma B.2 with $M = CK$ centered around p_k , we see that as long as $\eta_k < \eta_0$ for the $\eta_0 = \eta_0(K, m, \lambda)$ there, the points z_k and z_{k+1} may be connected by a curve γ_k of length at most $C\eta_k$ staying at least $c\eta_k$ away from $\partial\Omega$. Choose δ small enough that

$$\eta_k \leq \eta_J \leq Cd_2 \leq C\delta \leq \eta_0$$

and the assumption is verified for every k . Finally, apply Lemma B.2 one last time to connect z_J to x_2 by a similar curve.

We claim γ , the concatenation of all of these curves γ_k , works for the Harnack chain property. Indeed, we have that $\eta_k \leq \eta_J 2^{k-J}$ by construction, so the total length is controlled by $\sum \eta_k \leq 2\eta_J \leq Cd_2$, which is comparable to $|x_1 - x_2|$. By the same argument, the total length of the first k concatenated curves γ_k is bounded from above by $C\eta_k$, so for any point $z = \gamma(t)$ on γ_k , we have $d(z, \partial\Omega) \geq c\eta_k \geq cd(\gamma([0, t]))$.

This leaves only the case of $|x_1 - x_2| \geq 4Kd_2$. Pick a point $x_3 \in \Omega$ with $d(x_3, \partial\Omega) \geq |x_1 - x_2|$ and $|x_3 - x_2| \leq K|x_1 - x_2|$, using the clean inner ball property. Then $|x_3 - x_2| \leq 2K|x_1 - x_2| \leq 2Kd(x_3, \partial\Omega)$ and similarly $|x_3 - x_1| \leq 3Kd(x_3, \partial\Omega)$, so we may connect x_3 to x_2 and x_1 by curves satisfying the Harnack chain property. The concatenation of these curves then works to show that Ω satisfies the Harnack chain condition at x_1 and x_2 . \square

A direct application of this theorem proves Lemma 7.5.

Proof of Lemma 7.5. Apply Theorem B.3 to $u(x) = m_0[w_\Omega(r_0x) + \sqrt{\mathfrak{I}}u_\Omega(r_0x)]$ in normal coordinates around a point $x \in \partial\Omega$. By choosing r_0 and m_0 small, we have property (4) from the regularity of the metric. Property (2) follows from Corollary 5.5, while (3) comes from Theorem 4.1. The inner and outer ball conditions were verified in Lemma 7.1. \square

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